

# **Five-quark Components in Baryons and Pentaquark States**

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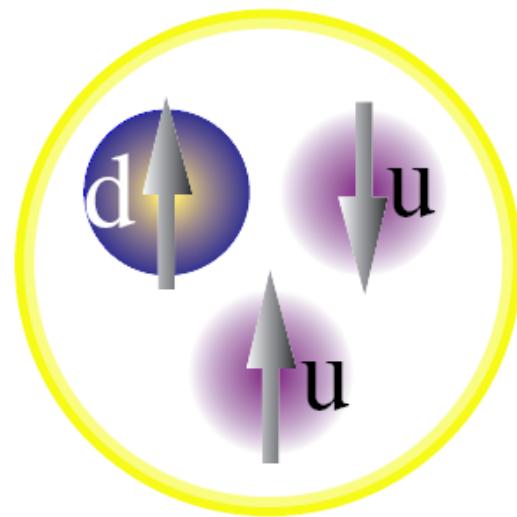
## **Outline:**

- **5-quark components in the proton**
- **Nature of  $1/2^-$  baryon nonet with strangeness**
- **Pentaquarks with hidden charm –  $P_c$  states**
- **Strange and beauty partners of  $P_c$  states**
- **Prospects**

# 1. Five-quark components in the proton

## Classical picture of the proton

Constituent Quarks



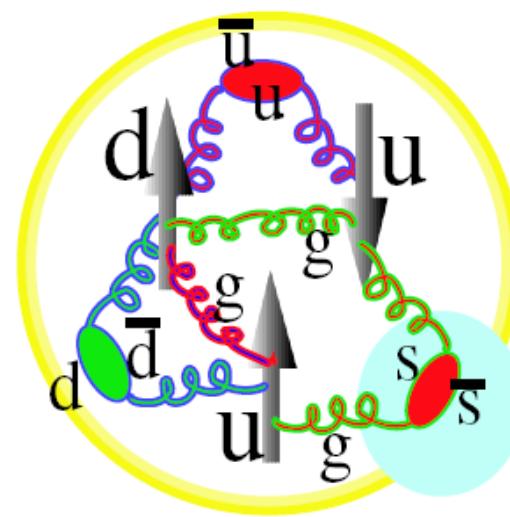
$$( Q^2 = 0 \text{ GeV}^2 )$$

baryon octet

masses, magn. momenta

1964–1974

Parton Distributions



$$( Q^2 > 1 \text{ GeV}^2 )$$

structure functions

momentum, spin

$$\bar{u}(x) = \bar{d}(x), \quad \bar{s}(x) = s(x)$$

1974–1992

# Flavor asymmetry of light quarks in the nucleon sea

Deep Inelastic Scattering (DIS) + Drell-Yan (DY) process



$$\bar{d} - \bar{u} \sim 0.12 \quad \text{for a proton}$$

Garvey&Peng, Prog. Part. Nucl. Phys. 47, 203 (2001)

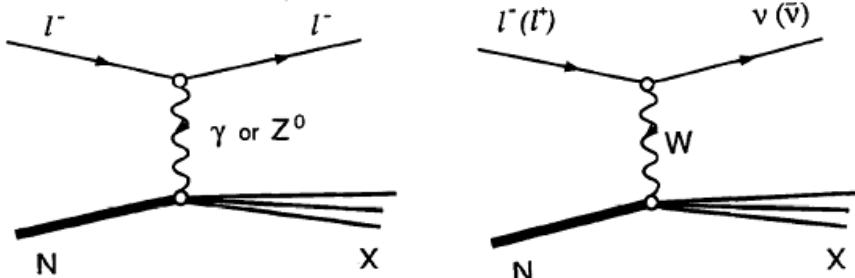
Table 1. Values of the integral  $\int_0^1 [\bar{d}(x) - \bar{u}(x)] dx$  determined from the DIS, semi-inclusive DIS, and Drell-Yan experiments.

Experiment	$\langle Q^2 \rangle$ (GeV $^2/c^2$ )	$\int_0^1 [\bar{d}(x) - \bar{u}(x)] dx$
NMC/DIS	4.0	$0.147 \pm 0.039$
HERMES/SIDIS	2.3	$0.16 \pm 0.03$
FNAL E866/DY	54.0	$0.118 \pm 0.012$

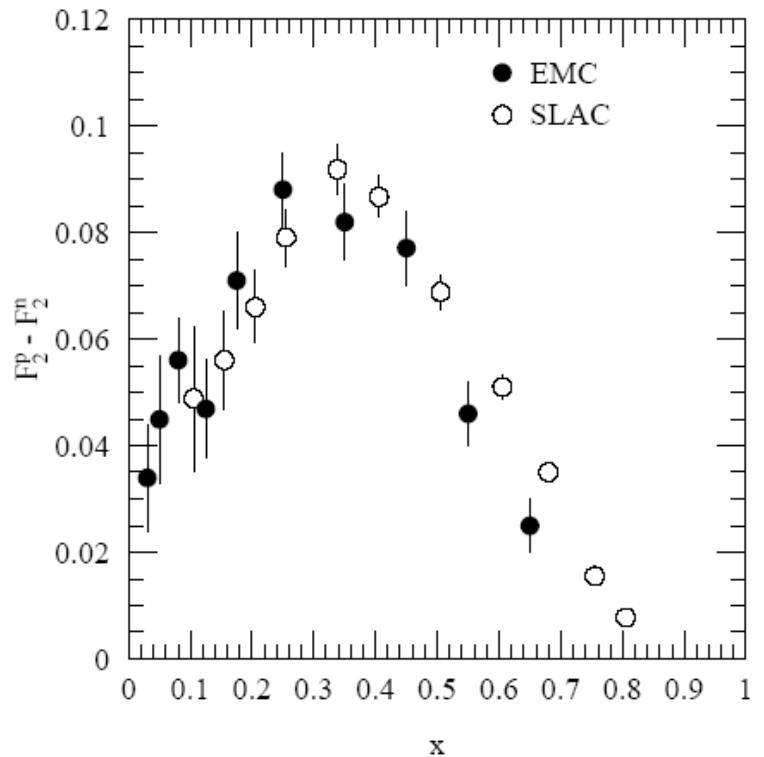
**DIS Gottfried Sum Rule :** assuming  $\bar{d} = \bar{u}$

$$I_2^p - I_2^n = \int_0^1 [F_2^p(x, Q^2) - F_2^n(x, Q^2)]/x \, dx = \sum_i [(Q_i^p)^2 - (Q_i^n)^2] = 1/3.$$

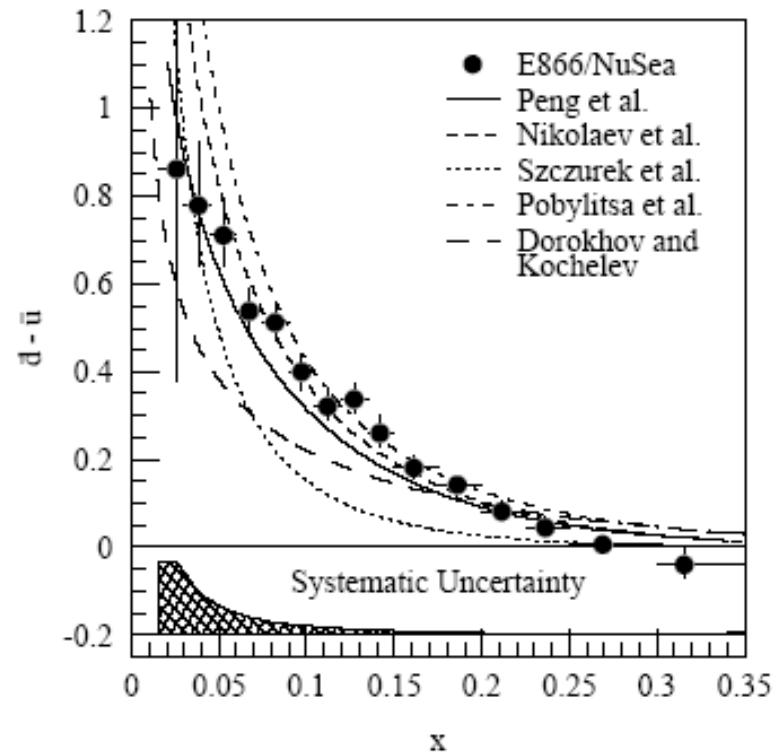
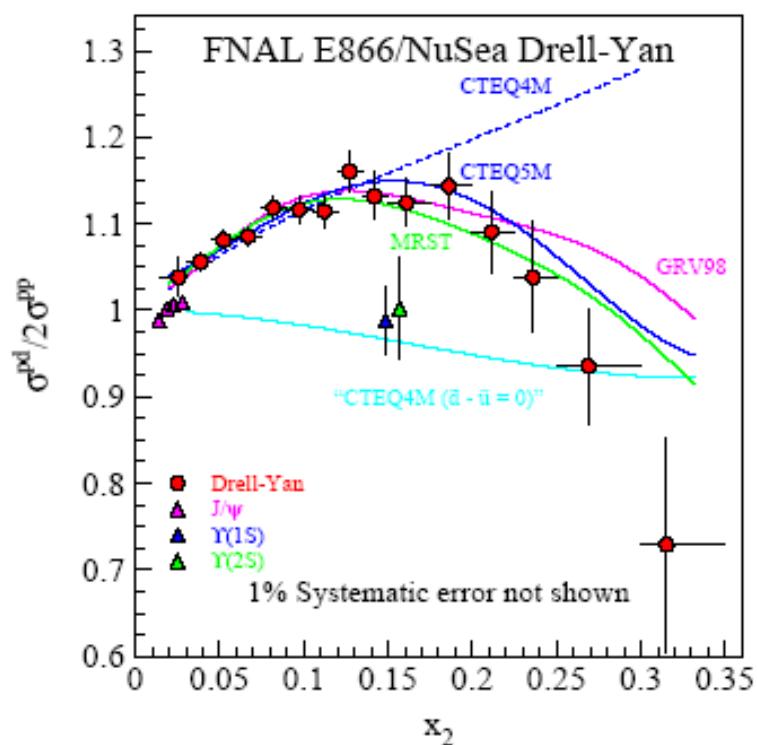
$$\int_0^1 [F_2^p(x, Q^2) - F_2^n(x, Q^2)]/x \, dx = \frac{1}{3} + \frac{2}{3} \int_0^1 [\bar{u}(x, Q^2) - \bar{d}(x, Q^2)] dx.$$



## Deep Inelastic Scattering (DIS)

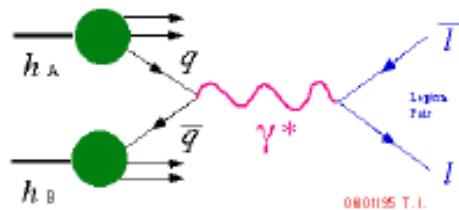


$$\sigma_{DY}(p+d)/2\sigma_{DY}(p+p) \simeq (1 + \bar{d}(x_2)/\bar{u}(x_2))/2.$$



**FIGURE 1.** Left panel: Cross section ratios of  $p+d$  over  $2(p+p)$  for Drell-Yan,  $J/\Psi$ , and  $\Upsilon$  production from FNAL E866. Right panel: Comparison of E866  $d - \bar{u}$  data with calculations from various models [2].

### The Drell-Yan Process



$$\frac{\bar{d}}{\bar{u}} > 1$$

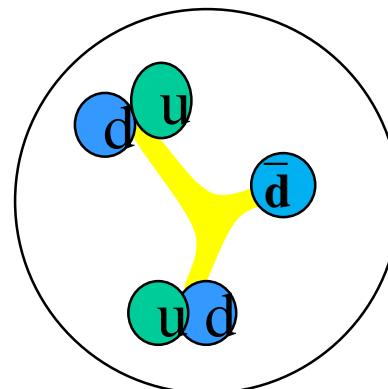
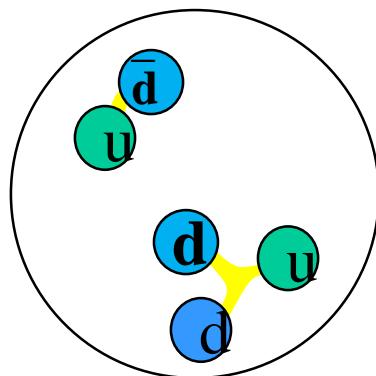
# Two major theoretical schemes for $\bar{d} - \bar{u} \sim 0.12$

**Meson cloud picture:** Thomas, Speth, Henley, Meissner, Miller, Weise, Oset, Brodsky, Ma, ...

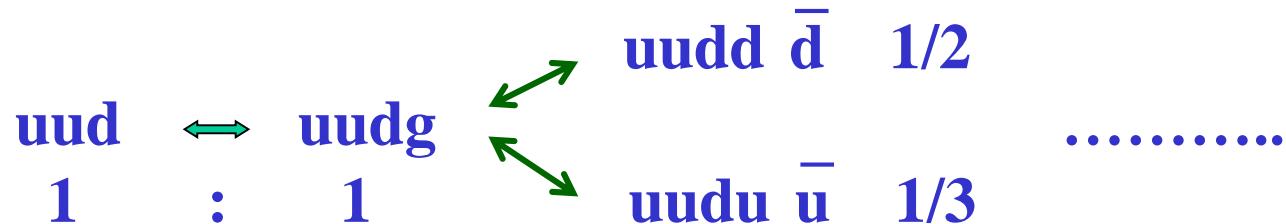
$$| p \rangle \sim | uud \rangle + \varepsilon_1 | n \text{ ( } udd \text{ ) } \pi^+ (\bar{d}u) \rangle + \varepsilon_2 | \Delta^{++} (\text{uuu}) \pi^- (\bar{u}d) \rangle + \varepsilon' | \Lambda \text{ (uds) } K^+ (\bar{s}u) \rangle \dots$$

**Penta-quark picture :** Riska, Zou, Zhu, ...

$$| p \rangle \sim | uud \rangle + \varepsilon_1 | [ud][ud] \bar{d} \rangle + \varepsilon' | [ud][us] \bar{s} \rangle + \dots$$



