



# Associated production of quarkonium in ATLAS

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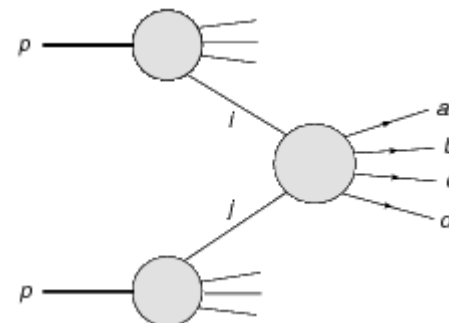
# `New observables' in quarkonium production

- ❑ Despite 40+ years' history, we still have no clear and reliable picture of quarkonium production in hadronic -- and other -- collisions
- ❑ New energy frontiers and higher luminosities at LHC allow exploration of other reactions that may help understand better the dynamics of quarkonium production
- ❑ Simply speaking, more equations (experimental constraints) may help determine unknowns better, even if some new unknowns are introduced
- ❑ Examples of these `new observables': associated production of quarkonium with other objects, such as:
  - other quarkonium (LHCb, CMS, now ATLAS)
  - W or Z bosons ( ATLAS )
  - others to come ?

The production of two objects in the same pp collision can be due to

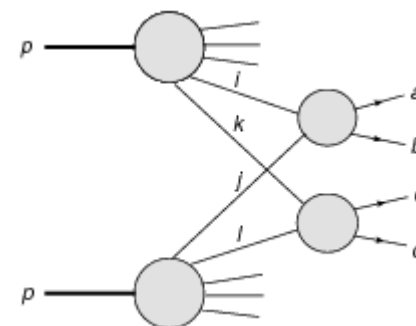
## Single-Parton Scattering (SPS):

the two objects are produced via a subprocess in a single interaction of two partons



## Double-Parton Scattering (DPS):

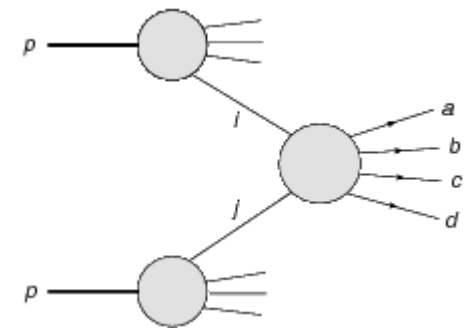
simultaneous interaction of two pairs of partons, each producing one of the two objects, assumed to be uncorrelated



*E. Berger et al. Phys.Rev. D84 (2011) 074021 arXiv:1107.3150 [hep-ph]*

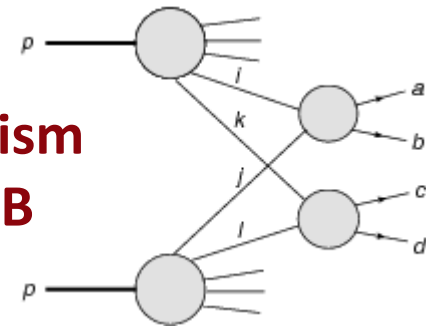


- Single-Parton Scattering (SPS) can be treated in the usual way, once the subprocess cross section is calculated
- After all, this is just “another subprocess”
- If the final particles are colourless, all the usual factorisation theorems are valid
- (at least as valid as for other subprocesses, may be more)
- Since the initial partons are unlikely to have large transverse momenta, the final state objects are expected to be produced back-to-back in transverse plane
- (up to some smearing due to initial- and final-state radiation etc)





- Double-Parton Scattering (DPS) can be treated in as two subprocesses happening simultaneously in the same pp collision
- Assuming the masses and momenta are much smaller than collision energy, one can ignore correlations due to overall Energy-momentum conservation
- So one can apply the usual parton model / QCD formalism TWICE, once for subprocess A and once for subprocess B
- The problem is the cross sections are dimensional, and one needs a dimensional factor  $\sigma_{\text{eff}}$  as a scale



$$\sigma_{A+B}^{\text{DPS}} = \frac{1}{1 + \delta_{AB}} \frac{\sigma_A \sigma_B}{\sigma_{\text{eff}}}$$

$\sigma_{\text{eff}} \sim (2 - 20) \text{ mb}$ , assumed (hoped?) to be independent of process and  $\sqrt{s}$



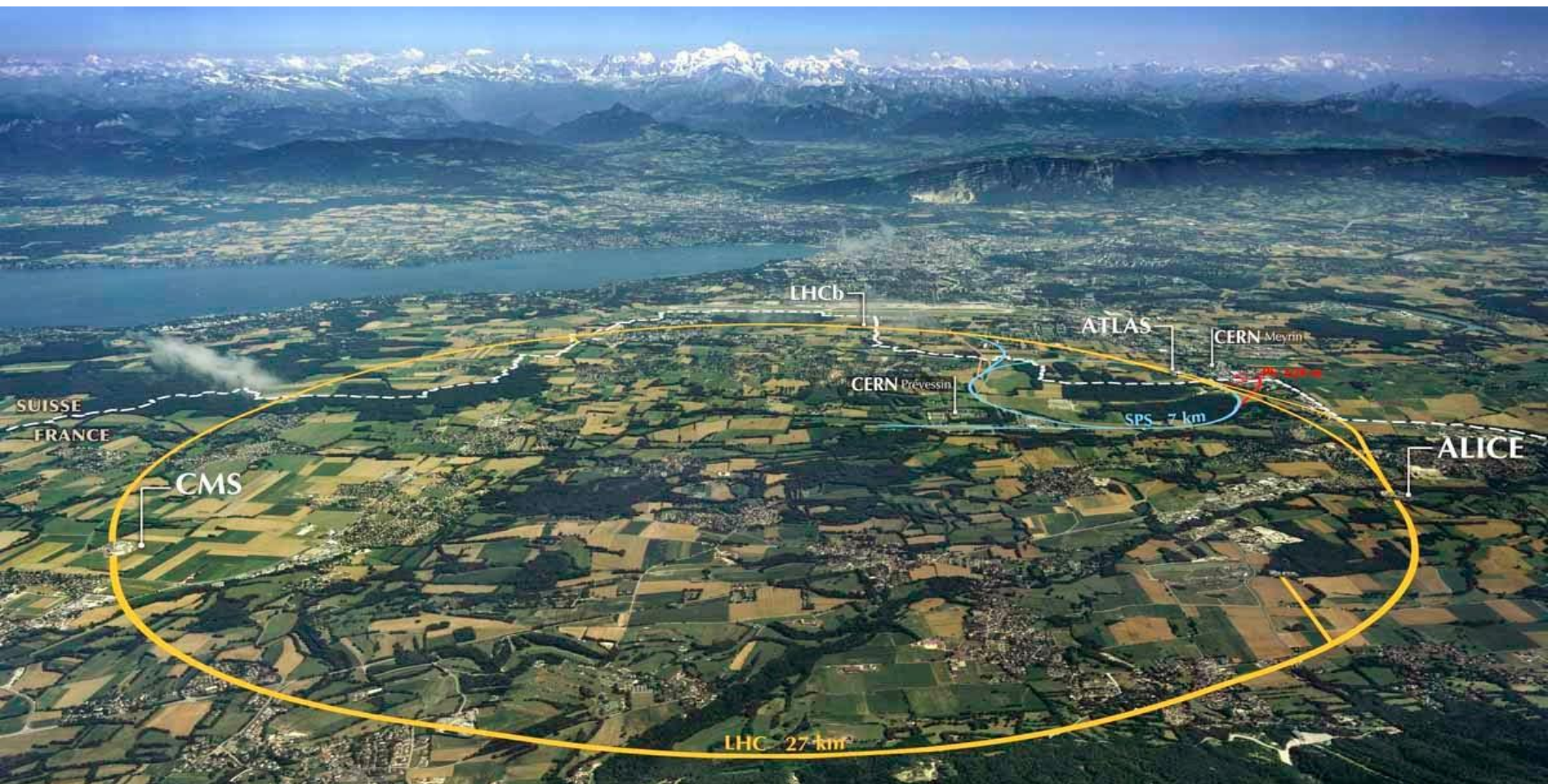
# Separating DPS and SPS

- ☐ DPS not distinguishable from SPS on event-by-event basis
- ☐ However, they are expected to differ in kinematic features, such as angular correlations, so can be separated on a statistical basis
- ☐ But there usually are large uncertainties in separation:
  - possible higher-order SPS contributions and feed-down
  - limited knowledge of proton's transverse profile
- ☐ At the end of the day, both SPS and DPS are part of the same QM amplitude, so need to be careful with things like unitarity and gauge invariance





