Small neutrino masses from chiral gravitational anomaly

Lena Funcke

Massachusett Institute of Technology

In collaboration with Gia Dvali, Georg Raffelt, Tanmay Vachaspati, and others (1602.03191, 1608.08969, 1811.01991, 1905.01264, 2102.13618, and 2112.02107)

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Massive pseudoscalar	$\eta^\prime = ar q \gamma_5 q \Rightarrow \partial^\mu \eta^\prime = \dot j_5^\mu = ar q \gamma^\mu \gamma_5 q$
Equations of motion	$d\left(E-\eta' ight)=0 \Rightarrow E=\eta'+ heta$
	$\Box \eta' + E = \Box \eta' + (\eta' + \theta) = 0$

Massive field: C "eats up" pseudoscalar η' , vacuum is at E = 0

Quantity	Gravity without neutrinos
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Chern-Simons three-form	$C_G \equiv \Gamma \mathrm{d}\Gamma - \frac{3}{2}\Gamma\Gamma\Gamma$
Topological susceptibility	$\langle R\tilde{R}, R\tilde{R} \rangle_{p \to 0} = \text{const.} \neq 0$
Effective Lagrangian	$\mathcal{L} = \frac{1}{2}E_G^2 + \theta_G E_G + \text{higher orders}$

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Non-perturbative topological effects in pure gravity.











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► Effective potential allows for neutrino mass hierarchy: $V(\hat{X}) = \sum_n \frac{1}{n} c_n \operatorname{Tr}[(\hat{X}^+ \hat{X})^n]$ with $\hat{X}_{\alpha_L}^{\alpha_R} \equiv \langle \bar{\nu}_{\alpha_L} \nu_{\alpha_R} \rangle$ $\rightarrow \partial V / \partial x_i = 0$ determines $\hat{X} = \operatorname{diag}(x_1, x_2, x_3)$.

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 Mechanism works for Dirac and Majorana masses.

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[9] Archidiacono, Hannestad (2014).Image credits: NASA / WMAP Science Team [http://map.gsfc.nasa.gov/]

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\rightarrow Neutrino vacuum condensate $\langle \bar{\nu} \nu \rangle$ on dark energy scale

[9] Archidiacono, Hannestad (2014). [10] Tanabashi et al. (Particle Data Group) (2018).Image credits: NASA / WMAP Science Team [http://map.gsfc.nasa.gov/] and Patterson (2005).

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Impact on other cosmic parameters.



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 Image credit: KATRIN [http://www.ikp.kit.edu/].
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 \Rightarrow Parameter degeneracy and/or new physics?



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- $\blacktriangleright\,$ Observable at Daya Bay? $\Delta t=6$ y, $\Delta d=v\Delta t=4\times 10^{13}~{\rm m}\sim\xi$



[15] Dvali, LF, Vachaspati (2021). Image credit: Roy Kaltschmidt, Berkeley Lab.

Astrophysical neutrinos:

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[16] Jackiw, Pi (2003). Image credits: The SXS Project [https://www.ligo.caltech.edu/]

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[16] Jackiw, Pi (2003). [17] Dvali, LF (2016b), "Domestic Axion" solution to strong CP problem. Image credits: The SXS Project [https://www.ligo.caltech.edu/] and Kim, Carosi (2008).

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Thanks for listening!



Do you have any questions?

Backup - String-Wall Network in Majorana Case

Topological defects

• Inter-string separation: $\xi = 10^{14} \text{ m} \left(\frac{\lambda}{1}\right) \left(\frac{\Lambda_G}{1 \text{ meV}}\right)^{\frac{1}{2}} \left(\frac{1}{a_G}\right)$

• Domain wall width:
$$\delta_{\text{DW}} = \frac{1}{m_{\phi_k}} = 8 \times 10^{14} \text{ m} \left(\frac{m_{\tau}}{m_l} \right) \left(\frac{1 \text{ meV}}{\Lambda_G} \right)^2$$

Impact on neutrino mixing angles

- Winding: $\langle \Phi(\theta) \rangle = \omega^{T}(\theta) \langle \Phi \rangle \omega(\theta) \neq \langle \Phi(0) \rangle$ changes flavor
- ► Time scale: $\Delta(\theta_{ij}, \alpha_k, \delta) = O(1)$ after $\Delta t = \xi/v$, v = 230 km/s
- ▶ Example: $\Delta t = d/\nu \sim 1$ h for $\xi < \delta_{\rm DW} = 8 \times 10^8$ m, $\Lambda_G = 1$ eV

Experimental constraints

- Daya Bay data: sin²(2θ₁₃) = 0.0856 ± 0.0029, angle could exhibit time-dependent change of sin²(2 × 2πνt/δ_{DW}) ≤ 2 × 0.0029
- ▶ Still viable: $\xi < \delta_{\rm DW}$ with $\Lambda_G \lesssim 0.2 \text{ meV}$ and $\xi > \delta_{\rm DW}$

Dvali, LF, Vachaspati (2021).