## Detailed balance model : Zhang, Ma, Zou, Yang, Alberg, Henley



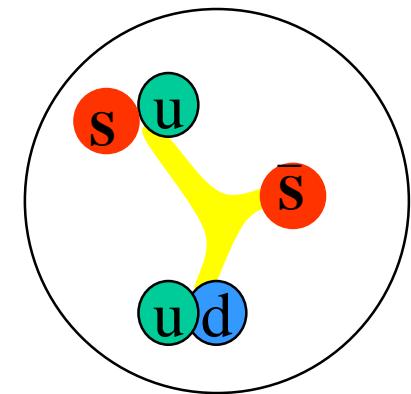
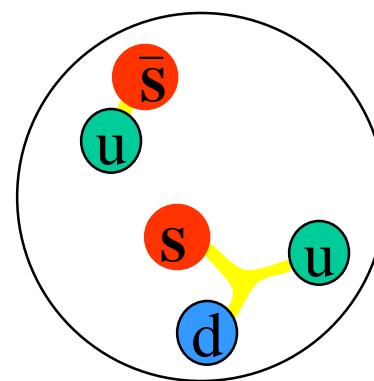
$$p = 0.168 (\text{uud}) + 0.168 (\text{uudg}) + 0.084 (\text{uudd } \bar{d}) + 0.056 (\text{uudu } \bar{u}) + 0.084 (\text{uudgg}) + \dots$$

$\bar{d} - \bar{u} \sim 0.124$

(uud+ng) 50% (uudd  $\bar{d}$ +ng) 22.4% (uudu  $\bar{u}$  +ng) 15.0%

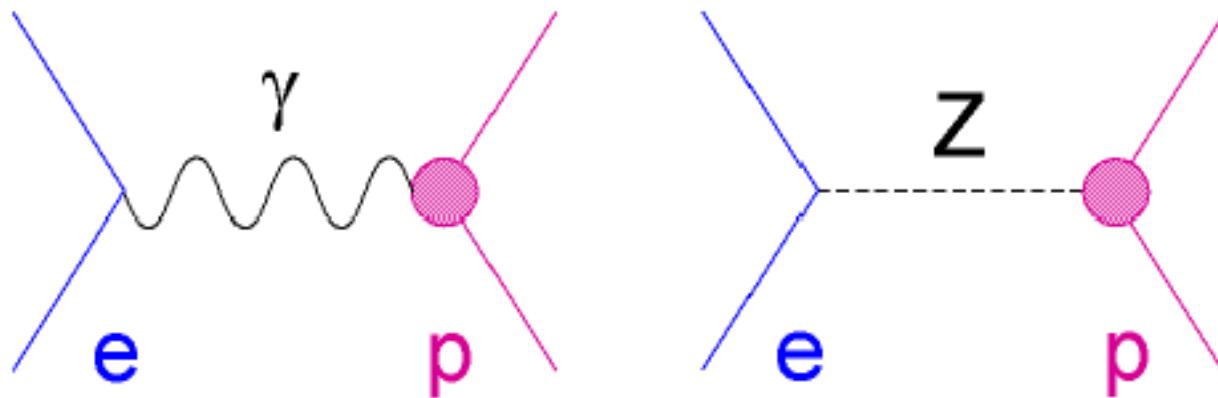
# Predictions for $\bar{s}/s$ asymmetry from two schemes :

	meson cloud	penta-quark
strange spin $\Delta s$ :	$< 0$	$< 0$
magnetic moment $\mu_s$ :	$< 0$	$> 0$
strange radius $r_s$ :	$< 0$	$> 0$



Expt:  $\Delta s = -0.05 \sim -0.1$  D. de Florian et al., PRD71 (2005) 094018

# The strange magnetic moment $\mu_s$ and radii $r_s$ from parity violating electron scattering



G0,HAPPEX/CEBAF, SAMPLE/MIT-Bates, A4/MAMI

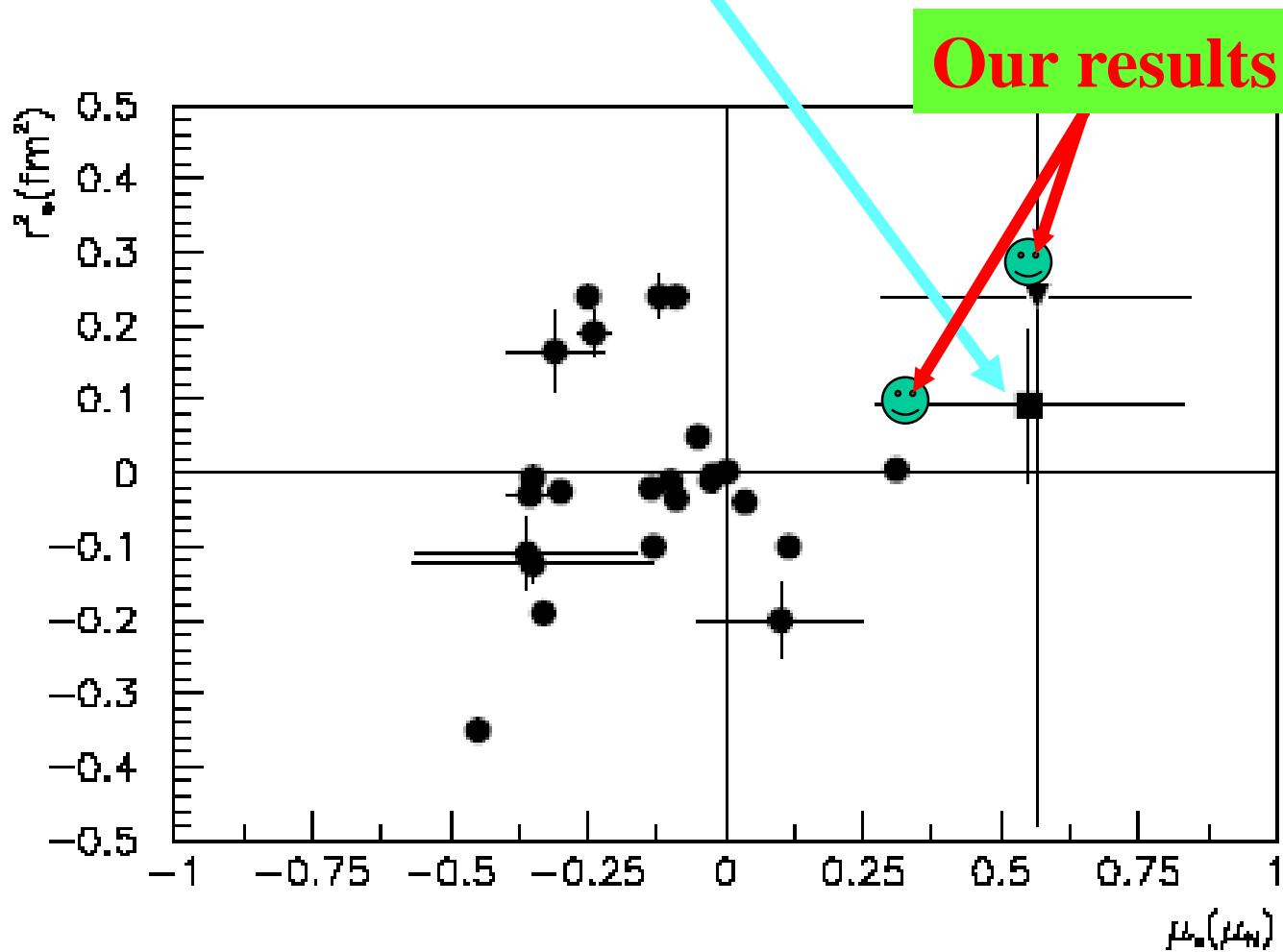
HAPPEX/CEBAF, Phys.Rev.Lett. 96 (2006) 022003

G0/CEBAF, Phys.Rev.Lett. 95 (2005) 092001

A4/MAMI, Phys.Rev.Lett. 94 (2005) 152001

SAMPLE/MIT-Bates: Phys.Lett.B583 (2004) 79

# Theory vs experiment for $\mu_s$ and $r_s$



Zou&Riska, PRL95(2005)072001; Riska&Zou, PLB636 (2006) 265  
An-Riska-Zou, PRC73 (2006) 035207

# Experiment extraction of $\mu_s$ and $r_s$ wrong?

R.Young et al., PRL97 (2006) 102002  $\rightarrow \mu_s \sim 0$

S.Baunack et al.(A4), PRL102(2009)151803

With  $\sim 25\%$   $\bar{q}qqqq$  components in the proton, the “spin crisis” and single spin asymmetry may also be naturally explained.

An-Riska-Zou, PRC73 (2006) 035207; F.X.Wei, B.S.Zou, hep-ph/0807.2324

$$\Delta_u = 0.85 \pm 0.17$$

$$\Delta_u = \frac{4}{3} |A_{3q}|^2$$

$$\Delta_d = -(0.33 \sim 0.56)$$

$$\Delta_d = -\frac{1}{3}(1 - P_{s\bar{s}})$$

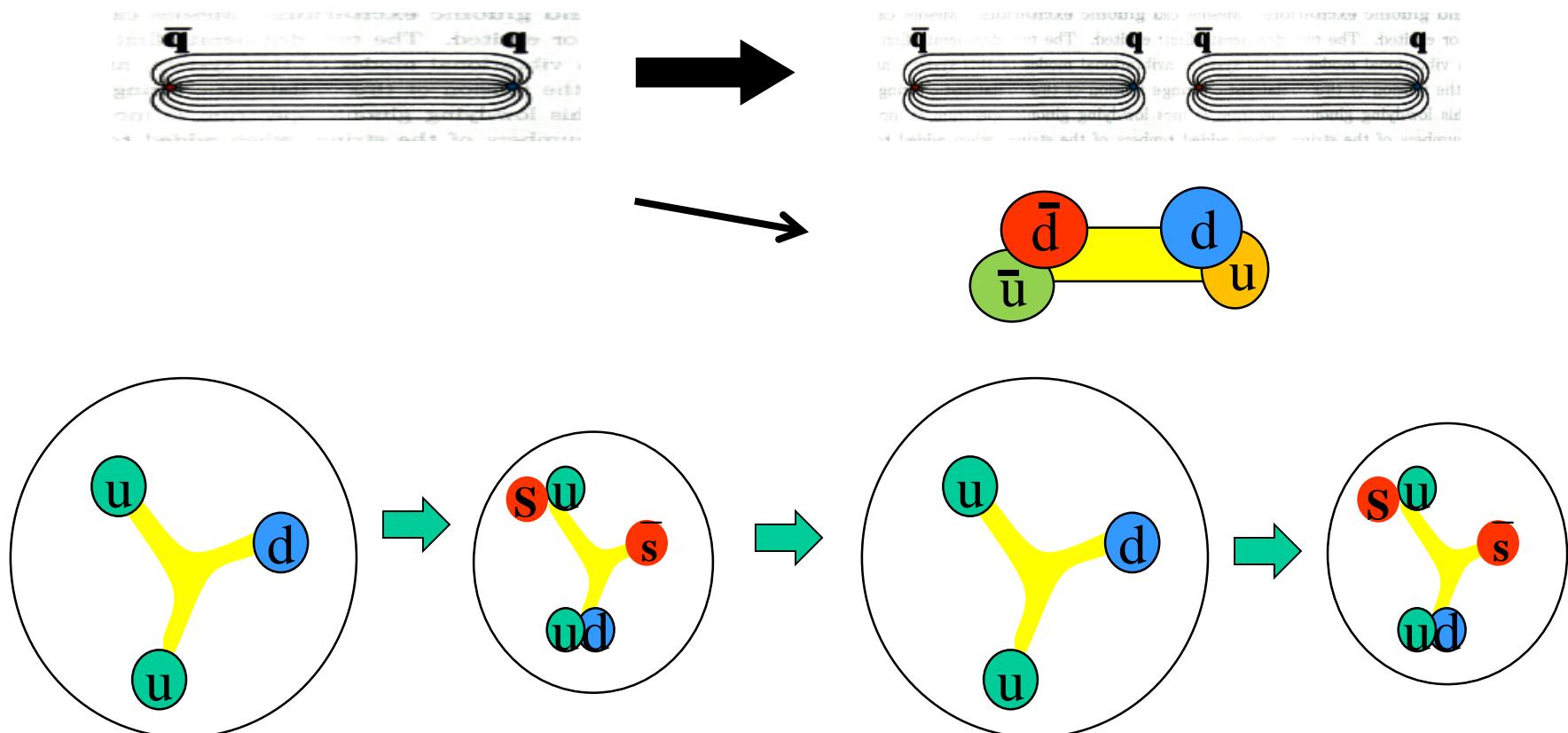
$$\Delta L_q = \frac{4}{3}(P_{d\bar{d}} + P_{s\bar{s}})$$

We must go beyond the simple 3q models,  
meson cloud vs penta-quark not settled yet.