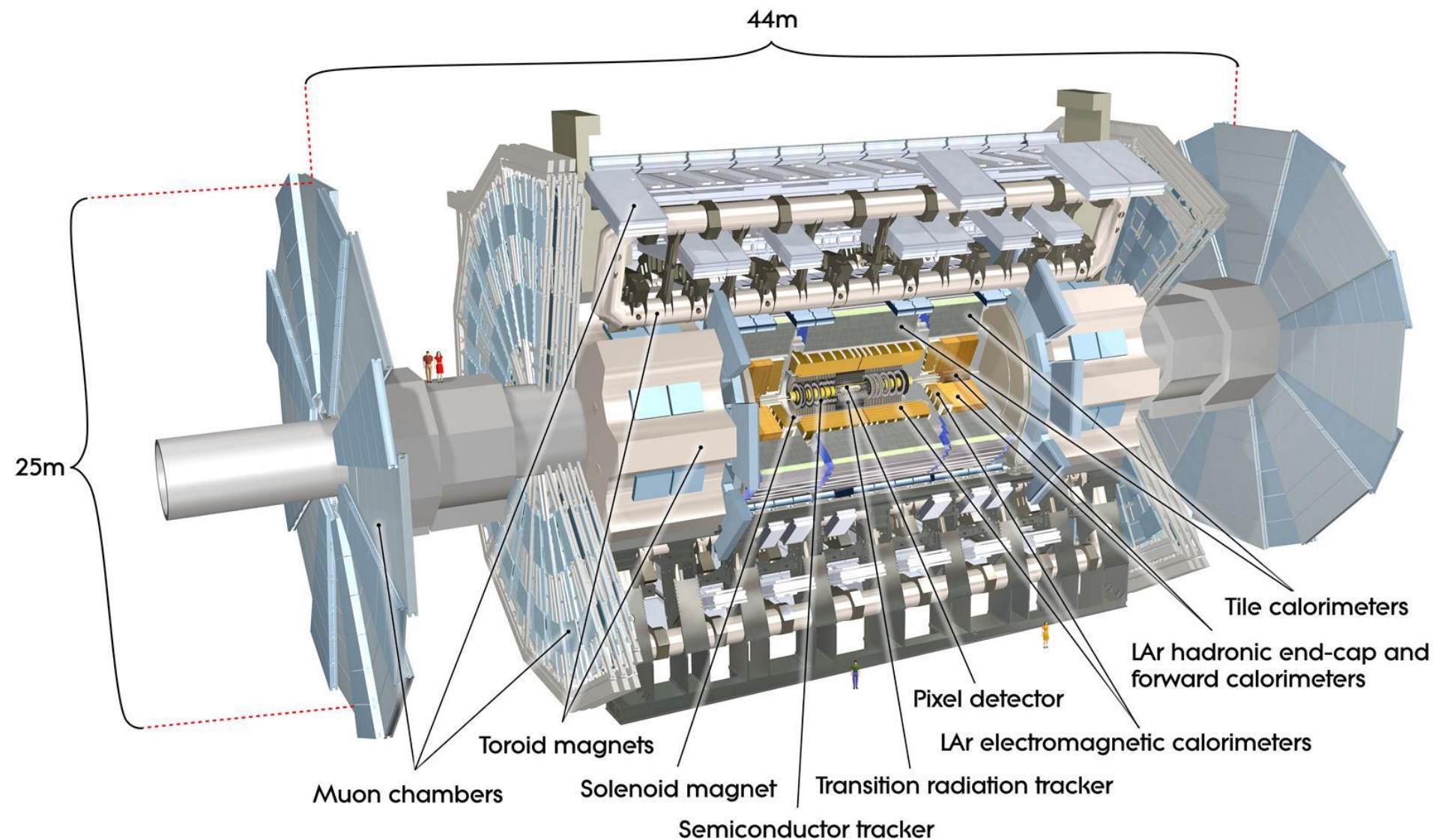
# ATLAS: experimental facility







# The ATLAS detector







## $J/\psi + W^\pm$

**Measurement of the production cross section of prompt  $J/\psi$  mesons in association with a  $W^\pm$  boson in pp collisions at  $\sqrt{s}=7$  TeV with the ATLAS detector**

JHEP 04 (2014) 172

arXiv:1401.2831

## $J/\psi + Z^0$

**Observation and measurements of the production of prompt and non-prompt  $J/\psi$  mesons in association with a Z boson in pp collisions at  $\sqrt{s} = 8$  TeV**

Eur. Phys. J. C75 (2015) 229

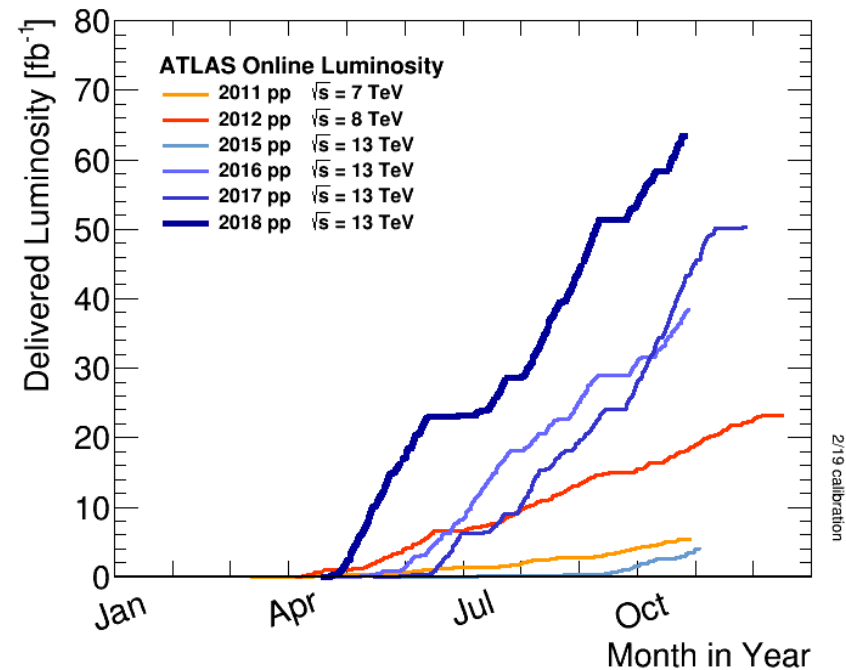
arXiv:1412.6428

## $J/\psi + J/\psi$

**Measurement of the prompt  $J/\psi$  pair production cross-section in pp collisions at  $\sqrt{s} = 8$  TeV with the ATLAS detector**