Backup – Binned Reconstruction of Neutrino Masses



Image credit: Lorenz, LF, Löffler, Calabrese (2021).

Backup – Reconstruction with Wider Redshift Bins



Image credit: Lorenz, LF, Löffler, Calabrese (2021).

Backup – Reconstruction with Linear Splines

Mass reconstruction with Bayesian regression splines and variable knots \Rightarrow Two new parameters: the knots (change points) z_1 and z_2



Image credit: Lorenz, LF, Löffler, Calabrese (2021).

Backup – Dark Energy and Neutrino Masses

Impose prior on Ω_{Λ} to test degeneracy between Ω_{Λ} and m_{ν}

ightarrow Constraint on $\sum m_
u(z=0)$ decreases by 42% but is still large



Backup – Constrain Extended Neutrino Models

Reconstruction results can be converted into neutrino mass bound for neutrino decay models [1]: $\sum m_{\nu} < 0.21$ eV (95% CL) [2]

ightarrow Bound too low for possible KATRIN detection of $m_eta \ge$ 0.2 eV



[1] e.g., Chacko et al. (2019), [2] Lorenz, LF, Löffler, Calabrese (2021).

Backup – Data Sets and Redshift Ranges

Data set	Redshift range
Planck 2018 CMB TTTEEE	mostly $z = 1100$
Planck 2018 CMB lowl	mostly $z = 1100$
Planck 2018 CMB lowE	mostly $z = 8,1100$
Planck 2018 CMB lensing	$0 \le z \le 1100$
BAO (6dF)	z = 0.106
BAO (SDSS DR7 BOSS MGS)	<i>z</i> = 0.15
BAO (SDSS DR12 BOSS)	z = 0.38, 0.51, 0.61
BAO (SDSS DR14 eBOSS quasars)	z = 1.52
BAO (SDSS DR14 eBOSS Ly- $lpha$)	<i>z</i> = 2.34
BAO (SDSS DR14 eBOSS cross Ly- α -QSO)	z = 2.35
SN (Pantheon)	0.01 < z < 2.3

Backup – Three-form Higgs Effect

- ► Chern-Pontryagin density: $E_G \equiv R\tilde{R} = \varepsilon^{\mu\nu\rho\sigma} R^{\alpha}_{\beta\mu\nu} R^{\beta}_{\alpha\rho\sigma} = dC_G$
- Chern-Simons three-form: $C_G \equiv \Gamma d\Gamma \frac{3}{2}\Gamma\Gamma\Gamma$

► Effective Lagrangian for C_G : $\mathcal{L} = \frac{1}{2\Lambda_G^4} E_G^2 + \theta_G E_G + \text{higher orders [3]}$ ► With neutrino: $\mathcal{L} = \frac{1}{2\Lambda_G^4} E_G^2 - \frac{1}{\Lambda_G} \eta_\nu E_G + \frac{1}{2} \partial_\mu \eta_\nu \partial^\mu \eta_\nu, \ \eta_\nu = \frac{1}{\Lambda_G^2} \bar{\nu} \gamma_5 \nu$

► Equation of motion for η_ν: □η_ν + 1/Λ_G E_G = 0, in vacuum: RR̃ = 0.
 ► Equation of motion for C_G: (□ + Λ²_G) E_G = 0, mass gap generated.

[3] Dvali (2005); Dvali, Jackiw, and Pi (2006); Dvali, Folkerts, and Franca (2014).

Backup – Formation of Condensate

- ► High-energy formulation: vacuum angle θ_G made unobservable due to arbitrary shift by chiral rotation of massless neutrino field
- Low-energy formulation: $E_G = R\tilde{R}$ screened in vacuum by massless neutrino, as C_G "eats up" pseudoscalar η_{ν} and becomes massive

► Integrated e.o.m. of E_G : $\langle R\tilde{R} \rangle_{q \to 0} \simeq -\theta_G m_\nu \Lambda_G^3 = -\theta_G m_\nu \langle \bar{\nu}\nu \rangle$ ► From Lagrangian: $\mathcal{L} = \frac{1}{2\Lambda_+^3} E_G^2 - \frac{1}{\Lambda_C} \eta_\nu E_G + \frac{1}{2} \partial_\mu \eta_\nu \partial^\mu \eta_\nu - \frac{1}{2} m_\nu \Lambda_G \eta_\nu^2$

► Effective potential: $V(\hat{X}) = \sum_{n} \frac{1}{n} c_{2n} \operatorname{Tr}[(\hat{X}^+ \hat{X})^n], \ \hat{X}_{\alpha_L}^{\alpha_R} \equiv \langle \bar{\nu}_{\alpha_L} \nu_{\alpha_R} \rangle$ ► Extrema: $\frac{\partial V}{\partial x_j} = x_j^* \left(\sum_n c_{2n} |x|_j^{2(n-1)} \right) = 0 \text{ for } \hat{X} = \operatorname{diag}(x_1, x_2, x_3)$

Dvali, LF (2016a).