# Quenched & unquenched quark models

Unquenching dynamics: gluons  $\rightarrow \bar{q}q$

crucial for quark confinement & hadron structure



quenched or unquenched quark models give very different predictions of baryon spectrum

## 2. Nature of $1/2^-$ baryon nonet with strangeness

- Mass pattern : quenched or unquenched ?

$$\text{uds (L=1) } 1/2^- \sim \Lambda^*(1670) \sim [\text{us}][\text{ds}] \bar{s}$$

$$\text{uud (L=1) } 1/2^- \sim N^*(1535) \sim [\text{ud}][\text{us}] \bar{s}$$

$$\text{uds (L=1) } 1/2^- \sim \Lambda^*(1405) \sim [\text{ud}][\text{su}] \bar{u}$$

$$\text{uus (L=1) } 1/2^- \sim \Sigma^*(1390) \sim [\text{us}][\text{ud}] \bar{d}$$

Zou et al, NPA835 (2010) 199 ; CLAS, PRC87(2013)035206

- Strange decays of  $N^*(1535)$  : **PDG  $\rightarrow$  large  $g_{N^*N\eta}$**

$$J/\psi \rightarrow \bar{p} N^* \rightarrow \bar{p} (K\Lambda) / \bar{p} (\rho\eta) \rightarrow \text{large } g_{N^*K\Lambda}$$

Liu&Zou, PRL96 (2006) 042002; Geng,Oset,Zou&Doring, PRC79 (2009) 025203

$$\gamma p \rightarrow p\eta' \& pp \rightarrow pp\eta' \rightarrow \text{large } g_{N^*N\eta'}$$

M.Dugger et al., PRL96 (2006) 062001; Cao&Lee, PRC78(2008) 035207

$$\pi^- p \rightarrow n\phi \& pp \rightarrow pp\phi \& pn \rightarrow d\phi \rightarrow \text{large } g_{N^*N\phi}$$

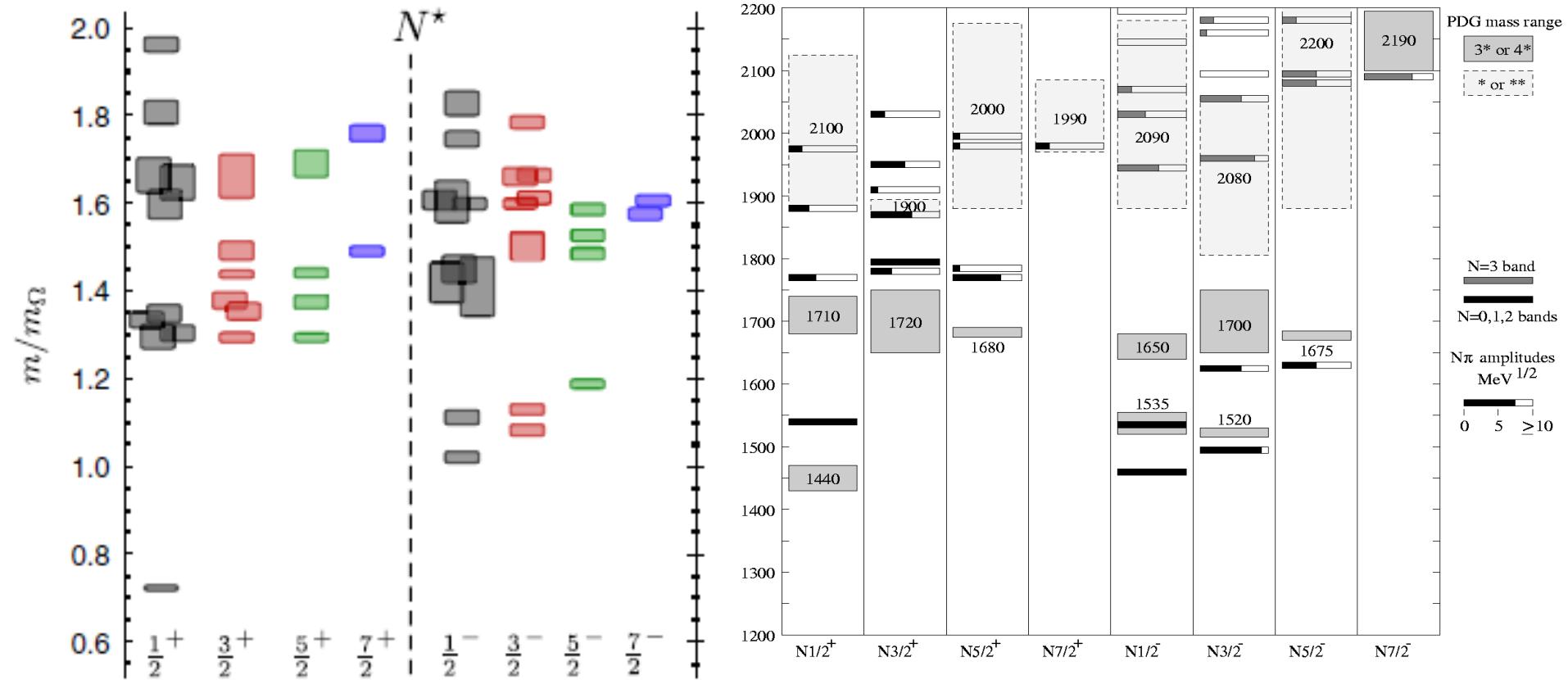
Xie, Zou & Chiang, PRC77(2008)015206; Cao, Xie, Zou & Xu, PRC80(2009)025203

- Strange decays of  $\Lambda^*(1670)$  : **PDG  $\rightarrow$  large  $g_{\Lambda^*\Lambda\eta}$**

narrower width (35MeV) than  $\Lambda^*(1405)$

# Problems for quenched $q^3$ picture for baryon resonances:

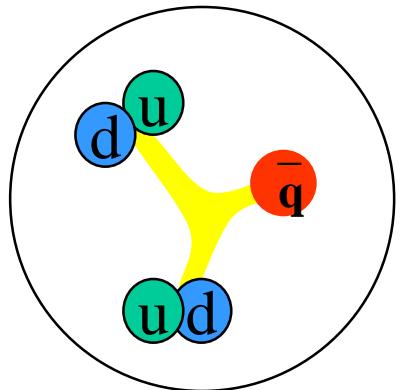
1) “Missing resonances” ; 2) mass ordering for the lowest ones



Lattice, R.Edwards et al,  
PRD84(2011)074508

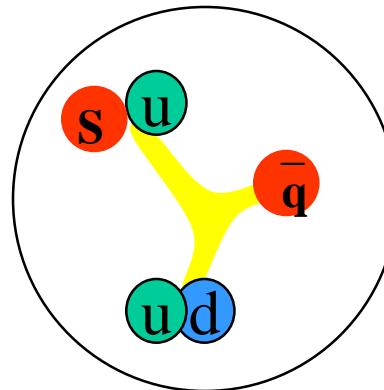
$q^3$  model, Capstick&Roberts,  
PPNP45(2000)S241

# Why $N^*(1535) \frac{1}{2}^-$ is heavier than $N^*(1440) \frac{1}{2}^+$ ?



$\bar{q}$   
 $[ud]$   
 $[ud]$

$\}^{1/2+}_{L=1}$



$\bar{q}$   
 $[ud]$   
 $[us]$

$\}^{1/2 -}_{L=0}$

$$N^*(1535) \sim uud (L=1) + \varepsilon [ud][us] \bar{s} + \dots$$

$$N^*(1440) \sim uud (n=1) + \xi [ud][ud] \bar{d} + \dots$$

$$\Lambda^*(1405) \sim uds (L=1) + \varepsilon [ud][su] \bar{u} + \dots$$

**$N^*(1535)$ :  $[ud][us] \bar{s} \rightarrow$  larger coupling to  $N\eta$ ,  $N\eta'$ ,  $N\phi$  &  $K\Lambda$ , weaker to  $N\pi$  &  $K\Sigma$ , and heavier !**

# quench vs un-quench for mesons

$\bar{q}q$   $^3S_1$  nonet

$\phi(1020)$      $\bar{s}s$

$K(892)$      $\bar{s}d$

$\omega(782)$      $\bar{u}u + \bar{d}d$   
 $\rho(770)$      $\bar{u}u - \bar{d}d$

$\bar{q}q$   $^3P_0$    or    $\bar{q}^2q^2$  nonet ?

$a_0(980)$      $\bar{u}u - \bar{d}d$ ,     $[\bar{u}\bar{s}][us] - [\bar{d}\bar{s}][ds]$

$f_0(980)$      $\bar{s}s$ ,     $[\bar{u}\bar{s}][us] + [\bar{d}\bar{s}][ds]$

$\kappa(800)$      $\bar{s}d$ ,     $[\bar{s}\bar{u}][ud]$

$f_0(600)$      $\bar{u}u + \bar{d}d$ ,     $[\bar{u}\bar{d}][ud]$

$D_{s0}^*(2317) \sim \bar{s}c$  ( $L=1$ ) +  $[\bar{q}\bar{s}][qc]$  + DK + ...

$D_{s1}^*(2460) \sim \bar{s}c$  ( $L=1$ ) +  $D^*K$  + ...

$X(3872) \sim \bar{c}c$  ( $L=1$ ) +  $[\bar{q}\bar{c}][qc]$  +  $D^*D$  + ...

# Alternative pictures :

## Hadronic molecules

$$N^*(1440) \sim N\sigma$$

$$N^*(1535) \sim K\Sigma - K\Lambda$$

$$\Lambda^*(1405) \sim KN - \Sigma\pi$$

## Penta-quark states

$$N^*(1440) \sim [ud][ud] \bar{q}$$

$$N^*(1535) \sim [ud][us] \bar{s}$$

$$\Lambda^*(1405) \sim [ud][sq] \bar{q}$$

Kaiser, Weise, Oset, Ramos,  
Oller, Meissner, Hyodo, Jido,  
Hosaka, ...

$\Sigma^* \sim 1430$  MeV      Oller, Meissner, Oh, Khemchandani, ...

$\Xi^* \sim 1610$  MeV      Oh, Ramos, Oset, ...

# Important implications:

- $\bar{q}qqqq$  in S-state more favorable than  $\underline{q}\underline{q}\underline{q}$  with  $L=1$  !  
&  $\bar{q}qqq$  in S-state more favorable than  $\underline{q}\underline{q}$  with  $L=1$  !

$1/2^-$  baryon nonet  $\sim \bar{q}q^2q^2$  state + ...

$0^+$  meson octet  $\sim \bar{q}^2q^2$  state + ...

Draging out  $\bar{q}q$  from gluon field –  
an important excitation mechanism for hadrons !  
multiquark components are important for hadrons !

## Totally different predictions for the lowest $\Xi^*$ & $\Omega^*$ :

$\bar{q}q^4$  ( $L=0$ )

$q^3$  ( $L=1$ )

$\Xi^* (1/2^-)$  [us][ds]  $\bar{d}$  ~ 1540 MeV

uss ~ 1800 MeV

$\Omega^*(1/2^-)$  [us] ss  $\bar{u}$  ~ 1840 MeV

sss ~ 2000 MeV

Yuan-An-Wei-Zou-Xu, PRC87(2013)025205 Capstick-Isgur, PRD34(1986)2809

$\Omega^*(3/2^-)$  as sss -  $\bar{u}u$ sss mixture : ~ 1780 MeV

by instanton/NJL interaction

An-Metsch-Zou, PRC87(2013) 065207; An-Zou, PRC89 (2014) 055209

**$\Sigma^*$  in PDG:** \*\*\*\*  $\Sigma(1189)1/2^+$   $\Sigma^*(1385)3/2^+$   $\Sigma^*(1670)3/2^-$   
 $\Sigma^*(1775)5/2^-$   $\Sigma^*(1915)5/2^+$   $\Sigma^*(2030)7/2^+$   
 \*\*\*  $\Sigma^*(1660)1/2^+$   $\Sigma^*(1750)1/2^-$   $\Sigma^*(1940)3/2^-$   $\Sigma^*(2250)??$   
 \*\*  $\Sigma^*(1620)1/2^-$   $\Sigma^*(1690)??$   $\Sigma^*(1880)1/2^+$   
 $\Sigma^*(2080)3/2^+$   $\Sigma^*(2455)??$   $\Sigma^*(2620)??$   
 \*  $\Sigma^*(1480)??$   $\Sigma^*(1560)??$   $\Sigma^*(1580)3/2^-$   $\Sigma^*(1770)1/2^+$   
 $\Sigma^*(1840)3/2^+$   $\Sigma^*(2000)3/2^-$   $\Sigma^*(2070)5/2^+$   $\Sigma^*(2100)7/2^-$   
 $\Sigma^*(3000)??$   $\Sigma^*(3170)??$

**$\Xi^*$  in PDG:** \*\*\*\*  $\Xi(1320)1/2^+$ ,  $\Xi(1530)3/2^+$   
 \*\*\*  $\Xi(1690)$ ,  $\Xi(1820)3/2^-$ ,  $\Xi(1950)$ ,  $\Xi(2030)$   
 \*\*  $\Xi(2250)$ ,  $\Xi(2370)$   
 \*  $\Xi(1620)$ ,  $\Xi(2120)$ ,  $\Xi(2500)$

**$\Omega^*$  in PDG:**  $\Omega(1672)3/2^+****$ ,  $\Omega(2250)***$ ,  $\Omega(2380)**$ ,  $\Omega(2470)**$

**Experiment knowledge on hyperon states still very poor !**

# Difficulties to pin down pentaquarks

Fate of the first pentaquark predicted and observed:

- 1959:  $\bar{K}N$  molecule predicted by Dalitz-Tuan, PRL2, 425
- 1961:  $\Lambda(1405) \rightarrow \Sigma\pi$  observed by Alston et al., PRL6, 698
- 1964: Quark model (uds) for  $\Lambda(1405)$
- 1995:  $\bar{K}N$  dynamically generated -- Kaiser et al., NPA954, 325
- 2001: 2 pole structure by  $\bar{K}N-\Sigma\pi$  -- Oller et al., PLB500, 263

PDG2010: “The clean  $\Lambda_c$  spectrum has in fact been taken to settle the decades-long discussion about the nature of the  $\Lambda(1405)$  —true 3-quark state or mere  $\bar{K}N$  threshold effect?— unambiguously in favor of the first interpretation.”

