

Localized States in Engineered Carbon Nanoribbons

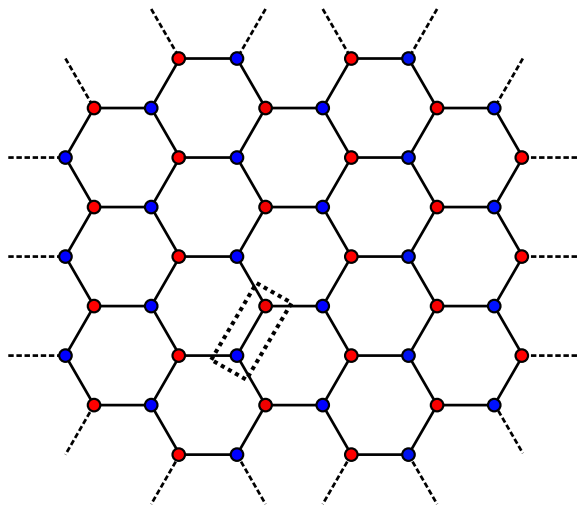
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September 26, 2022

Introduction

Graphene



Introduction

Graphene Nanoribbons

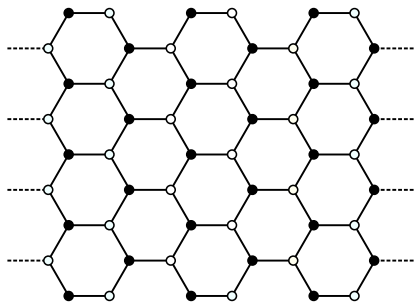


Figure: Armchair nanoribbon (ANR)

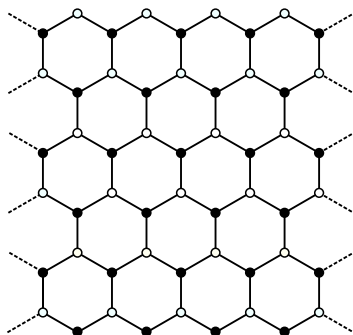


Figure: Zigzag nanoribbon (ZNR)

Introduction

Dispersion Relations: Armchair[5]

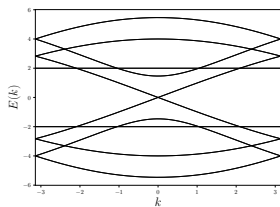


Figure: $W = 5$

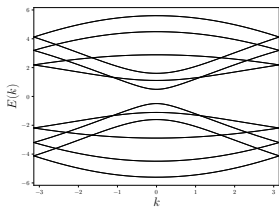


Figure: $W = 6$

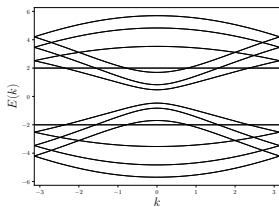
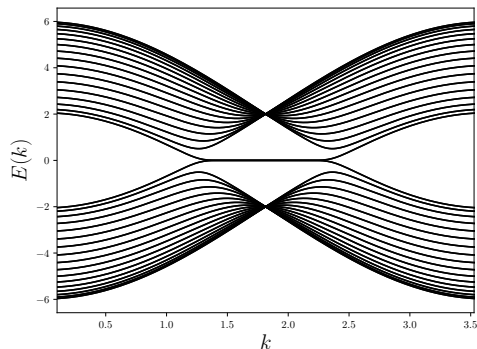


Figure: $W = 7$

- Dispersion depends on widths
- Zero crossing for $3m - 1$
- Gap $\sim 1/W$

Introduction

Dispersion Relations: Zigzag[5]



- No width dependence
- Flatband in the middle third of the BZ
- Exponentially localized edge states

Hybrid Armchair Nanoribbon

Reasoning

- Less covered in literature (arXiv search results: armchair/zigzag $\approx 1/2$)
- Can be synthesized from ground up [2]
- Has been shown to have topological properties