Backup – Structure of Fermion Condensates

 Flavor structure of condensate determined by minimization of effective potential for following order parameters:

$$\hat{X}_{i\bar{j}} \equiv \psi_i \psi_{c\bar{j}} , \ X_{ij} \equiv \psi_i C \psi_j , \ \bar{X}_{\bar{i}\bar{j}} \equiv \psi_{c\bar{i}} C \psi_{c\bar{j}} ,$$

- ► 1) Flavor-invariant terms, Tr(X⁺X), Tr(X⁺XX⁺X), ..., Tr(X⁺X), Tr(X⁺X), Tr(X⁺X), Tr(XX⁺XX), In effective potential: infinite polynomial of invariants scaled by powers of Λ_G.
- ► 2) Explicitly flavor-breaking terms, break anomalous U(1)_A, leave invariant anomaly-free subgroup Z_{2N_F}. Operators:

$$\epsilon^{i_1 \cdots \, i_{N_F}} \epsilon^{\overline{j}_1 \cdots \, \overline{j}_{N_F}} \hat{X}_{i_1 \overline{j}_1} \cdots \hat{X}_{i_{N_F} \overline{j}_{N_F}}$$

► Effective potential invariant under symmetry group SU(N_F)× SU(N_F) × U(1)_V × Z_{2N_F}, spontaneously broken by condensate.

Backup – Gravitational Wave Propagation

- ▶ Non-dynamical Chern-Simons modification of GR, proposed in [16]
- ▶ Angle θ_G assumed to be time-dependent, non-dynamical quantity
- ► Deformation of equations of motion of GR: $G_{\mu\nu} + C_{\mu\nu} = -8\pi G T_{\mu\nu}$, where $D_{\mu}C^{\mu\nu} = \frac{1}{8\sqrt{-g}}v^{\nu}R\tilde{R}$, $v_{\nu} \equiv \partial_{\mu}\theta_{G}$, and $v_{\nu} = (1/\mu, \mathbf{0})$
- ► Diffeomorphism breaking $(D_{\mu}C^{\mu\nu} \neq 0)$ dynamically suppressed, Schwarzschild solution persists, gravitational waves travel at *c*

► Radiated power P per unit angle Ω for different polarizations:

$$\frac{dP_{\pm}}{d\Omega} \propto \left(1 \pm \frac{k}{\mu}\right)^{-2}$$

Strong *CP* problem [1]: $\mathcal{L}_{\rm QCD} \supset \theta G \tilde{G}, \ \bar{\theta} < 10^{-10}.$



[1] Recent review: Kim (2016).









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Backup – The CP Problem of QCD



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Backup – Generation of Light Fermion Masses







[2] Dvali, LF (2016a). [3] Dvali, LF (2016b).

Backup – Generation of Light Fermion Masses

- ▶ Neutrino [2] and up quark [3] mass generation through non-perturbative coupling to neutrino condensate.
- ▶ Coupling analogous to 't Hooft vertex in QCD [4].





[2] Dvali, LF (2016a). [3] Dvali, LF (2016b). [4] 't Hooft (1986).

Backup – Generation of Light Fermion Masses

- Neutrino [2] and up quark [3] mass generation through non-perturbative coupling to neutrino condensate.
- ▶ Coupling analogous to 't Hooft vertex in QCD [4].
- No new scale or additional tuning: m_u ~ ξ⟨νν⟩, highmultiplicity parameter ξ ~ 10⁷ replaces Y_u ~ 10⁻⁵ [3].





[2] Dvali, LF (2016a). [3] Dvali, LF (2016b). [4] 't Hooft (1986).

Neutrino and up-quark masses generated by IR $\langle \bar{\nu}\nu \rangle \equiv$ "second Higgs".



Neutrino and up-quark masses generated by IR $\langle \bar{\nu}\nu \rangle \equiv$ "second Higgs". $U(1)_{A\nu}$ and $U(1)_{Au}$ mix to $U(1)_{PQ}$ and $U(1)_{G}$.

















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Flavor-violation by infrared "Higgs boson" suppressed.