## Fate of the last famous fading pentaquark $\Theta^+(1540)$ :

- 1997:  $Z^+(1530)$  predicted by Diakonov et al., ZPA359, 305
- 2003:  $\Theta^+(1540) \rightarrow K^+ n$  claimed by LEPS, PRL91, 012002
- 2003:  $\bar{s}(ud)(ud)$  for  $\Theta(1540)$  by Jaffe&Wilczek, PRL91, 232003
- 2003:  $\bar{s}ud(ud)$  for  $\Theta(1540)$  by Karliner&Lipkin, PLB575, 249
- 2004: supported by 10 expts  $\rightarrow \Theta(1540)$  well-established by PDG
- 2004: not supported by BESII, PRD70, 012004
- 2005: not supported by many high stats experiments
- 2006: removed from PDG

Note:  $\Theta^+(1540)$  is not supported by hadronic molecule model & chiral quark model by Huang, Zhang, Yu, Zou, PLB586(2004)69

### 3. Pentaquarks with hidden charm – $P_c$ states

New direction for hunting for pentaquarks:

Extension to hidden charm and beauty for baryons

$N^*(1535)$   $\bar{s}suud$

$N^*(4260)$   $\bar{c}cuud$  J.J.Wu, R.Molina, E.Oset, B.S.Zou.  
Phys.Rev.Lett. 105 (2010) 232001

$N^*(11050)$   $\bar{b}buud$  J.J.Wu, L.Zhao, B.S.Zou. PLB709(2012)70

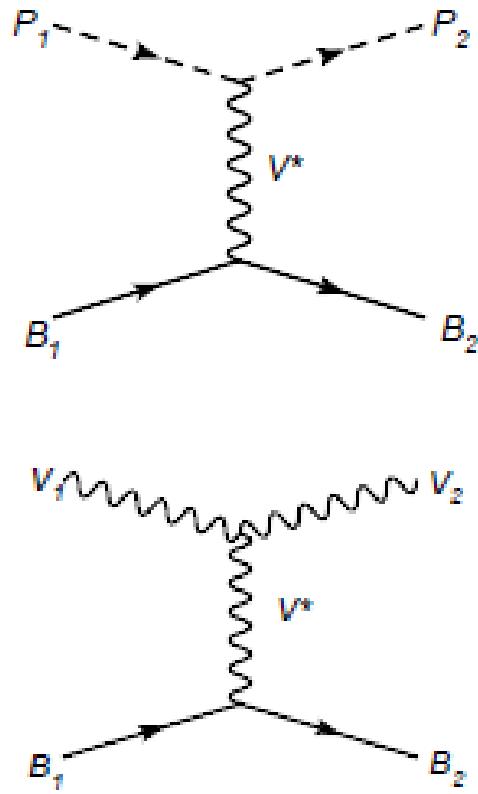
$\Lambda^*(1405)$   $\bar{q}quds$

$\Lambda^*(4210)$   $\bar{c}cuuds$  J.J.Wu, R.Molina, E.Oset, B.S.Zou.  
Phys.Rev.Lett. 105 (2010) 232001

$\Lambda^*(11020)$   $\bar{b}bbuds$  J.J.Wu, L.Zhao, B.S.Zou. PLB709(2012)70

# From $K\Sigma$ , $\bar{K}N \rightarrow \bar{D}\Sigma_c$ , $\bar{D}_s\Lambda_c \rightarrow B\Sigma_b$ , $B_s\Lambda_b$ bound states

**“Prediction of narrow  $N^*$  and  $\Lambda^*$  resonances with hidden charm above 4 GeV”,  
Wu, Molina, Oset, Zou, PRL105 (2010) 232001**



$$\mathcal{L}_{VVV} = ig \langle V^\mu [V^\nu, \partial_\mu V_\nu] \rangle$$

$$\mathcal{L}_{PPV} = -ig \langle V^\mu [P, \partial_\mu P] \rangle$$

$$\mathcal{L}_{BBV} = g(\langle \bar{B} \gamma_\mu [V^\mu, B] \rangle + \langle \bar{B} \gamma_\mu B \rangle \langle V^\mu \rangle)$$

$$V_{ab(P_1 B_1 \rightarrow P_2 B_2)} = \frac{C_{ab}}{4f^2} (E_{P_1} + E_{P_2}),$$

$$V_{ab(V_1 B_1 \rightarrow V_2 B_2)} = \frac{C_{ab}}{4f^2} (E_{V_1} + E_{V_2}) \vec{\epsilon}_1 \cdot \vec{\epsilon}_2,$$

$$T = [1 - VG]^{-1}V$$

$$T_{ab} = \frac{g_a g_b}{\sqrt{s} - z_R}$$

	$(I, S)$	$M$	$\Gamma$	$\Gamma_i$				$J^P$
$N^* - \bar{D}\Sigma_c$	$(1/2, 0)$			$\pi N$	$\eta N$	$\eta' N$	$K\Sigma$	$\eta_c N$
		4261	56.9	3.8	8.1	3.9	17.0	23.4
$\Lambda^*$	$\bar{K}N$				$\pi\Sigma$	$\eta\Lambda$	$\eta'\Lambda$	$K\Xi$
	$(0, -1)$	4209	32.4	15.8	2.9	3.2	1.7	2.4
		4394	43.3	0	10.6	7.1	3.3	5.8
								16.3

TABLE V: Mass ( $M$ ), total width ( $\Gamma$ ), and the partial decay width ( $\Gamma_i$ ) for the states from  $PB \rightarrow PB$ , with units in MeV.

	$(I, S)$	$M$	$\Gamma$	$\Gamma_i$				
$N^* - \bar{D}^*\Sigma_c$	$(1/2, 0)$	4412	47.3	$\rho N$	$\omega N$	$K^*\Sigma$		$J/\psi N$
				3.2	10.4	13.7		19.2
$\Lambda^*$	$K^*N$				$\rho\Sigma$	$\omega\Lambda$	$\phi\Lambda$	$K^*\Xi$
	$(0, -1)$	4368	28.0	13.9	3.1	0.3	4.0	1.8
		4544	36.6	0	8.8	9.1	0	5.0
								13.8

TABLE VI: Mass ( $M$ ), total width ( $\Gamma$ ), and the partial decay width ( $\Gamma_i$ ) for the states from  $VB \rightarrow VB$  with units in MeV.

# Further studies support such hidden charm N\*

W.L.Wang, F.Huang, Z.Y.Zhang, B.S.Zou, PRC84(2011)015203:

**Chiral quark model** →  $\bar{D}\Sigma_c$  state  $\sim 4.3$  GeV

J.J.Wu, T.S.H.Lee, B.S.Zou, PRC85(2012)044002:

**EBAC-DCC model** →  $\bar{D}\Sigma_c (1/2^-) \sim 4.3$  GeV,  
 $\bar{D}^*\Sigma_c (1/2^-, 3/2^-) \sim 4.4 - 4.5$  GeV -

C.W.Xiao, J.Nieves, E.Oset, PRD 88 (2013) 056012:

**Heavy quark spin symmetry** → 7 such N\* molecules

$\bar{D}\Sigma_c (1/2^-) \sim 4.26$  GeV,  $\bar{D}\Sigma_c^* (3/2^-) \sim 4.33$  GeV,  
 $\bar{D}^*\Sigma_c (1/2^-, 3/2^-) \sim 4.41, 4.42$  GeV,  
 $\bar{D}^*\Sigma_c^* (1/2^-, 3/2^-, 5/2^-) \sim 4.48 - 4.49$  GeV

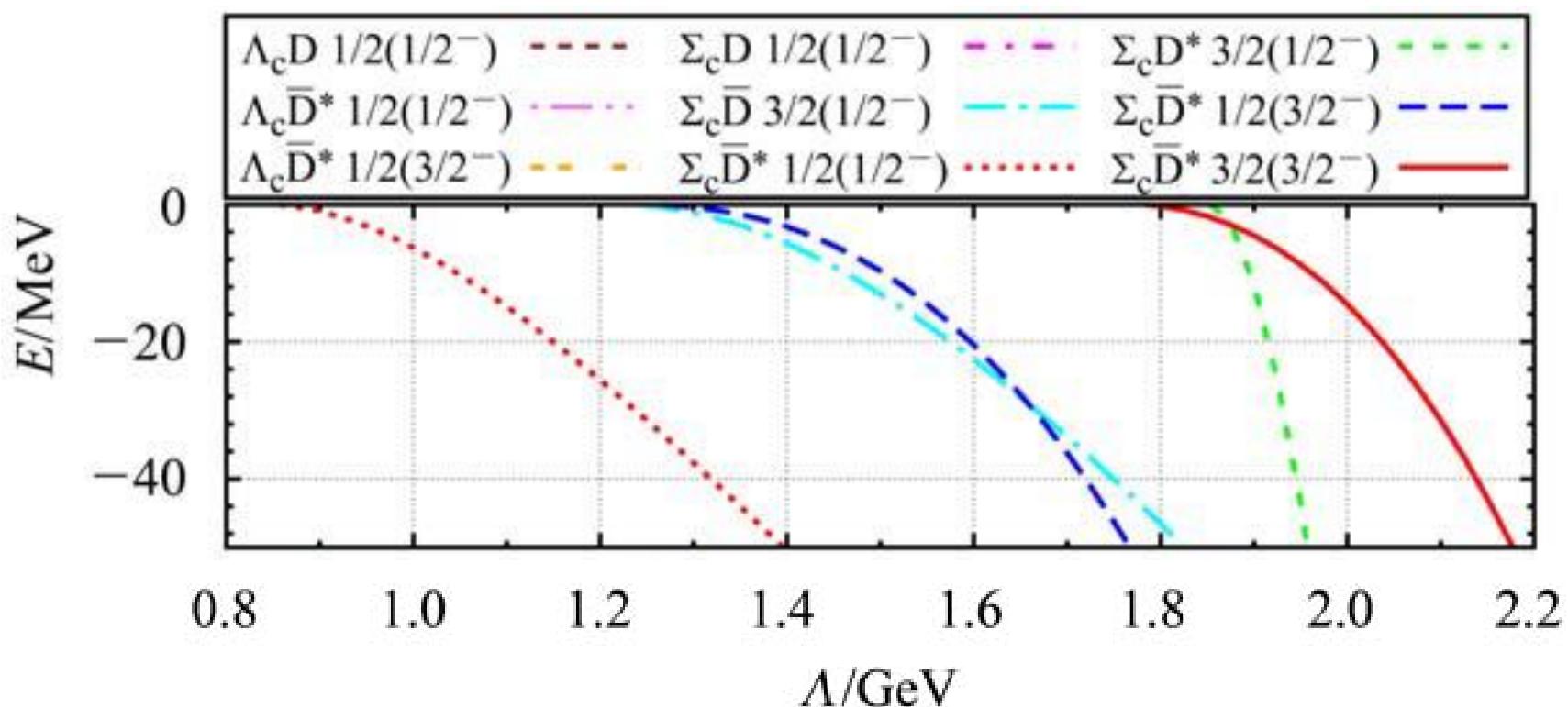
M.Karliner, J.L.Rosner, PRL115(2015)122001:

**Pion exchange** →  $\bar{D}^*\Sigma_c (1/2^-, 3/2^-) \sim 4.5$  GeV

Schoedinger Equation method with  $\pi, \eta, \rho, \omega, \sigma$  exchanges:

$\bar{D}^*\Sigma_c (1/2^-, 3/2^-)$  N\* state --  $4360 \sim 4460$  MeV

$\bar{D}\Sigma_c (I=3/2)$   $\Delta^*$  state --  $\sim 4300$  MeV



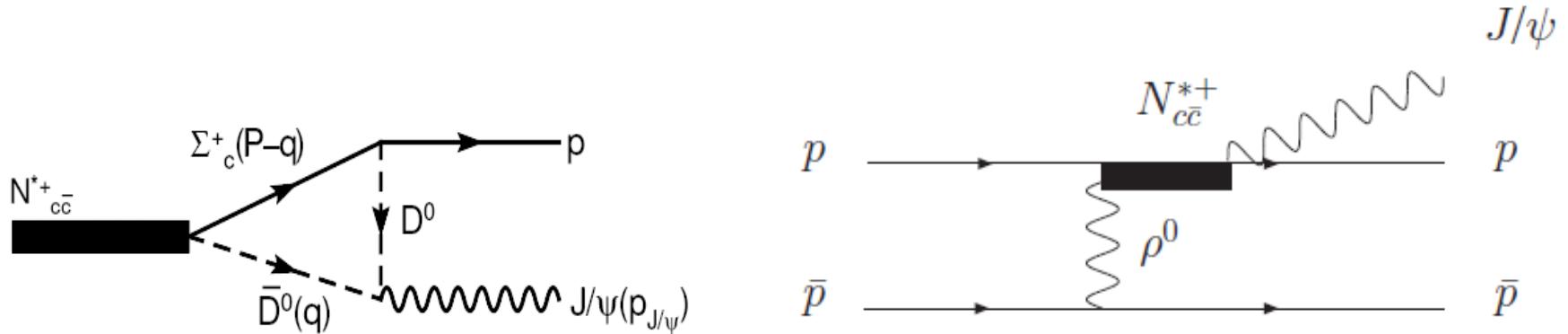
**S.G.Yuan, K.W.Wei, J.He, H.S.Xu, B.S.Zou, “Study of  $\bar{c}cqqq$  five quark system with three kinds of quark-quark hyperfine interaction,”**  
**Eur. Phys. J. A48 (2012) 61**

$J^P$	<i>CM</i>		<i>FS</i>		<i>Inst.</i>	
	<i>udsc̄</i>	<i>uudc̄</i>	<i>udsc̄</i>	<i>uudc̄</i>	<i>udsc̄</i>	<i>uudc̄</i>
$\frac{1}{2}^-$	4273	4267	4084	3933	4209	4114
$\frac{1}{2}^-$	4377	4363	4154	4013	4216	4131
$\frac{1}{2}^-$	4453	4377	4160	4119	4277	4204
$\frac{1}{2}^-$	4469	4471	4171	4136	4295	4207
$\frac{1}{2}^-$	4494	4541	4253	4156	4360	4272
$\frac{1}{2}^-$	4576		4263		4362	
$\frac{1}{2}^-$	4649		4278		4416	
$\frac{3}{2}^-$	4431	<u>4389</u>	4154	4013	4216	4131
$\frac{3}{2}^-$	4503	<u>4445</u>	4171	4119	4295	4204
$\frac{3}{2}^-$	4549	4476	4263	4136	4362	4272
$\frac{3}{2}^-$	4577	4526	4278	4236	4416	<u>4322</u>
$\frac{3}{2}^-$	4629		4362		4461	
$\frac{5}{2}^-$	4719	4616	4362	4236	4461	4322