Eur. Phys. J. C77 (2017) 76

arXiv:1612.02950





$J/\psi + W^\pm$

Trigger: single muon,  $p_T > 18$  GeV

$$\begin{aligned} \sqrt{s} &= 7 \text{ TeV} & \text{fiducial phase space } & 8.5 < p_T^{J/\psi} < 30 \text{ GeV} & |y^{J/\psi}| < 2.1 \\ \mathcal{L} &= 4.51 \text{ fb}^{-1} & p_T^\mu > 3.5 \text{ GeV} & |\eta^\mu| < 1.3 & |\eta^\mu| < 2.5 & \text{at least one } p_T^\mu > 4 \text{ GeV} \\ J/\psi &\rightarrow \mu^+ \mu^- & p_T^\mu > 2.5 \text{ GeV} & |\eta^\mu| > 1.3 & p_T^{\mu(W)} > 25 \text{ GeV} & |\eta^{\mu(W)}| < 2.4 \\ W^\pm &\rightarrow \mu \nu_\mu \end{aligned}$$

$J/\psi + Z^0$

Trigger: single muon or electron,  $p_T > 24$  GeV

$$\begin{aligned} \sqrt{s} &= 8 \text{ TeV} & \text{fiducial phase space } & 8.5 < p_T^{J/\psi} < 100 \text{ GeV} & |y^{J/\psi}| < 2.1 \\ \mathcal{L} &= 20.3 \text{ fb}^{-1} & p_T^\mu > 3.5 \text{ GeV} & |\eta^\mu| < 1.3 & |\eta^\mu| < 2.5 \\ J/\psi &\rightarrow \mu^+ \mu^- & p_T^\mu > 2.5 \text{ GeV} & |\eta^\mu| > 1.3 & \text{at least one } p_T^\mu > 4 \text{ GeV} \\ Z &\rightarrow \ell\ell, \ell = \mu, e & p_T^{\mu(Z)} > 15 \text{ GeV} & |\eta^{\mu(Z)}| < 2.5 \\ & & p_T^{e(Z)} > 15 \text{ GeV} & |\eta^{e(Z)}| < 2.47 \end{aligned}$$

$J/\psi + J/\psi$

Trigger: 2 muons,  $p_T > 4$  GeV, around  $J/\psi$  mass

$$\begin{aligned} \sqrt{s} &= 8 \text{ TeV} & \text{fiducial phase space } & p_T^{J/\psi} > 8.5 \text{ GeV} & |y^{J/\psi}| < 2.1 \\ \mathcal{L} &= 11.4 \text{ fb}^{-1} & p_T^\mu > 2.5 \text{ GeV} & |\eta^\mu| < 2.3 \end{aligned}$$



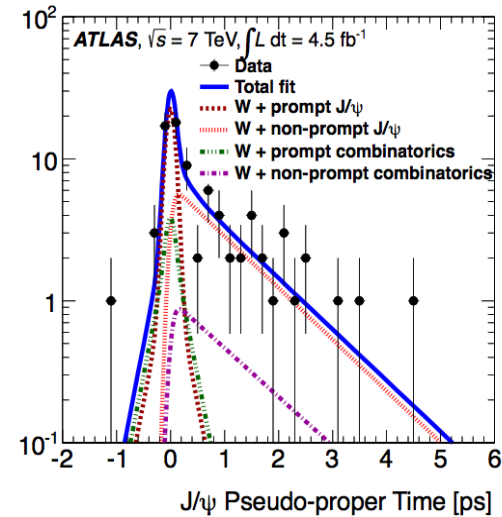
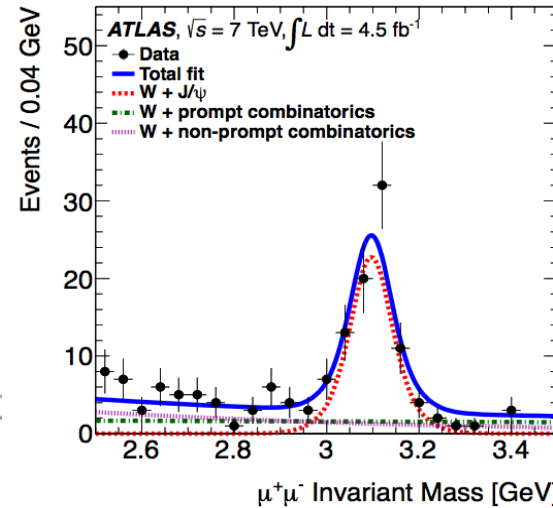
# J/ψ + W<sup>±</sup> : selection and yields

Unbinned maximum likelihood (ML) fit  
in the  $J/\psi$  invariant mass and  
pseudo-proper time  $\tau$  to obtain yields for  
prompt and non-prompt  $J/\psi$  and background

Assign weights with sPlot

arXiv:physics/0402083

$$\tau \equiv \frac{\vec{L} \cdot \vec{p}_T^{J/\psi}}{p_T^{J/\psi}} \frac{m_{\mu^+\mu^-}}{p_T^{J/\psi}}$$



Fit templates to the weighted W boson  
transverse mass to extract

signal yield:  $29.2^{+7.5}_{-6.5}$

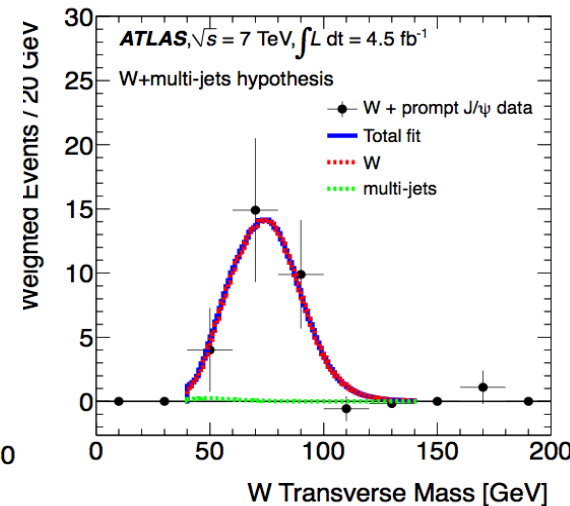
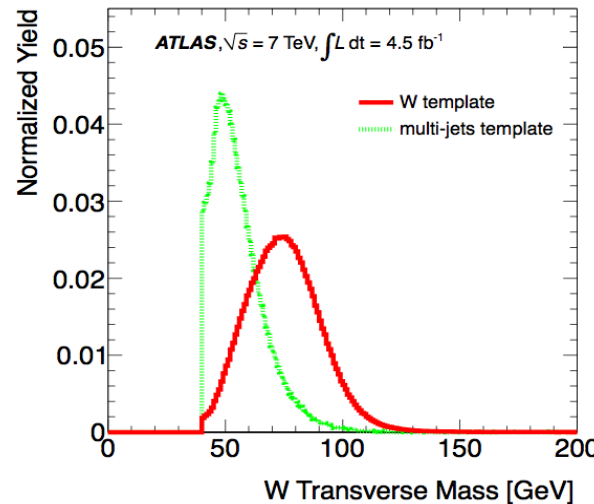
including  $1.8 \pm 0.2$  from pile-up:

DPS yield:  $10.8 \pm 4.2$

assuming  $\sigma_{\text{eff}} = 15 \pm 3$  (stat)  $^{+5}_{-3}$  (syst) mb

arXiv:1301.6872

and  $\sigma_{J/\psi}$  from arXiv:1104.3038





# J/ψ + W<sup>±</sup> : results

Ratios of the W + J/ψ prompt cross section to the inclusive W cross section

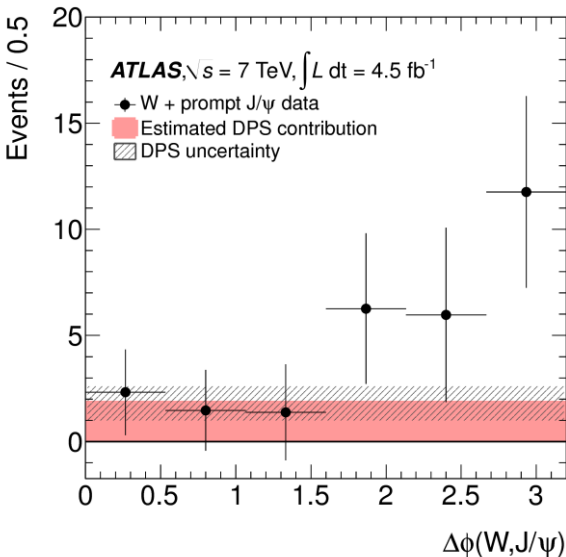
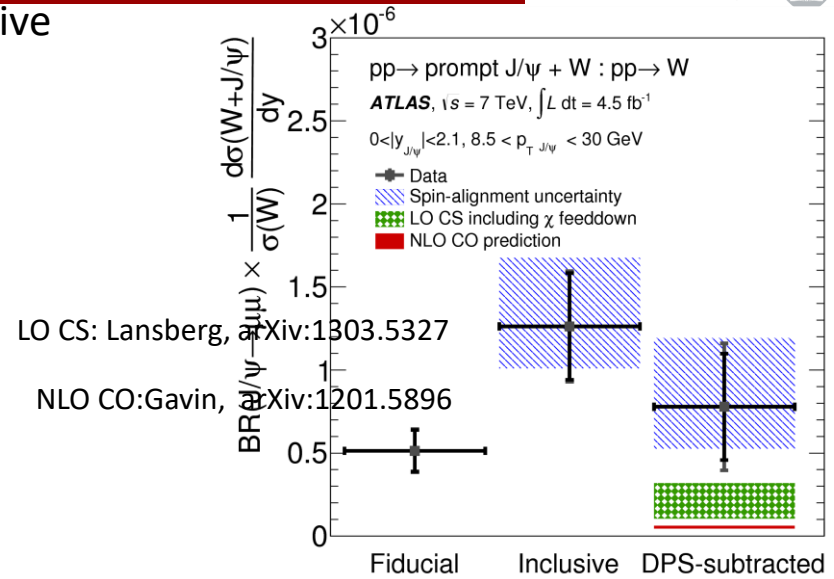
fiducial  $R_{J/\psi}^{\text{fid}} = (51 \pm 13 \pm 4) \times 10^{-8}$

inclusive  $R_{J/\psi}^{\text{incl}} = (126 \pm 32 \pm 9^{+41}_{-25}) \times 10^{-8}$

corrected for the fiducial acceptance of the muons from J/ψ  
isotropic spin-alignment assumed  
last uncertainty from variations with 5 extreme scenarios

DPS subtracted  $R_{J/\psi}^{\text{DPS sub}} = (78 \pm 32 \pm 22^{+41}_{-25}) \times 10^{-8}$

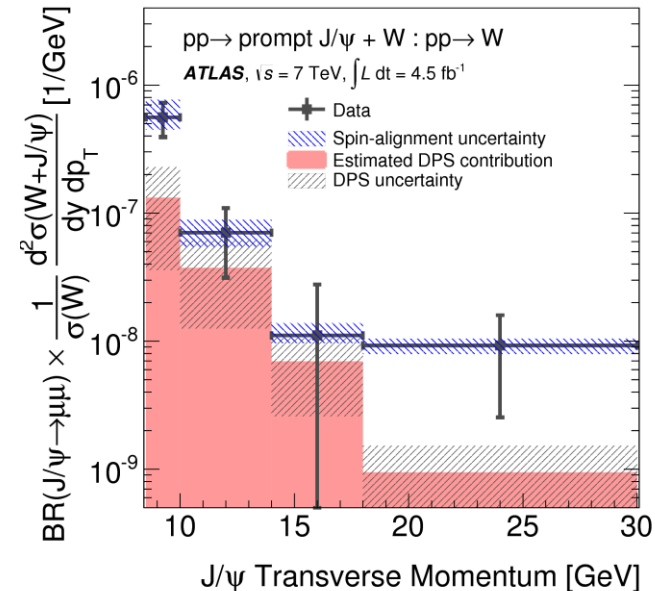
W + J/ψ dominated by CS production



Presence of both SPS and DPS contributions

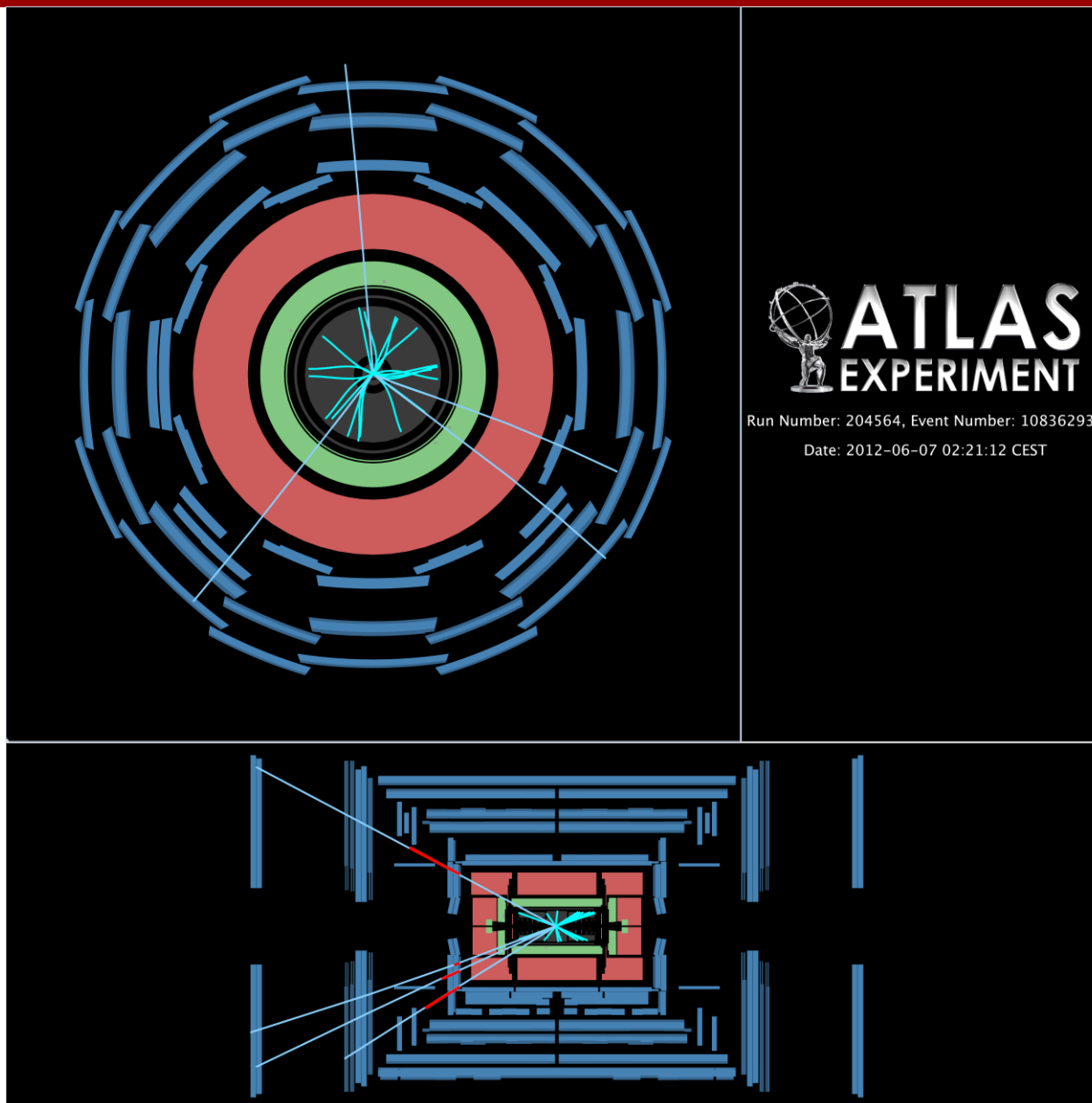
Inclusive differential cross section ratio