Termination type	Zigzag ($N = \text{Odd}$)	Zigzag' ($N = \text{Odd}$)	Zigzag ($N = \text{Even}$)	Bearded ($N = \text{Even}$)
Unit cell shape				
Bulk Symmetry	Inversion/mirror	Inversion/mirror	Mirror	Inversion
Z_2	$\frac{1 + (-1)^{\lfloor \frac{N}{3} \rfloor + \lfloor \frac{N+1}{2} \rfloor}}{2}$	$\frac{1 - (-1)^{\lfloor \frac{N}{3} \rfloor + \lfloor \frac{N+1}{2} \rfloor}}{2}$	$\frac{1 - (-1)^{\lfloor \frac{N}{3} \rfloor + \lfloor \frac{N+1}{2} \rfloor}}{2}$	$\frac{1 - (-1)^{\lfloor \frac{N}{3} \rfloor}}{2}$

Figure: Z_2 invariant of an armchair nanoribbon [1]

Hybrid Armchair Nanoribbon

Reasoning

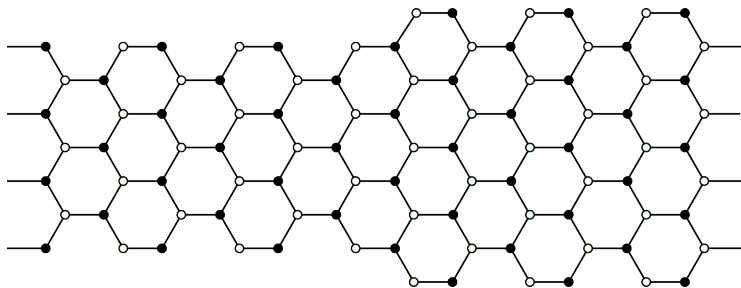


Figure: 7/9 Junction Armchair Nanoribbon

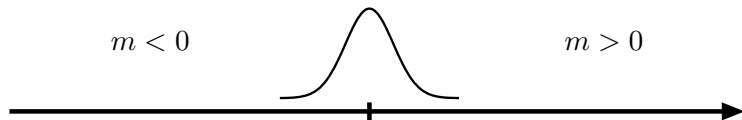


Figure: Jackiw-Rebbi Zero Mode[4]

Non-interacting Dispersion

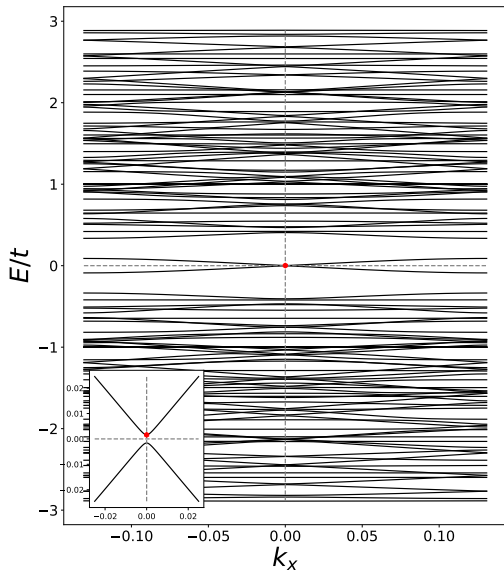


Figure: Low Lying $k_x = 0$ state is highlighted in red

Hopping + Hubbard

Monte Carlo Simulation Results

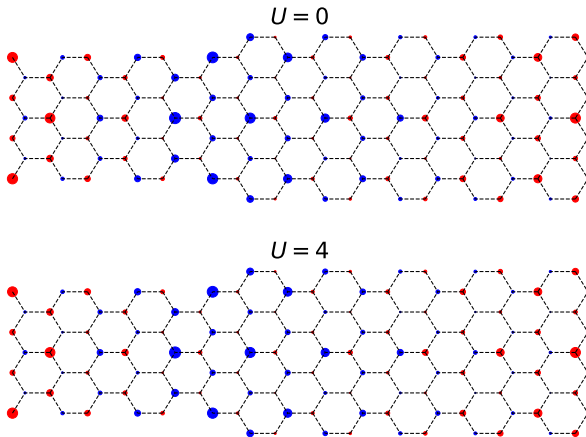


Figure: GS densities are effectively independent of U

BCS-Hubbard Hamiltonian[3]

$$H = - \sum_{\langle i,j \rangle \sigma} (t c_{i\sigma}^\dagger c_{j\sigma} + \Delta c_{i\sigma} c_{j\sigma} + h.c.) + U \sum_k \left(n_{k\uparrow} - \frac{1}{2} \right) \left(n_{k\downarrow} - \frac{1}{2} \right)$$