$J^P$	<i>CM</i>		<i>FS</i>		<i>Inst.</i>	
	<i>udsc̄c̄</i>	<i>uudc̄c̄</i>	<i>udsc̄c̄</i>	<i>uudc̄c̄</i>	<i>udsc̄c̄</i>	<i>uudc̄c̄</i>
$\frac{1}{2}^+$	4622	4456	4291	4138	4487	4396
$\frac{1}{2}^+$	4636	4480	4297	4140	4501	4426
$\frac{1}{2}^+$	4645	4557	4363	4238	4520	4426
$\frac{1}{2}^+$	4658	4581	4439	4320	4540	4470
$\frac{1}{2}^+$	4690	4593	4439	4367	4557	4482
$\frac{1}{2}^+$	4696	4632	4467	4377	4587	4490
$\frac{1}{2}^+$	4714	4654	4469	4404	4590	4517
$\frac{1}{2}^+$	4728	4676	4486	4489	4614	4518
$\frac{1}{2}^+$	4737	4714	4492	4508	4616	4549
$\frac{1}{2}^+$	4766	4720	4510	4515	4626	4566
$\frac{3}{2}^+$	4623	<u>4457</u>	4291	4138	4487	4396
$\frac{3}{2}^+$	4638	4515	4297	4140	4501	4426
$\frac{3}{2}^+$	4680	4561	4363	4238	4520	4426
$\frac{3}{2}^+$	4692	4582	4439	4320	4540	4470
$\frac{3}{2}^+$	4695	4625	4439	4367	4557	4482
$\frac{5}{2}^+$	4705	4539	4297	4140	4501	<u>4426</u>
$\frac{5}{2}^+$	4719	4649	4439	4320	4540	4470
$\frac{5}{2}^+$	4773	4689	4467	4367	4587	4482
$\frac{5}{2}^+$	4793	4696	4486	4404	4615	4490
$\frac{5}{2}^+$	4821	4710	4492	4515	4632	4517
$\frac{7}{2}^+$	4945	4841	4638	4508	4698	4566
$\frac{7}{2}^+$	4955	4862	4671	4551	4712	4634
$\frac{7}{2}^+$	4974	4919	4705	4587	4765	4669
$\frac{7}{2}^+$	5010		4759		4797	

**M(5/2<sup>+</sup>) – M(3/2<sup>-</sup>) : 130 ~300 MeV**

# Prediction for PANDA



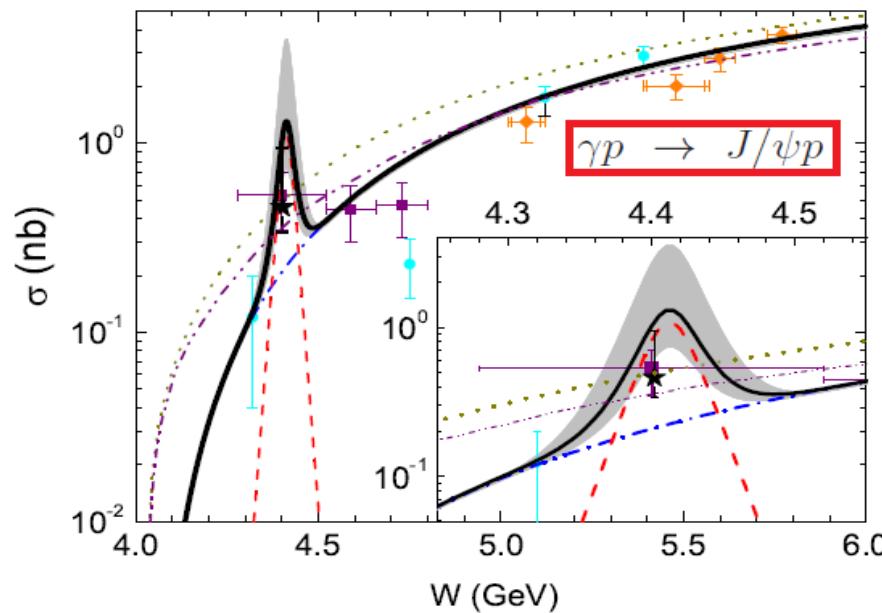
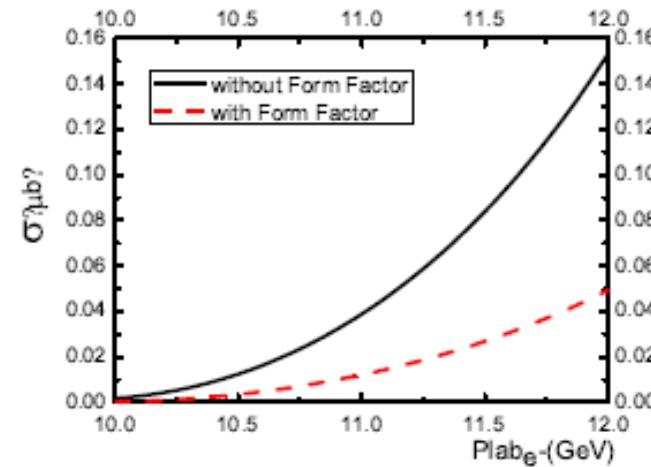
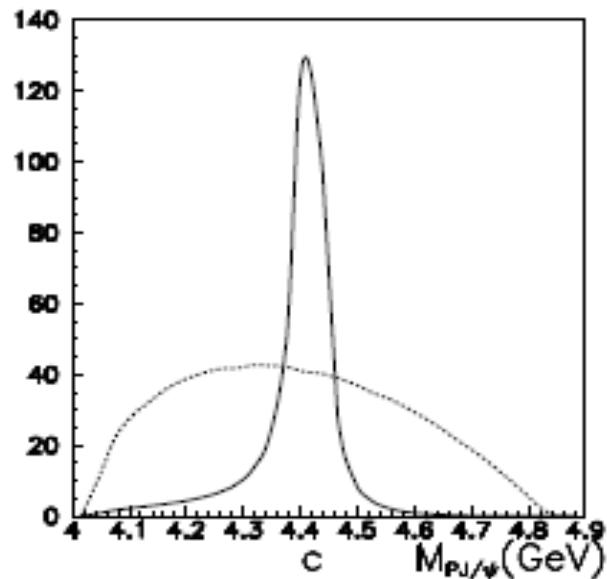
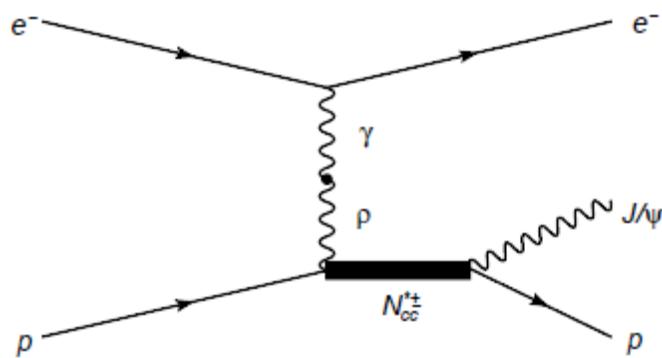
$\bar{p}p \rightarrow \bar{p}p J/\psi > 0.1 \text{ nb}$

> 100 events per day at PANDA/FAIR by  $L=10^{31} \text{ cm}^{-2}\text{s}^{-1}$

These Super-heavy narrow  $N^*$  and  $\Lambda^*$   
can be found at PANDA !

Albrecht Gillitzer@Juelich had a plan to find them at PANDA

# Prediction for 12GeV@JLab



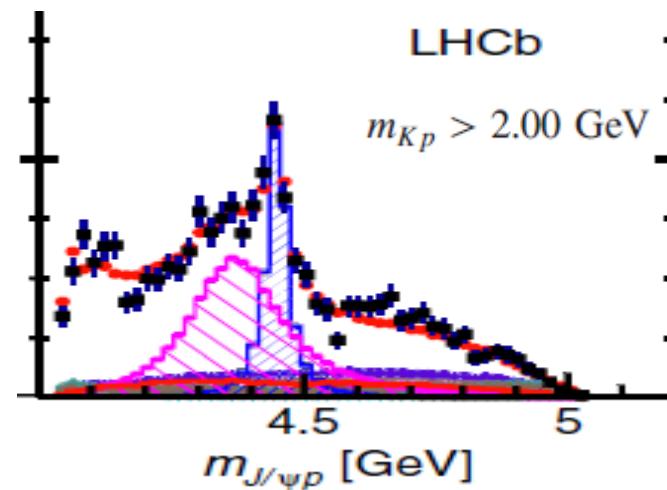
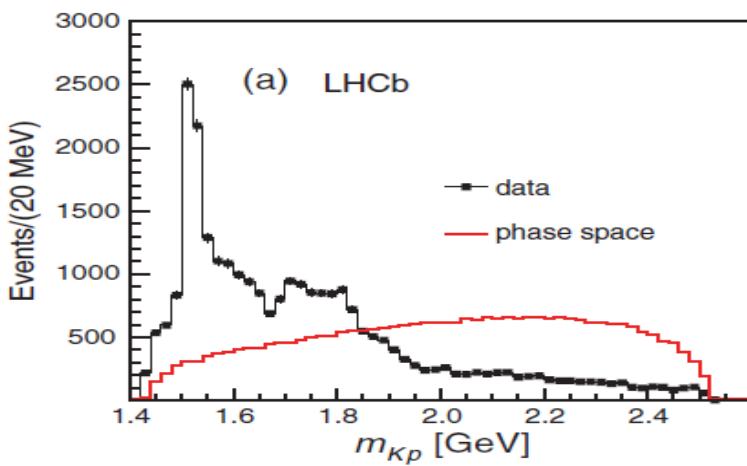
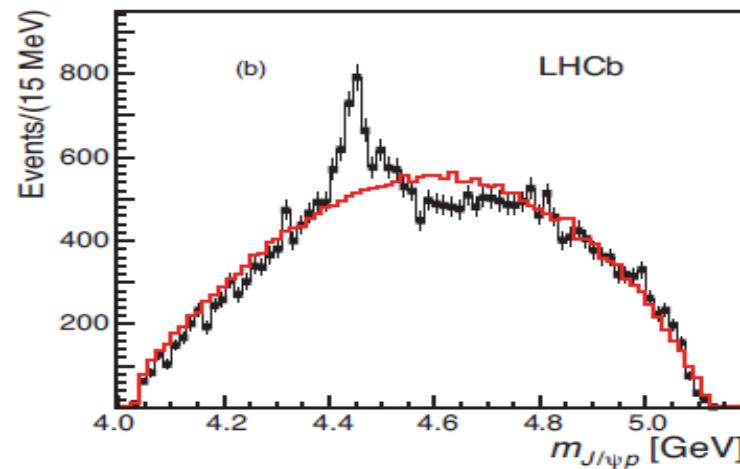
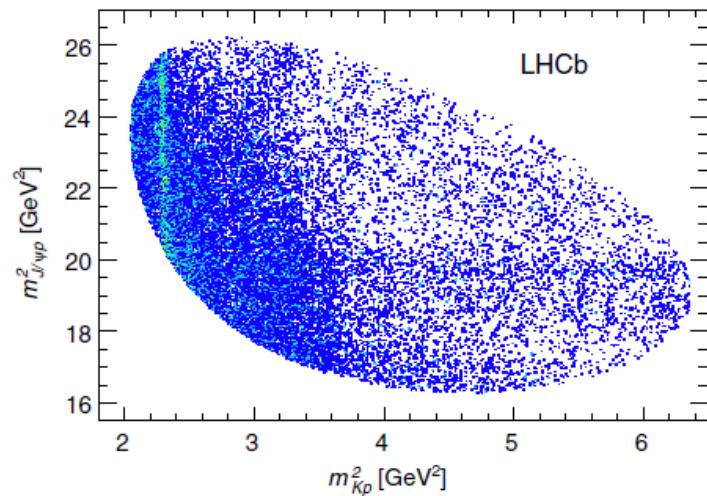
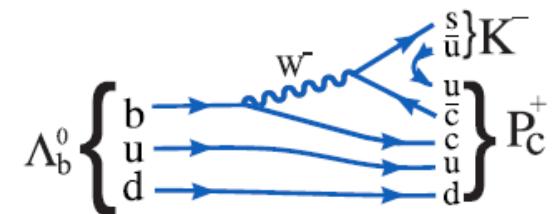
## **Proposals for looking for $N_{cc}^-$ & $\Lambda_{cc}^-$ with $\pi^-$ , K beams at JPARC**

- a) X.Y.Wang, X.R.Chen, “The production of hidden charm baryon  $N^*(4261)$  from  $\pi^- p \rightarrow \eta_c^- n$  reaction”, EPL109 (2015) 41001.
- b) E.J.Garzon, J.J.Xie, “Effects of a  $N_{cc}^-$  resonance with hidden charm in the  $\pi^- p \rightarrow D^- \Sigma_c^+$  reaction near threshold”, PRC 92 (2015) 035201
- c) X.Y.Wang, X.R.Chen, “Production of the superheavy baryon  $\Lambda^*(4209)$  in kaon-induced reaction”, EPJA51 (2015) 85