SPS is the dominant contribution to the total rate at low p<sub>T</sub>





# $J/\psi + Z^0$ : event candidate

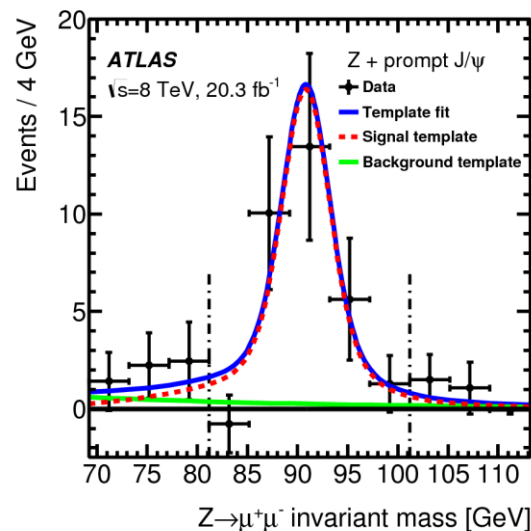
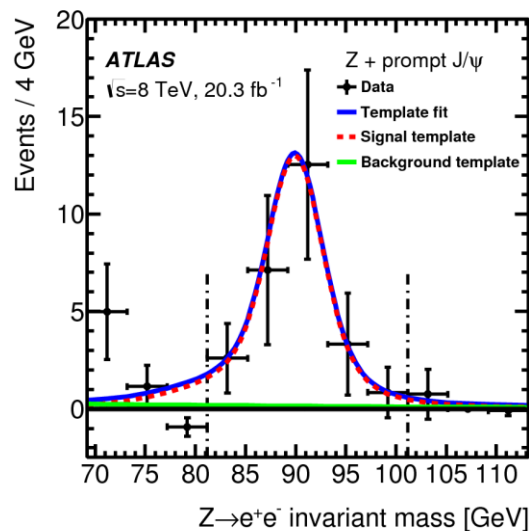
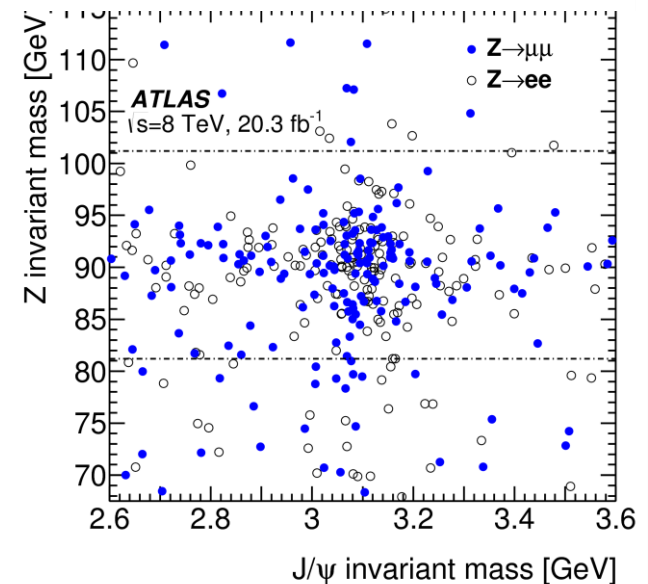
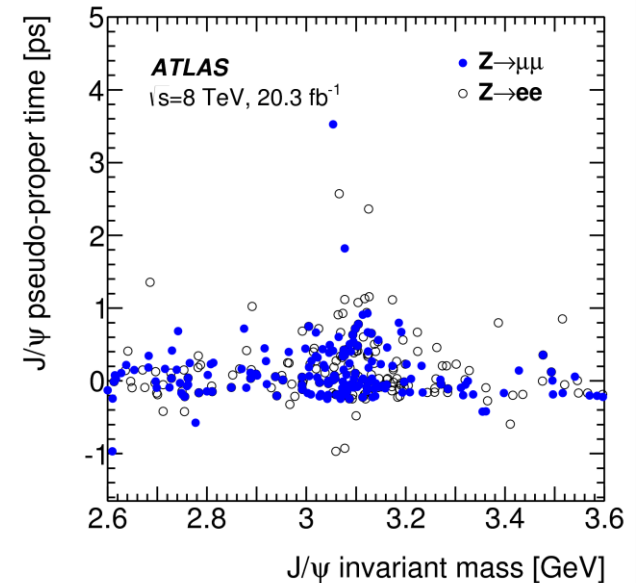






# $J/\psi + Z^0$ : masses and lifetimes

- ❑ Identify events with a Z boson (decaying into electrons or muons) AND another pair of muons around the  $J/\psi$  mass range
- ❑ 2D  $J/\psi$  mass and lifetime distribution fits used to assign sPlot weights to events with prompt and non-prompt  $J/\psi$  signal candidates and backgrounds
- ❑ Weighted Z candidates fitted with Z signal and multi-jet background templates





# J/ψ + Z<sup>0</sup> : event selection and yields

Some J/ψ are prompt, some are non-prompt

Unbinned ML fit in J/ψ mass and lifetime is used to extract prompt and non-prompt yields

Yields:      prompt                      non-prompt  
 $56 \pm 10 \pm 5$  (5σ)     $95 \pm 12 \pm 8$  (9σ)

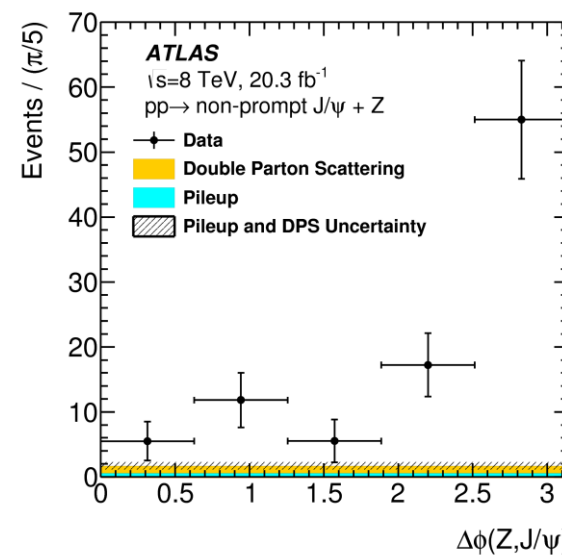
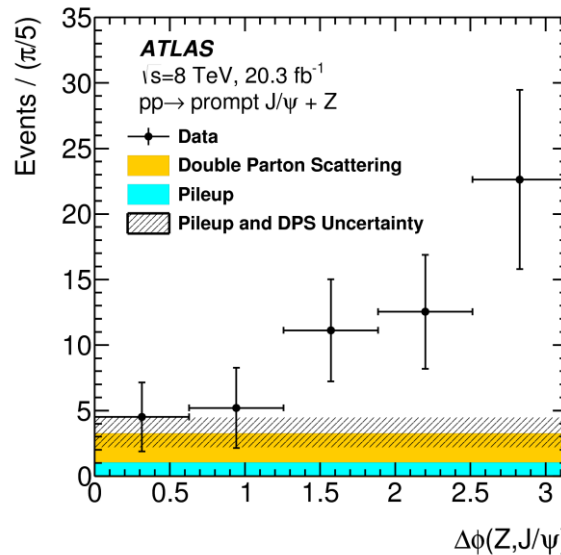
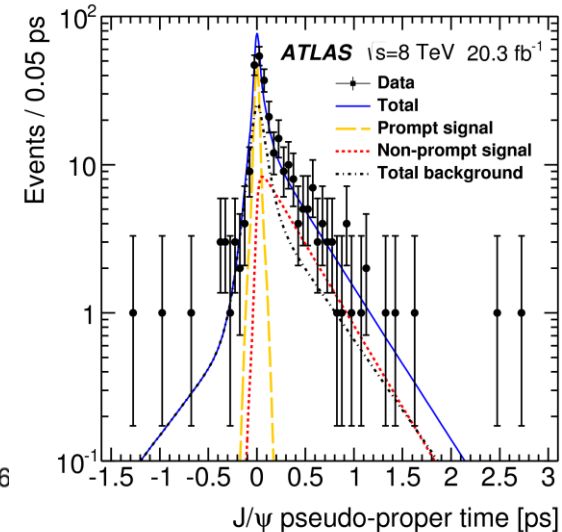
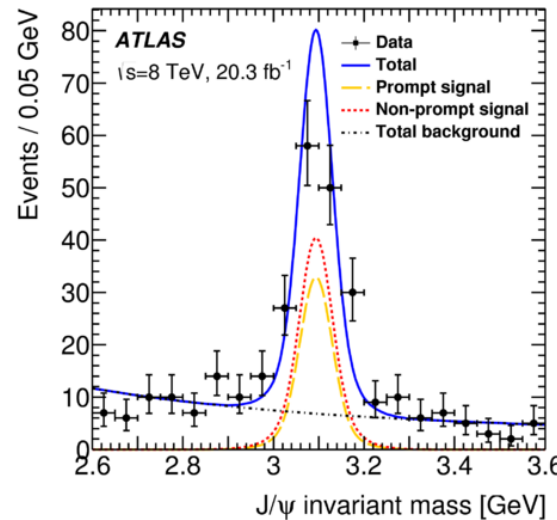
DPS:             $11.1^{+5.7}_{-5.0}$                        $5.8^{+2.8}_{-2.6}$