Symmetric lines condition ($t = \Delta$)

- For any bipartite lattice is exactly solvable
- Investigate behaviour for arbitrary U
- 2^N disjoint sectors
 - Ferromagnetic & Antiferromagnetic

Hopping + SC pairing + Hubbard

Results

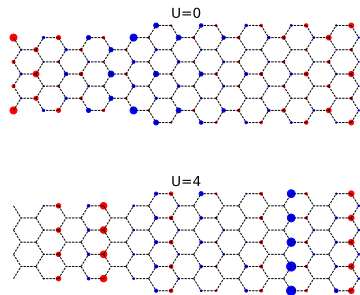


Figure: Ferromagnetic

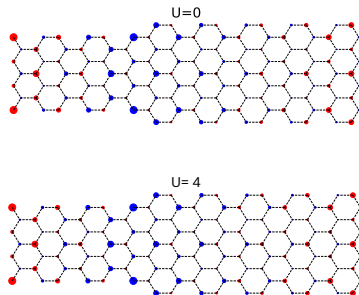


Figure: Antiferromagnetic

Hybrid ANR's

- Low-lying energy state
- Exponentially localized
- Resilient against Hubbard interaction
- Locality persists even with SC pairs

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Localization of Electronic States in Hybrid Nano-Ribbons in the Non-Perturbative Regime

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Abstract

We investigate the localization of low-energy single quasi-particle states in the 7/9- and 13/15-hybrid nanoribbon systems in the presence of strong interactions and within a finite volume. We consider two scenarios, the first being the Hubbard model at half-filling and perform quantum Monte Carlo simulations for a range U that includes the strongly correlated regime. In the second case we add a nearest-neighbor superconducting pairing Δ and take the symmetric line limit, where Δ is equal in magnitude to the hopping parameter t . In this limit the quasi-particle spectrum and wavefunctions can be directly solved for general onsite interaction U . In both cases we extract the site-dependent quasi-particle wavefunction densities and demonstrate that localization persists in these non-perturbative regimes under particular scenarios. Our findings suggest that such localization under the presence of interactions and within a finite volume is a generic feature of hybrid nanoribbons composed of topologically distinct regions.

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Thank You

Thank You

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- [1] Ting Cao, Fangzhou Zhao, and Steven G. Louie. Topological phases in graphene nanoribbons: Junction states, spin centers, and quantum spin chains. *Phys. Rev. Lett.*, 119:076401, Aug 2017.
- [2] Y. Chen, T. Cao, C. Chen, Z. Pedramrazi, D. Haberer, D. de Oteyza, F. Fischer, S. Louie, and M. Crommie. Molecular bandgap engineering of bottom-up synthesized graphene nanoribbon heterojunctions. *Nature Nanotechnology*, 10, 2015.
- [3] Zewei Chen, Xiaohui Li, and Tai Kai Ng. Exactly solvable bcs-hubbard model in arbitrary dimensions. *Phys. Rev. Lett.*, 120:046401, Jan 2018.
- [4] R. Jackiw and C. Rebbi. Solitons with fermion number $\frac{1}{2}$. *Phys. Rev. D*, 13:3398–3409, Jun 1976.
- [5] Katsunori Wakabayashi, Ken ichi Sasaki, Takeshi Nakanishi, and Toshiaki Enoki. Electronic states of graphene nanoribbons and analytical solutions. *Science and Technology of Advanced Materials*, 11(5):054504, Oct 2010.