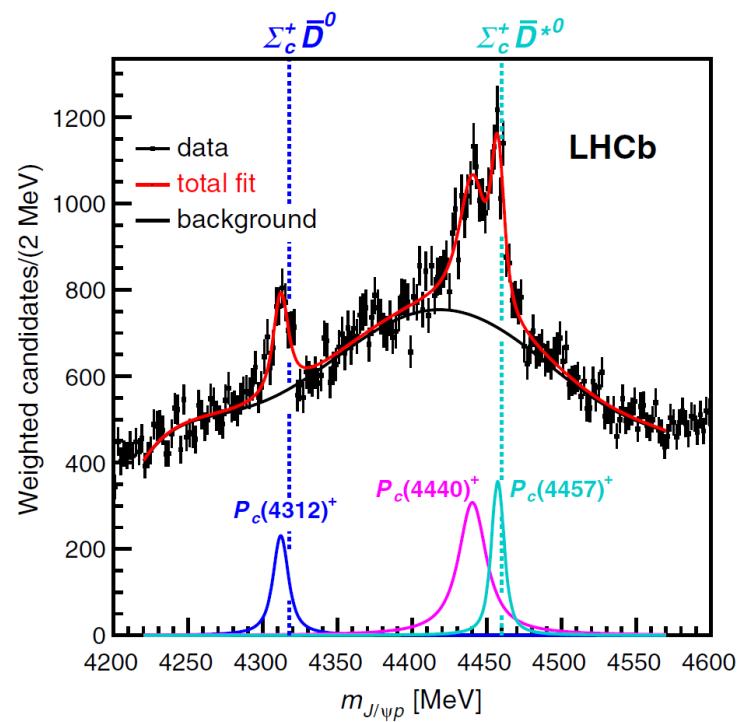
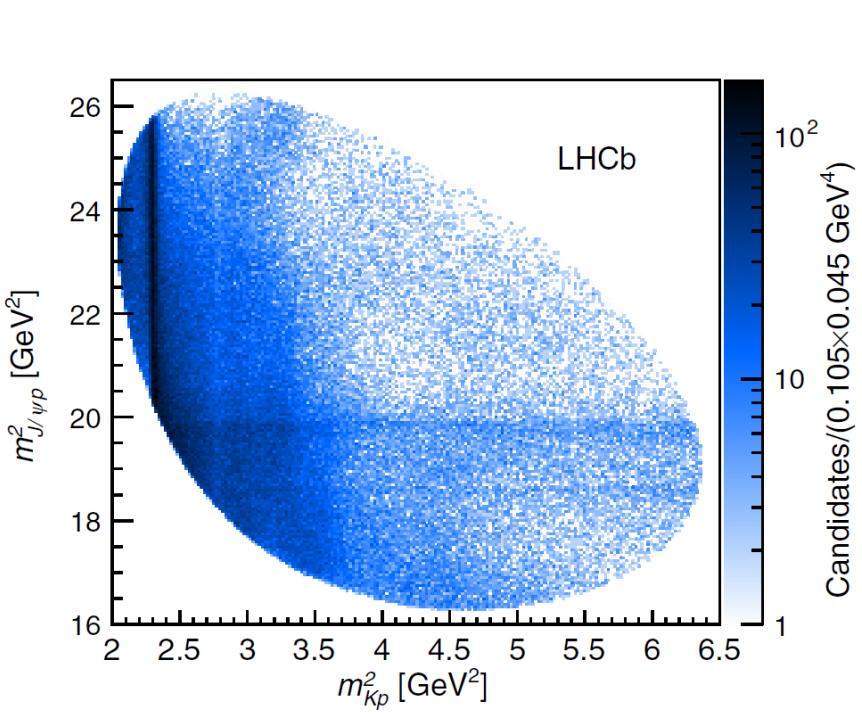
# LHCb observation of $P_c$ pentaquarks

LHCb, Phys.Rev.Lett. 115 (2015) 072001 :

Observation of two  $N^*$  from  $\Lambda_b^0 \rightarrow J/\psi K^- p$



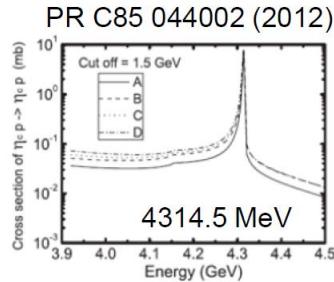
# LHCb, Phys.Rev.Lett. 122 (2019) 222001



State	$M$ [MeV]	$\Gamma$ [MeV]	(95% C.L.)	$\mathcal{R}$ [%]
$P_c(4312)^+$	$4311.9 \pm 0.7^{+6.8}_{-0.6}$	$9.8 \pm 2.7^{+3.7}_{-4.5}$	( $<27$ )	$0.30 \pm 0.07^{+0.34}_{-0.09}$
$P_c(4440)^+$	$4440.3 \pm 1.3^{+4.1}_{-4.7}$	$20.6 \pm 4.9^{+8.7}_{-10.1}$	( $<49$ )	$1.11 \pm 0.33^{+0.22}_{-0.10}$
$P_c(4457)^+$	$4457.3 \pm 0.6^{+4.1}_{-1.7}$	$6.4 \pm 2.0^{+5.7}_{-1.9}$	( $<20$ )	$0.53 \pm 0.16^{+0.15}_{-0.13}$

## Comparison to numerical predictions

- Many theoretical predictions for  $\Sigma_c^+ \bar{D}^{(*)0}$  published before 2015, some in quantitative agreement with the LHCb data
  - Wu,Molina,Oset,Zou, PRL105, 232001 (2010),
  - Wang,Huang,Zhang,Zou, PR C84, 015203 (2011),
  - Yang,Sun,He,Liu,Zhu, Chin. Phys. C36, 6 (2012),
  - Wu, Lee, Zou, PR C85 044002 (2012),
  - Karliner,Rosner, PRL 115, 122001 (2015)



$\Delta E$  – binding energy

Example:

Nucleon resonances with hidden charm in coupled-channels models

Jia-Jun Wu, T.-S. H. Lee, and B. S. Zou  
Phys. Rev. C 85, 044002 – Published 17 April 2012

arXiv:1202.1036

TABLE III: The pole position ( $M - i\Gamma/2$ ) and “binding energy” ( $\Delta E = E_{thr} - M$ ) for different cut-off parameter  $\Lambda$  and spin-parity  $J^P$ . The threshold  $E_{thr}$  is 4320.79 MeV of  $\bar{D}\Sigma_c$  in PB system and 4462.18 MeV of  $D^*\Sigma_c$  in VB system. The unit for the listed numbers is MeV.

$J^P = \frac{1}{2}^-$	$\Lambda$	PB System		VB System	
		$M - i\Gamma/2$	$\Delta E$	$M - i\Gamma/2$	$\Delta E$
650	-	-	-	4462.178 – 0.002 <i>i</i>	0.002
800	$\Delta E(4312) = 5.8^{+1.0}_{-6.8}$ MeV	4318.964 – 0.362 <i>i</i>	1.826	4459.513 – 0.417 <i>i</i>	2.667
1200	4314.531 – 1.448 <i>i</i>	6.259	4454.088 – 1.662 <i>i</i>	8.092	
1500	4301.115 – 5.835 <i>i</i>	19.68	4438.277 – 7.115 <i>i</i>	23.90	
2000	-	-	-	-	-
$J^P = \frac{3}{2}^-$	$\Lambda$	PB System		VB System	
		650	-	-	-
		800	-	-	4462.178 – 0.002 <i>i</i>
		1200	-	-	4459.507 – 0.420 <i>i</i>
		1500	-	-	4454.057 – 1.681 <i>i</i>
		2000	-	-	4438.039 – 7.268 <i>i</i>

$\Lambda$  – cut off on exchanged meson mass.

$\Delta E(4440) = 19.5^{+4.9}_{-4.3}$  MeV

# Progress on $P_c$ states after LHCb observation

Thresholds  $\bar{D}\Sigma_c^*$  (4383MeV),  $\bar{D}^*\Sigma_c$  (4460MeV),  $p\chi_{c1}$  (4449MeV)

## 1) $\bar{D}\Sigma_c^*$ , $\bar{D}^*\Sigma_c$ , $\bar{D}^*\Sigma_c^*$ molecular states

R.Chen, X.Liu, X.Q.Li, S.L.Zhu, PRL115 (2015) 132002;  
L.Roca, J.Nieves, E.Oset, PRD92 (2015) 094003;  
J.He, PLB 753 (2016) 547 ;

## 2) diquark $cu$ & triquark $\bar{c}(ud)$ states

L.Maiani, A.D.Polosa, V.Riquer, PLB749 (2015) 289;  
R.Lebed, PLB749 (2015) 454;  
G.N.Li, M.He, X.G.He, JHEP 1512 (2015) 128;  
R.Zhu, C.F.Qiao, PLB756 (2016) 259;

## 3) Kinematic triangle-singularity

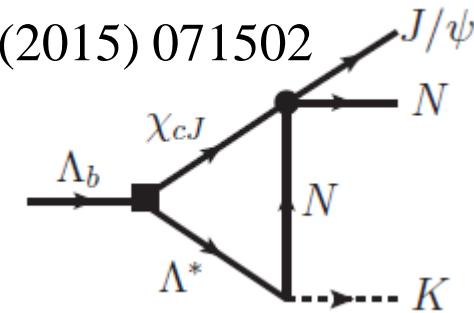
F.K.Guo, Ulf-G.Meißner, W.Wang, Z.Yang, PRD92 (2015) 071502  
X.H.Liu, Q.Wang, Q.Zhao, PLB757 (2016) 231

For comprehensive reviews, cf.:

H.X.Chen, W.Chen, X.Liu, S.L.Zhu, Phys.Rept. 639 (2016) 1

F.K.Guo, C.Hanhart, U.Meissner, Q.Wang, Q.Zhao, B.S.Zou, RMP 90 (2018) 015004

Y.R.Liu, H.X.Chen, W.Chen, X.Liu, S.L.Zhu, Prog.Part.Nucl.Phys. 107 (2019) 237



# Different predictions from various models :

**Diquark model by  
A.Ali et al., PLB793 (2019) 365**

$J^P$	This work	Refs. [22, 23]
$S_{ld} = 0, L = 0$		
$1/2^-$	$3830 \pm 34$	$4086 \pm 42$
	$4150 \pm 29$	$4162 \pm 38$
$3/2^-$	$4240 \pm 29$	$4133 \pm 55$
$S_{ld} = 0, L = 1$		
$1/2^+$	$4030 \pm 39$	$4030 \pm 62$
	$4351 \pm 35$	$4141 \pm 44$
	$4430 \pm 35$	$4217 \pm 40$
$3/2^+$	$4040 \pm 39$	
	$4361 \pm 35$	
$S_{ld} = 1$		
	$4440 \pm 35$	
$5/2^+$	$4457 \pm 35$	$4510 \pm 57$

**Hadrocharmonium states by  
M.Eides et al., arXiv:1904.11616**

Constituents	Binding energy [MeV]	Mass [MeV]	Spin-parity
$\eta_c(2S)N$	176.1	4401	$1/2^-$
$\chi_{c1}(1P)N$	44.2	4406	$3/2^+, 1/2^+$
$h_c(1P)N$	43.9	4421	$1/2^+, 3/2^+$
$\chi_{c2}(1P)N$	43.7	4452	$5/2^+, 3/2^+$

## $\bar{D}\Sigma_c^*$ , $\bar{D}^*\Sigma_c$ , $\bar{D}^*\Sigma_c^*$ bound states

[1] R.Chen, X.Liu, X.Q.Li, S.L.Zhu, PRL115 (2015) 132002;

$$P_c^+(4380) \quad \text{--} \quad \bar{D}^*\Sigma_c \quad 3/2^- ; \quad P_c^+(4450) \quad \text{--} \quad \bar{D}^*\Sigma_c^* \quad 5/2^-$$

[2] Y.Yamaguchi, E. Santopinto, PRD96 (2017) 014018

$$P_c^+(4380) \quad \text{--} \quad \bar{D}^{(*)}\Sigma_c^{(*)-} \quad \bar{D}^{(*)}\Lambda_c \quad 3/2^+ ; \quad P_c^+(4450) \quad \text{--} \quad \bar{D}^{(*)}\Sigma_c^{(*)-} \quad \bar{D}^{(*)}\Lambda_c \quad 5/2^-$$

[3] J.He, PLB 753 (2016)547 ; PRD95 (2017)074004

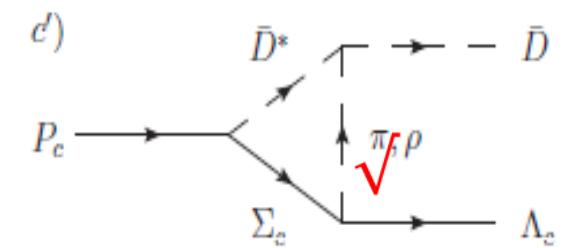
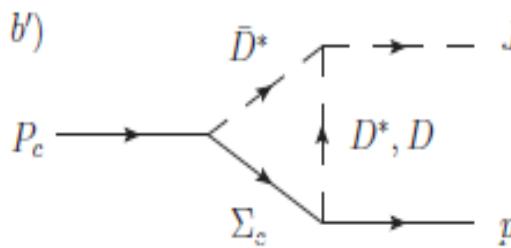
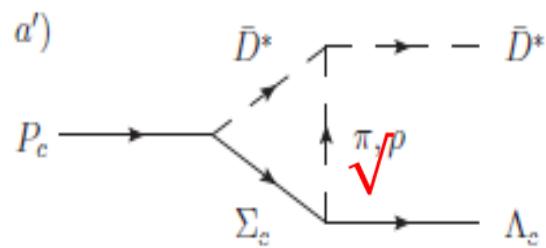
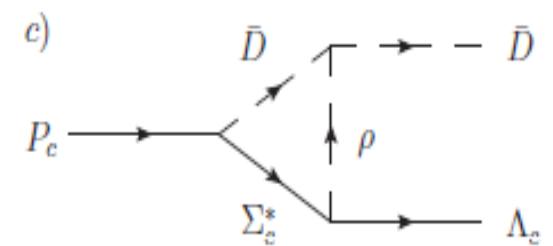
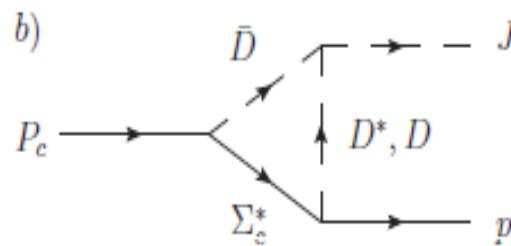
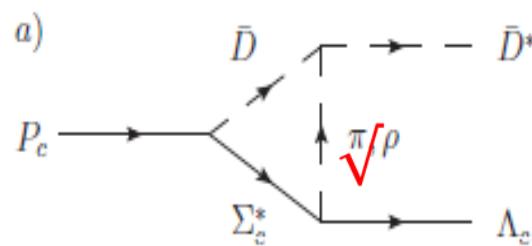
Y.H.Lin, C.W.Shen, F.K.Guo, B.S.Zou, PRD95(2017)114017

$$P_c^+(4380) \quad \text{--} \quad \bar{D}\Sigma_c^*/ \quad \bar{D}^*\Sigma_c \quad 3/2^- ; \quad P_c^+(4450) \quad \text{--} \quad \bar{D}^*\Sigma_c \quad 5/2^+$$

→ Different predictions to be checked !