(assuming  $\sigma_{\text{eff}} = 15 \pm 3$  (stat)  $^{+5}_{-3}$  (syst) mb  
 arXiv:1301.6872

and  $\sigma_{J/\psi}$  from arXiv:1104.3038

If all signal in the first  $\Delta\phi$  bin is due to DPS, a lower limit is set:

$$\sigma_{\text{eff}} > 5.3 \text{ mb}$$

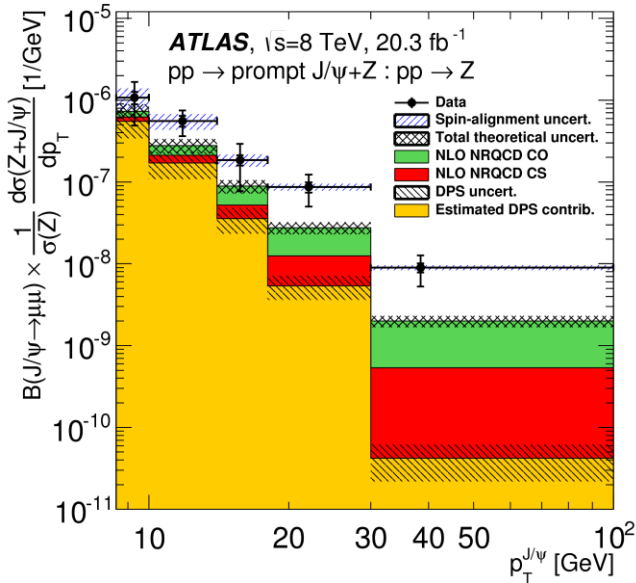




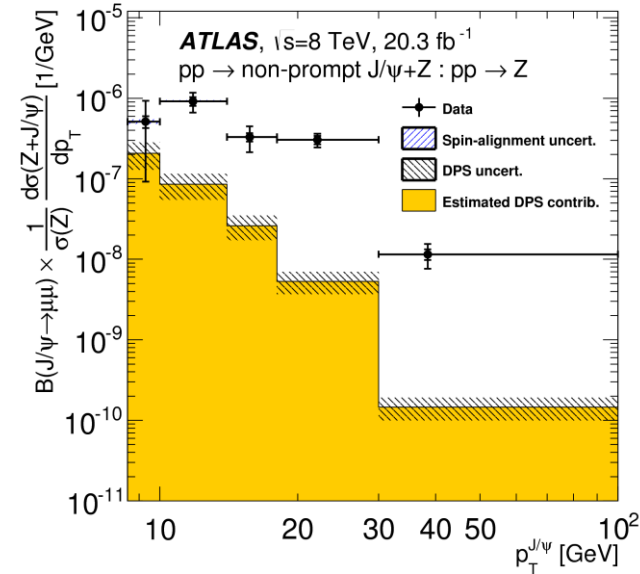
# J/ψ + Z<sup>0</sup> : cross sections

Ratios of the Z + J/ψ cross section to the inclusive Z cross section  
prompt non-prompt

fiducial	${}^p R_{Z+J/\psi}^{\text{fid}} = (36.8 \pm 6.7 \pm 2.5) \times 10^{-7}$	${}^{\text{np}} R_{Z+J/\psi}^{\text{fid}} = (65.8 \pm 9.2 \pm 4.2) \times 10^{-7}$
inclusive	${}^p R_{Z+J/\psi}^{\text{incl}} = (63 \pm 13 \pm 5 \pm 10) \times 10^{-7}$	${}^{\text{np}} R_{Z+J/\psi}^{\text{incl}} = (102 \pm 15 \pm 5 \pm 3) \times 10^{-7}$
DPS subtracted	${}^p R_{Z+J/\psi}^{\text{DPS sub}} = (45 \pm 13 \pm 6 \pm 10) \times 10^{-7}$	${}^{\text{np}} R_{Z+J/\psi}^{\text{DPS sub}} = (94 \pm 15 \pm 5 \pm 3) \times 10^{-7}$
DPS fraction	${}^p f_{\text{DPS}} = (29 \pm 9)\%$	${}^{\text{np}} f_{\text{DPS}} = (8 \pm 2)\%$



LO: Gong, arXiv:1210.2430  
NLO: Mao, arXiv:1102.0398



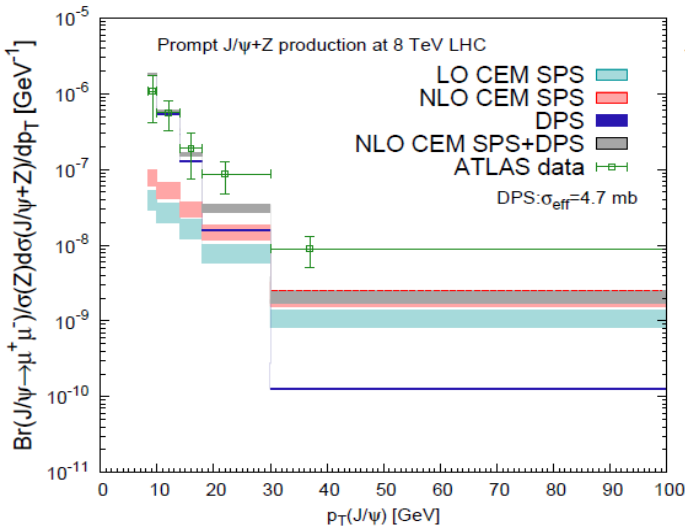
- ❑ A higher production rate predicted through CO than CS, CO dominant at high transverse momentum
- ❑ The expected production rate from the sum of CO and CS is lower than the data by a factor of 2 to 5
- ❑ Discrepancy increasing with transverse momentum



# J/ψ + Z<sup>0</sup> : recent theoretical developments

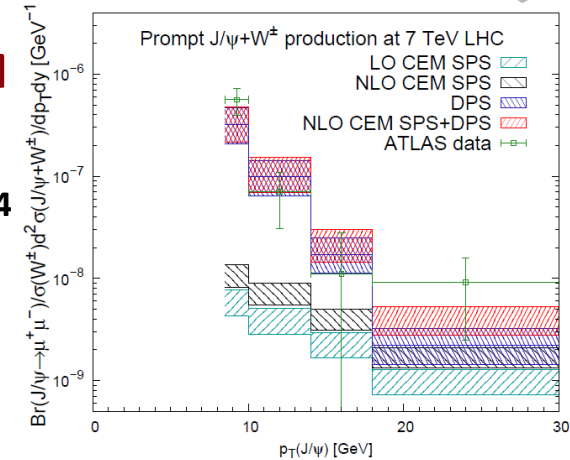
**Analysis of prompt J/ψ + W production, using improved CEM for SPS, may indicate that smaller  $\sigma_{\text{eff}}$  is needed for DPS**