# Disentangling $\bar{D}\Sigma_c^*$ / $\bar{D}^*\Sigma_c$ nature of $P_c^+$ states from their decays

Y.H.Lin, C.W.Shen, F.K.Guo, B.S.Zou, PRD95(2017)114017



**One pion exchange is very important !**

$\bar{D}\Sigma_c^*$  &  $\bar{D}^*\Sigma_c^*$  are much broader than  $\bar{D}\Sigma_c$  &  $\bar{D}^*\Sigma_c$  states

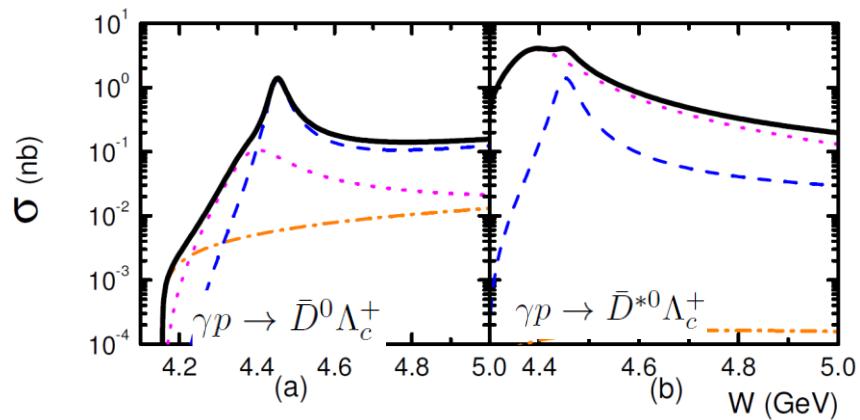
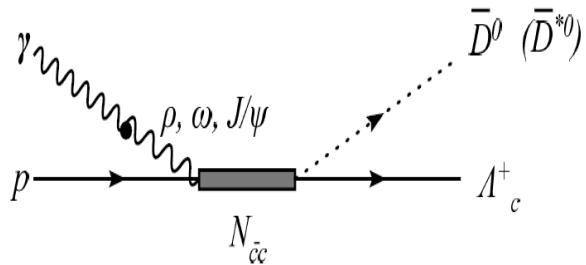
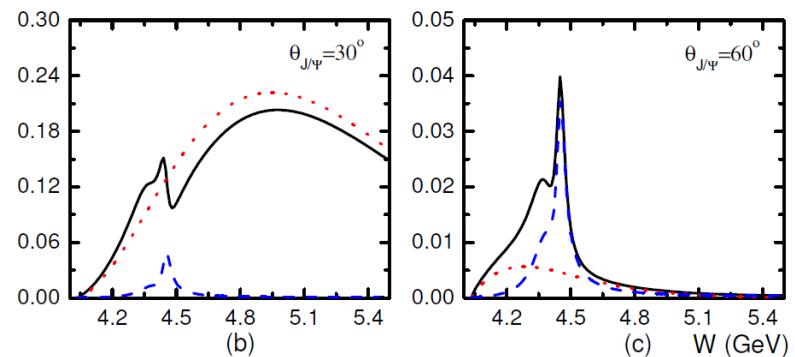
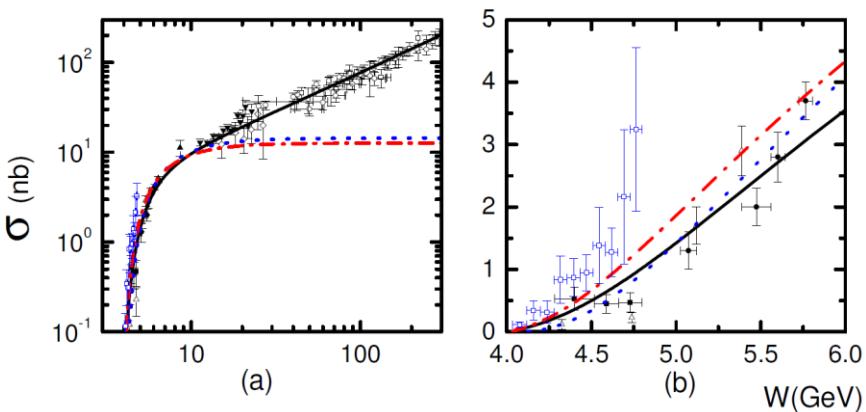
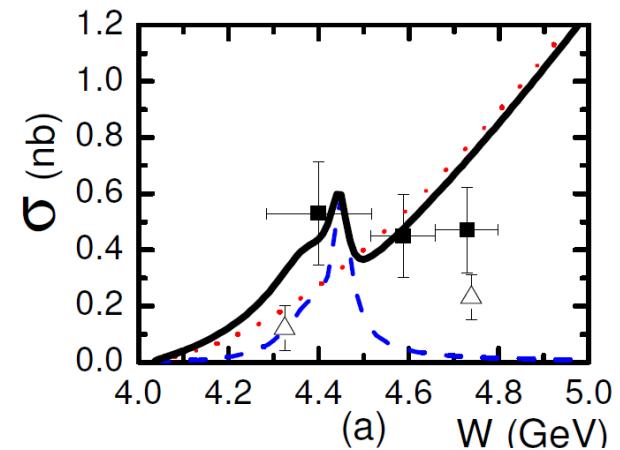
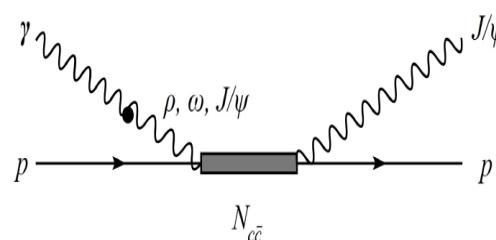
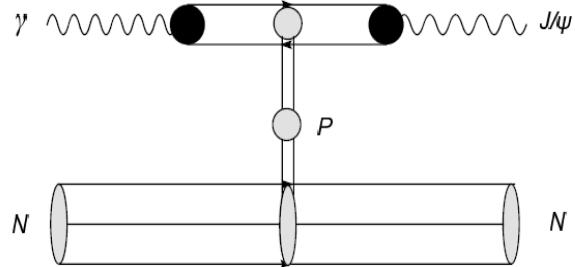
## Partial decay widths of $P_c^+(4380)$ & $P_c^+(4450)$

Mode	Widths (MeV)			
	$P_c(4380)$		$P_c(4450)$	
	$\bar{D}\Sigma_c^*(\frac{3}{2}^-)$	$\bar{D}^*\Sigma_c(\frac{3}{2}^-)$	$\bar{D}^*\Sigma_c(\frac{3}{2}^-)$	$\bar{D}^*\Sigma_c(\frac{5}{2}^+)$
$\bar{D}^*\Lambda_c$	131.3 ✓	35.3 ✓	72.3 ✓	20.5 ✓
$J/\psi p$	3.8	16.6	16.3	4.0
$\bar{D}\Lambda_c$	1.2	17.0 ✓	41.4 ✓	18.8 ✓
$\pi N$	0.06	0.07	0.07	0.2
$\chi_{c0}p$	0.9	0.004	0.02	0.002
$\eta_c p$	0.2	0.09	0.1	0.04
$\rho N$	1.4	0.15	0.14	0.3
$\omega p$	5.3	0.6	0.5	0.3
$\bar{D}\Sigma_c$	0.01	0.1	1.2	0.8
$\bar{D}\Sigma_c^*$	...	...	7.7	1.4
$\bar{D}\Lambda_c\pi$	11.6	...	...	...
Total	144.3	69.9	139.8	46.4

It is very important to study  $P_c \rightarrow \bar{D}^*\Lambda_c$  &  $\bar{D}\Lambda_c$  !

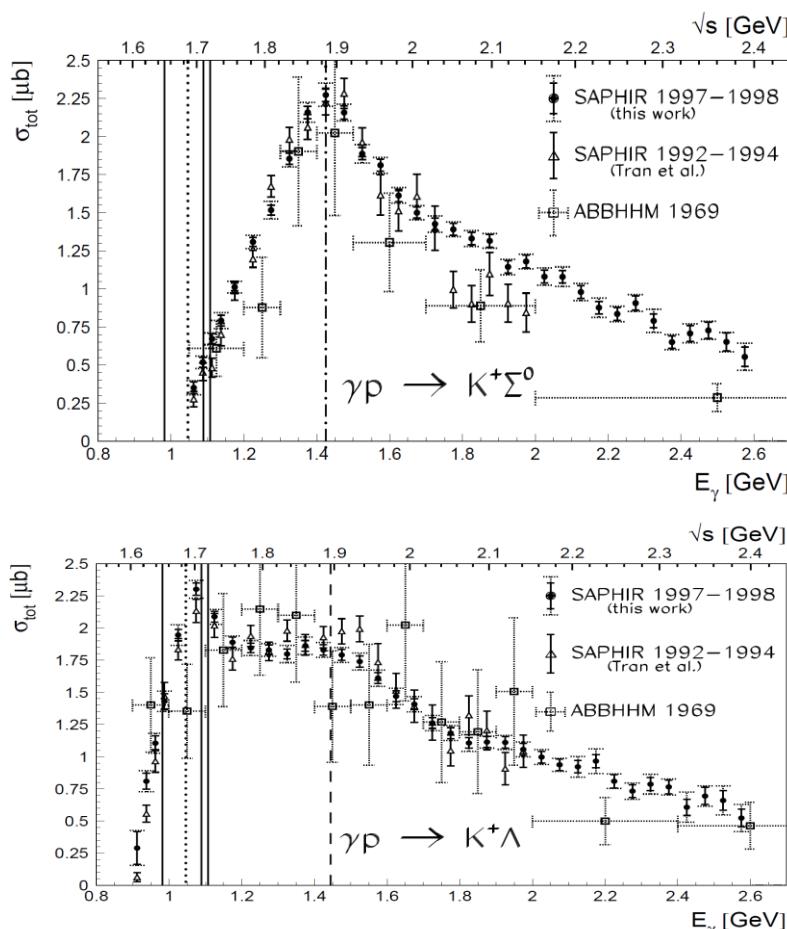
# Pin down $P_c^+$ states from their photo-production

J.J.Wu, T.S.H.Lee, B.S.Zou, arXiv:1906.05375

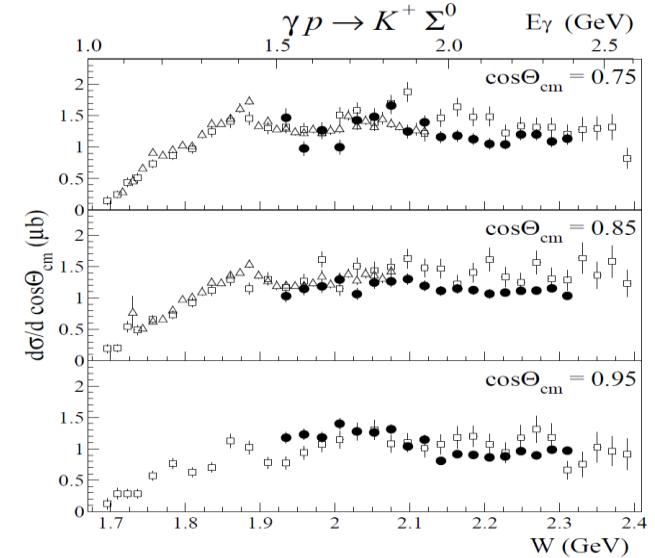


## 4. Strange & beauty partners of $P_c$ states

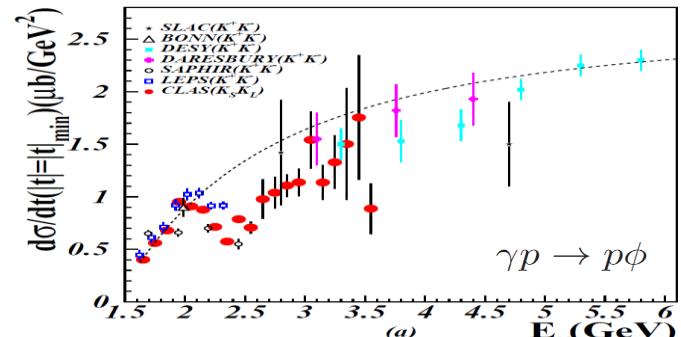
Strangeness partners of  $P_c$  states:  $N^*(1875)$  &  $N^*(2080)$   
 $K\Sigma^* \sim 1880$        $K^*\Sigma \sim 2086$