Lansberg et al arXiv:1707:04



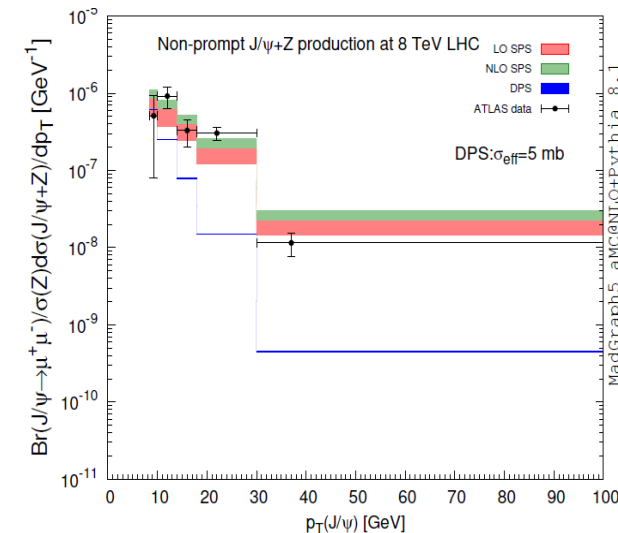
**Analysis of prompt J/ψ + Z production also 'needs' larger DPS (smaller  $\sigma_{\text{eff}}$ ), but even that may not be enough at high  $p_T$**

Lansberg et al arXiv:1608.03198



**However, non-prompt Z + J/ψ production seems to be saturated by SPS, leaving little room for large DPS contributions**

Lansberg et al arXiv:1611.09303





# J/ψ + J/ψ: yields and cross sections

- Unbinned ML fit to the two dimuon invariant masses to extract di- J/ψ signal
- Signal used to create prompt-prompt event weights from a 2D fit to the transverse decay length distributions of the two J/ψ

- Cross sections reported for two rapidity regions based on the sub-leading J/ψ rapidity

$$|y_{J/\psi_2}| < 1.05$$

$$N_{J/\psi J/\psi} = 3310 \pm 330$$

$$\sigma_{J/\psi J/\psi}^{\text{fid}} = 15.6 \pm 1.3 \pm 1.2 \pm 0.2 (\mathcal{B}) \pm 0.3(\mathcal{L}) \text{ pb}$$

$$1.05 < |y_{J/\psi_2}| < 2.1$$

$$N_{J/\psi J/\psi} = 3140 \pm 370$$

$$\sigma_{J/\psi J/\psi}^{\text{fid}} = 13.5 \pm 1.3 \pm 1.1 \pm 0.2 (\mathcal{B}) \pm 0.3(\mathcal{L}) \text{ pb}$$

- Correcting for muon acceptance and assuming unpolarised production

$$\sigma_{J/\psi J/\psi} = 82.2 \pm 8.3 \pm 6.3 \pm 0.9 (\mathcal{B}) \pm 1.6(\mathcal{L}) \text{ pb}$$

$$\sigma_{J/\psi J/\psi} = 78.3 \pm 9.2 \pm 6.6 \pm 0.9 (\mathcal{B}) \pm 1.5(\mathcal{L}) \text{ pb}$$

- the fraction of DPS events is determined by fitting DPS and SPS templates in  $\Delta y, \Delta \phi$  to the data, assign DPS and SPS event weights

$$f_{\text{DPS}} = (9.2 \pm 2.1 \pm 0.5)\%$$

$$\sigma_{J/\psi J/\psi}^{\text{DPS}} = 14.8 \pm 3.5 \pm 1.5 \pm 0.2 (\mathcal{B}) \pm 0.3(\mathcal{L}) \text{ pb}$$

- $\sigma_{\text{eff}}^{\text{eff}}$  measured from prompt di- J/ψ is lower than from other final states:

$$\sigma_{\text{eff}}^{J/\psi J/\psi} = 6.3 \pm 1.6(\text{stat}) \pm 1.0(\text{syst}) \pm 0.1(\text{BF}) \pm 0.1(\text{lumi}) \text{ mb}$$

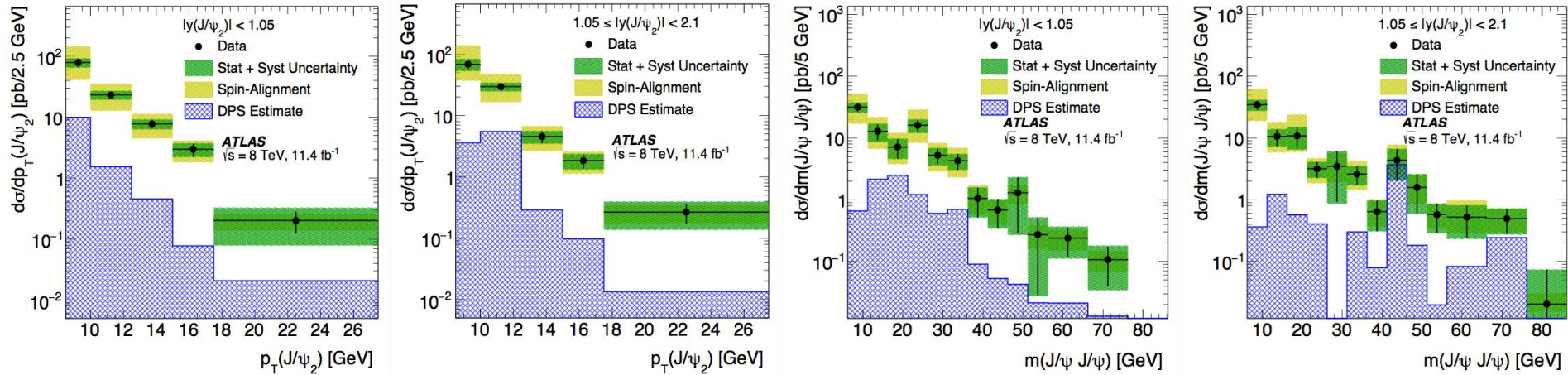




# $J/\psi + J/\psi$ : differential cross sections

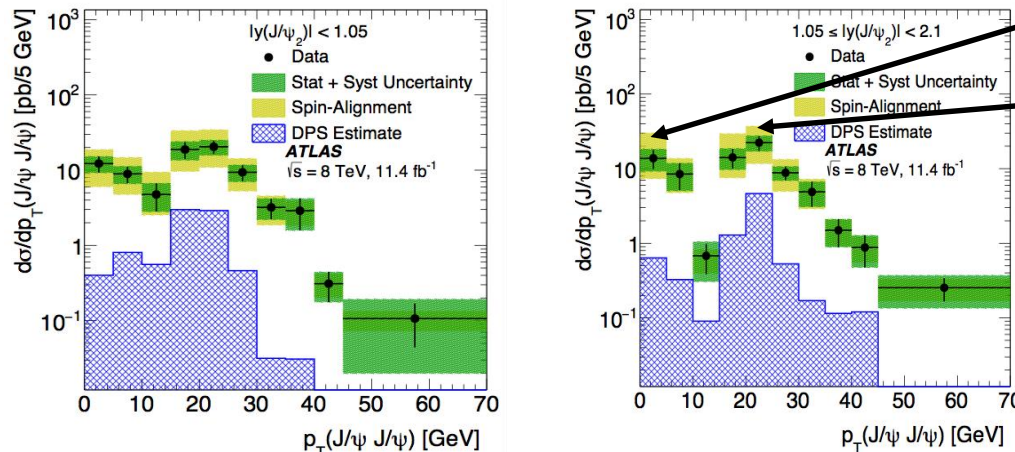
as a function of the sub-leading  $J/\psi$   $p_T$