Glander, K.H. et al. EPJA19 (2004) 251-273



LEPS, PRC73 (2006) 035214



CLAS, PRC89 (2014) 055206

# Strangeness partners of $P_c$ states at BES ?

$N^*(1875)$

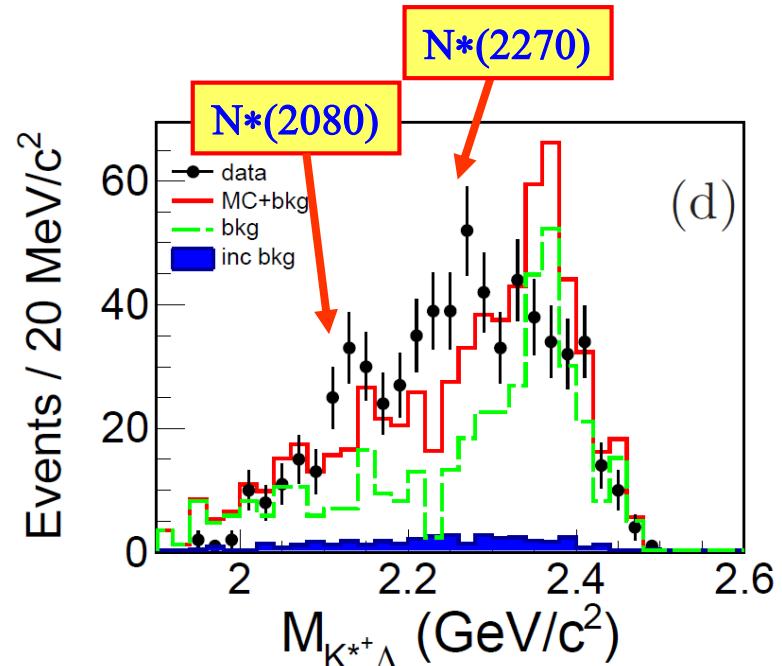
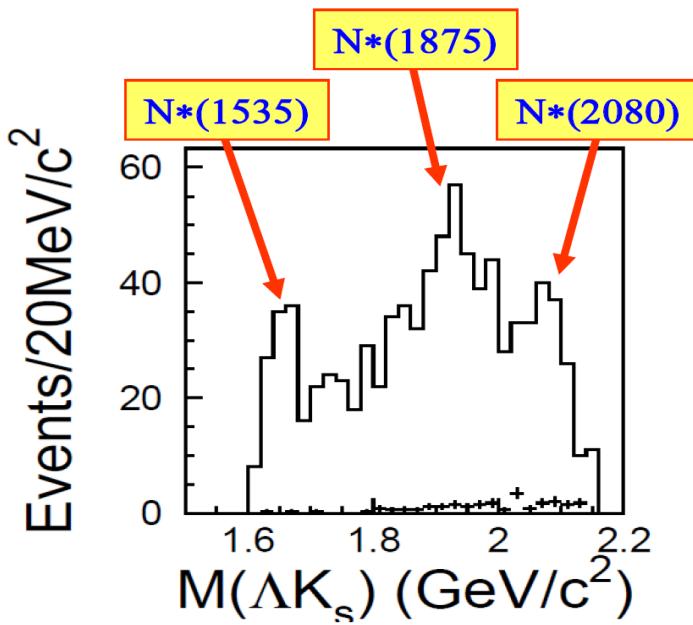
$K\Sigma^* \sim 1880$

$N^*(2080)$

$K^*\Sigma \sim 2086$

$N^*(2270)$

$K^*\Sigma^* \sim 2280$



$$J/\psi \rightarrow n K_s^0 \bar{\Lambda}$$

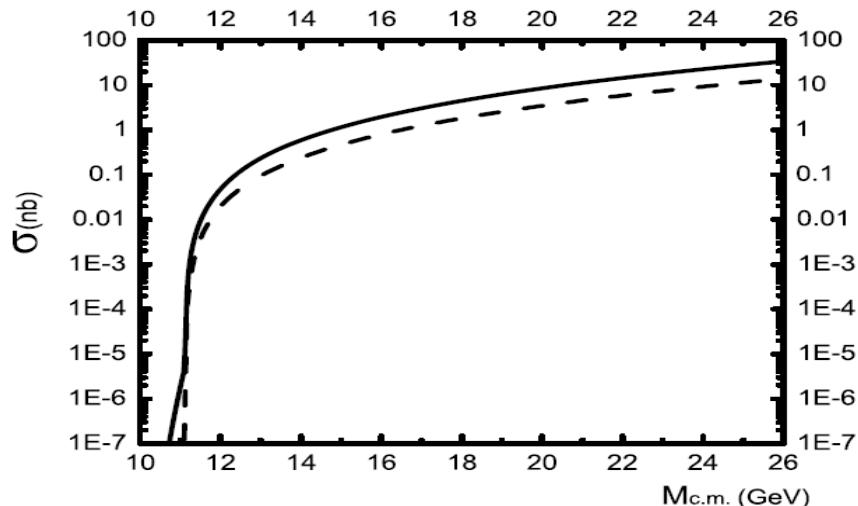
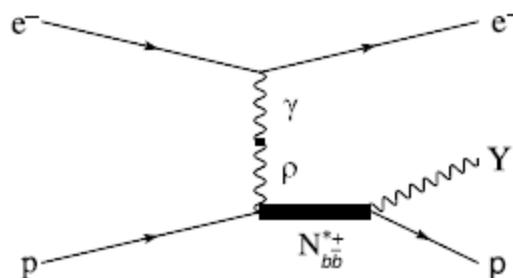
BESII, PLB659 (2008) 789

$$\chi_{c0} \rightarrow \bar{p} K^{*+} \Lambda + c.c.$$

BESIII, arXiv:1908. 02979

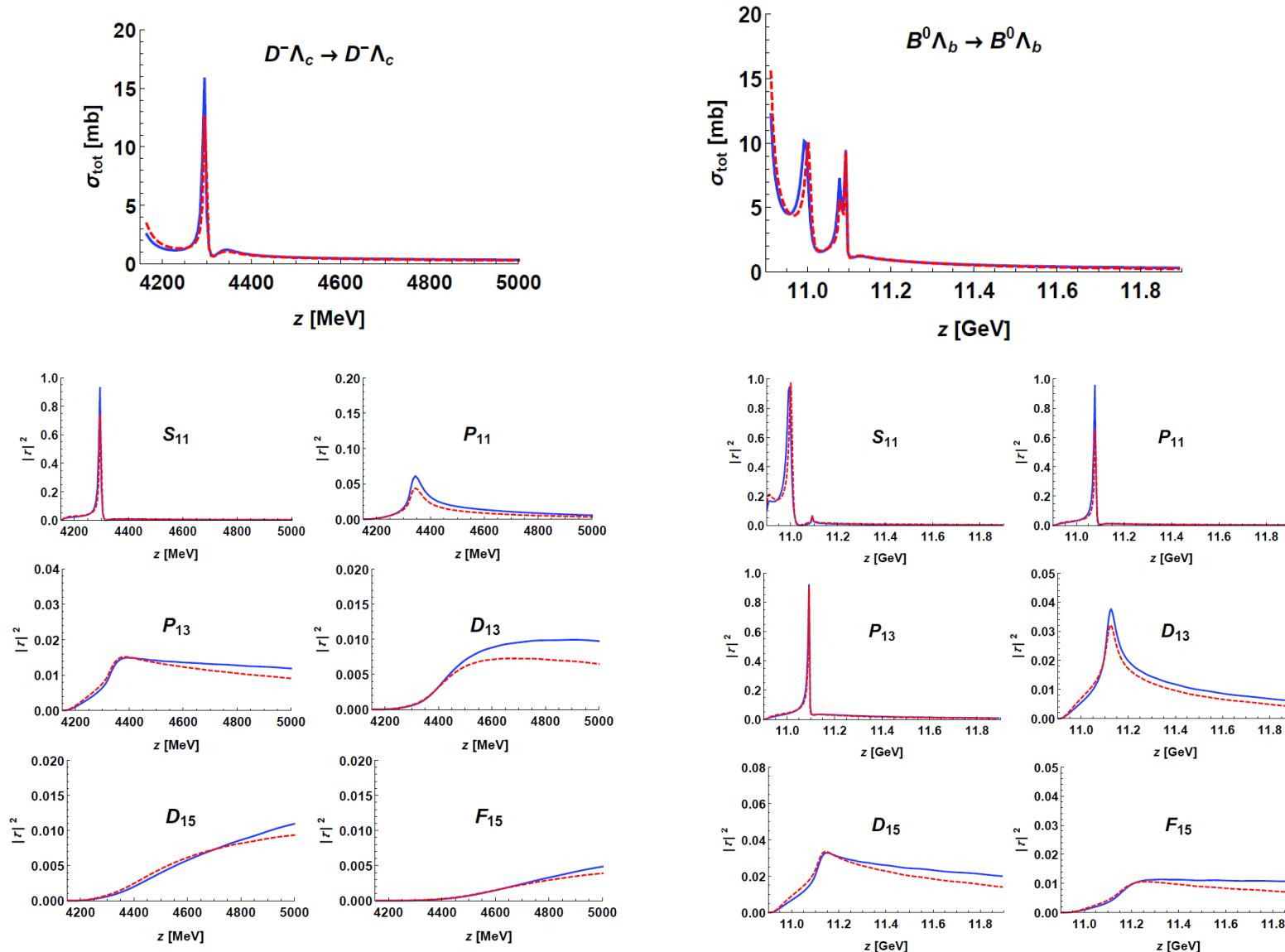
Prediction of super-heavy  $N^*$  and  $\Lambda^*$  resonances with hidden beautyJia-Jun Wu <sup>a,\*</sup>, Lu Zhao <sup>a</sup>, B.S. Zou <sup>a,b</sup>

$M$ (MeV)	$\Gamma$ (MeV)	$\Gamma_i$ (MeV)	$\pi N$	$\eta N$	$\eta' N$	$K \Sigma$	$\eta_b N$	
11 052	1.38		0.10	0.21	0.11	0.42	0.52	<b>1/2<sup>-</sup></b>
11 100	1.33		$\rho N$	$\omega N$	$K^* \Sigma$	$\gamma N$	0.51	<b>1/2<sup>-</sup>, 3/2<sup>-</sup></b>



# $\bar{D}\Lambda_c - \bar{D}\Sigma_c$ and $B\Lambda_b - B\Sigma_b$ dynamical coupled channel study

C.W.Shen, Roechen, Meissner, Zou, CPC42(2018) 023106



More pentaquarks with hidden beauty than with hidden charm

# Decay behavior of $P_s$ & $P_b$ pentaquark states

Y.H.Lin, C.W.Shen, B.S.Zou, NPA980(2018)21

Mode	Widths (MeV)					
	$J^P = 3/2^-$		$J^P = 1/2^-$		$N(1875)$	$K\Sigma^*$
	$N(2080)$	$K^*\Sigma$	$N(2080)$	$K^*\Sigma$		
$N\sigma(500)$	2.6	0.05	0.3			
$\pi N$	3.8	0.2	22.7			
$\rho N$	2.3	3.8	6.1			
$\omega p$	6.6	11.3	18.2			
$K\Sigma$	0.03	1.4	9.1			
$K\Lambda$	0.7	3.7	19.3			
$\eta p$	0.6	0.4	1.8			
$\pi\Delta$	201.4	82.6	46.9			
$K^*\Lambda$	-	2.4	7.9			
$\phi p$	-	19.2	27.0			
$K\Sigma^*$	-	7.3	1.3			
$K\Lambda(1520)$	-	0.1	1.3			
$K\Lambda(1405)$	-	8.0	8.8			
$K\pi\Lambda$	10.1	-	-			
$K\pi\Sigma$	-	41.3	46.1			
Total	228.2	181.7	216.8			

Mode	Widths (MeV)		
	$J^P = 3/2^-$		$J^P = 1/2^-$
	$B\Sigma_b^*$	$B^*\Sigma_b$	$B^*\Sigma_b$
$B^*\Lambda_b$	271.1	19.9	167.0
$\Upsilon p$	0.3	0.04	0.1
$\rho N$	5.5	0.02	0.1
$\omega p$	20.9	0.07	0.4
$B\Lambda_b$	-	7.3	135.9
$B\Sigma_b$	-	-	-
$\eta_b p$	0.02	0.0001	0.0009
$\chi_{b0} p$	1.4	0.0008	0.2
$\pi N$	0.7	0.005	0.003
$B\Sigma_b^*$	-	-	-
Total	299.9	27.4	303.8

## Guidance for $P_s$ & $P_b$ search

Decay behaviors of possible  $\Lambda_{c\bar{c}}$  states in hadronic molecule pictures

C.W.Shen, J.J.Wu, B.S.Zou (2019), ArXiv:1906.03896

## Guidance for $P_{sc}$ search

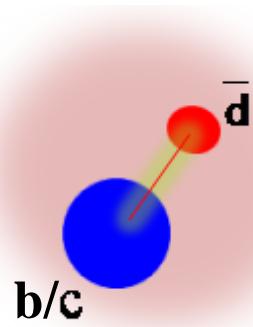
## 5. Prospects

◆ my favorite strategy for hadron spectroscopy:

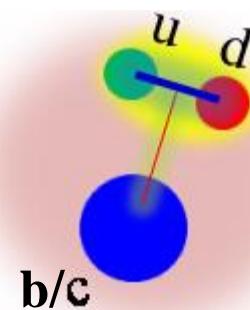
$\bar{c}cuud$  &  $\bar{c}cuds$   $\rightarrow$  sss -  $\bar{q}qsss \rightarrow cqq$  -  $\bar{q}qcqq$   
 $\rightarrow$  hyperons  $\rightarrow$  light baryons

$\bar{c}\bar{c}$   $\bar{u}d$  &  $\bar{c}\bar{s}$   $\bar{u}d$   $\rightarrow$   $\bar{c}\bar{c}$  -  $\bar{q}q$   $\bar{c}\bar{c} \rightarrow \bar{c}q$  -  $\bar{c}q$   $\bar{q}q$   
 $\rightarrow$  K mesons  $\rightarrow$  light mesons

s  $\rightarrow$  c  $\rightarrow$  b



charm & beauty meson



charm & beauty baryon

- New penta-quark spectroscopy provides a new ideal platform for understanding multiquark dynamics
- Further experimental confirmation and extension for whole penta-quark spectroscopy from  $\gamma N$ ,  $\pi N$ ,  $KN$ ,  $e^+e^- \rightarrow \bar{\Lambda}_b \Lambda_b$ , etc.  
 $ep/\gamma p @ JLab$ ,  $\pi 10/K10 @ JPARC$ ,  $BelleII$ ,  $BESIII$ ,  $Eic/EicC$ ,  
 $PANDA @ FAIR$ ,  $STCF$  etc. may play important role here!

Thank you for  
your attention!