as a function of the invariant mass



All for central (left) and forward (right) rapidity regions, with data-driven DPS estimates shown in blue

as a function of di- $J/\psi$   $p_T$



onia produced back-to-back

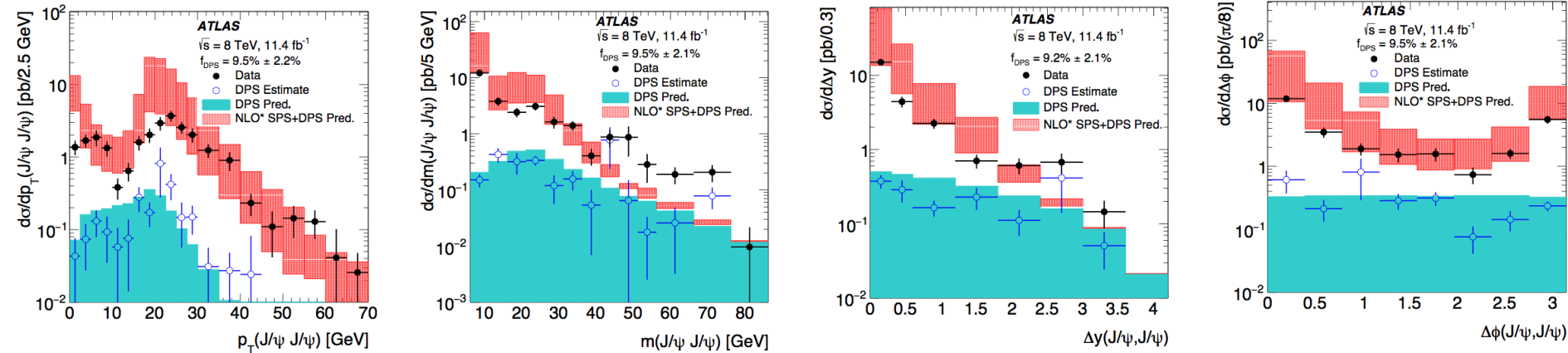
onia produced together,  
and back-to-back to a jet



# J/ψ + J/ψ: more differential distributions

(DPS+SPS) and DPS cross sections (full rapidity range) in the muon fiducial volume:

di- $J/\psi$   $p_T$  and invariant mass,  $\Delta y$  and  $\Delta\phi$



NLO\* SPS (with a feed-down correction factor)

Data points are compared to:

obtained using HELAC-Onia with matrix elements from  
 Lansberg, Shao arXiv:1410.8822, 1308.0474

LO DPS (normalised to measured)

Borschensky arXiv:1610.00666

Data largely in agreement with NLO\* SPS + LO DPS

Some localised disagreements for large invariant mass, large  $\Delta y$  and low  $p_T$

More realistic predictions for feed-down and a better treatment of parton transverse motion are needed



# Double onia production

Measurement of prompt  $J/\psi$  pair production in pp collisions at  $\sqrt{s} = 7$  TeV

JHEP 09 (2014) 094

arXiv:1406.0484

CMS

Measurement of the prompt  $J/\psi$  pair production cross-section in pp collisions at  $\sqrt{s} = 8$  TeV with the ATLAS detector

Eur. Phys. J. C77 (2017) 76

arXiv:1612.02950

ATLAS

Measurement of the  $J/\psi$  pair production cross-section in  $pp$  collisions at  $\sqrt{s} = 13$  TeV

Submitted to JHEP

arXiv:1612.07451

LHCb

Observation of  $Y(1S)$  pair production in proton-proton collisions at  $\sqrt{s} = 8$  TeV

Accepted by JHEP

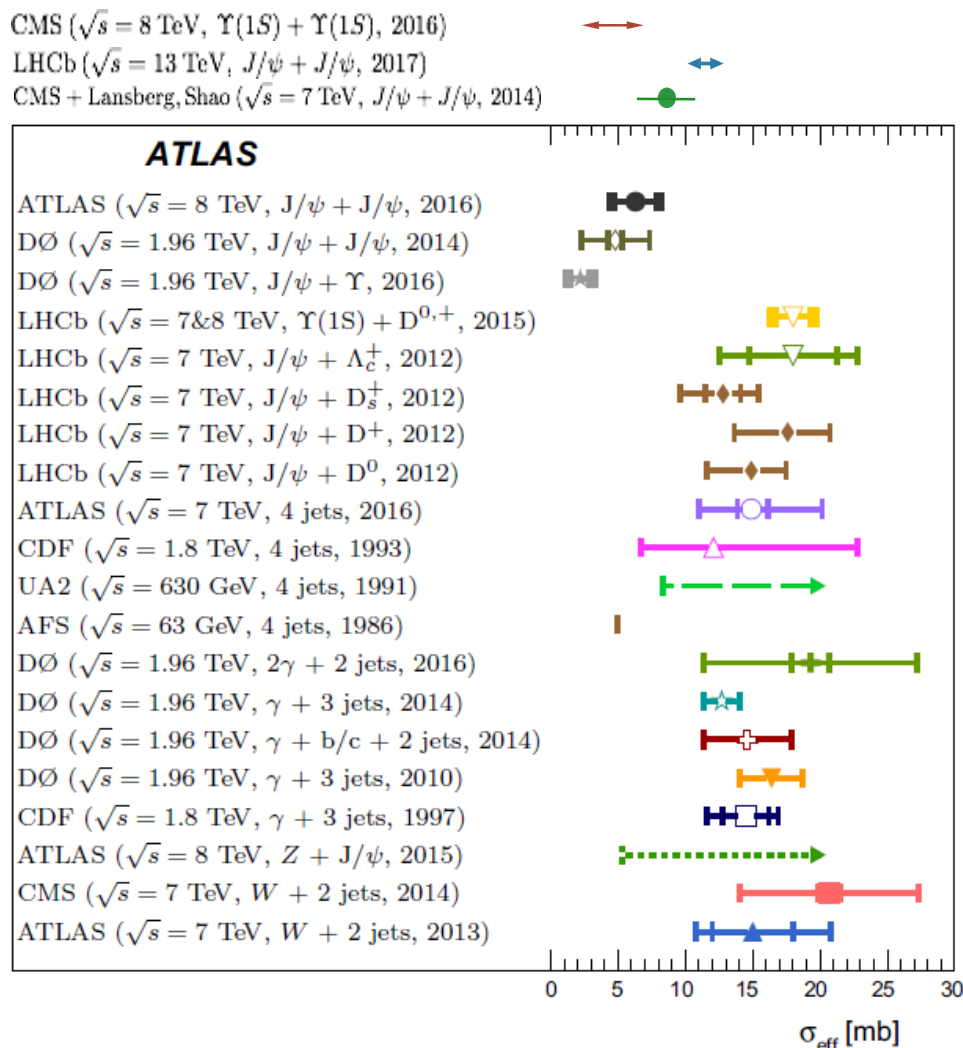
arXiv:1610.07095

CMS



# Summary

- Many results from the LHC experiments are now shedding light on double onia and associated onia production -- expect a lot more using 13 TeV data



- Some measured SPS contributions are well above theoretical predictions
- DPS contributions provide insight into the transverse profile of the proton, but our understanding is somewhat limited
- $\sigma_{\text{eff}}$  measured from prompt di-onia lower than from other final states
- Theoretical predictions of the dependence of  $\sigma_{\text{eff}}$  on the process and energy are needed
- There are some signs of improved understanding, but more work still to be done



**THANKS!**