QCD phase diagram and its dualities

Roman N. Zhokhov
IZMIRAN, IHEP

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Russian Science Foundation

БАЗИС
Фонд развития теоретической физики и математики

K.G. Klimenko, IHEP
T.G. Khunjua, MSU, University of Georgia, GAS

in the broad sense our group stems from
Department of Theoretical Physics, Moscow State University
Prof. V. Ch. Zhukovsky

strong connections with
prof. D. Ebert, Humboldt University of Berlin

details can be found in

Phys.Rev. D100 (2019) no.3, 034009 arXiv:1907.04151 [hep-ph]
JHEP 1906 (2019) 006 arXiv:1901.02855 [hep-ph]
Eur.Phys.J. C79 (2019) no.2, 151, arXiv:1812.00772 [hep-ph],
Phys.Rev. D98 (2018) no.5, 054030 arXiv:1804.01014 [hep-ph],
Phys.Rev. D97 (2018) no.5, 054036 arXiv:1710.09706 [hep-ph]
Phys.Rev. D95 (2017) no.10, 105010 arXiv:1704.01477 [hep-ph]

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- Russian Science Foundation (RSF)

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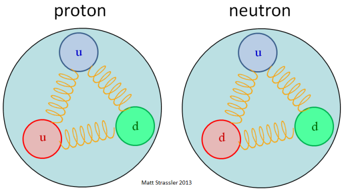
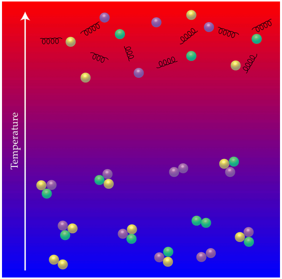
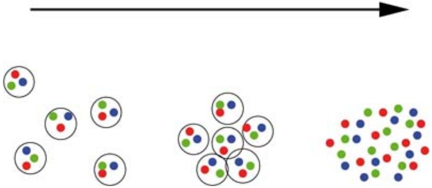
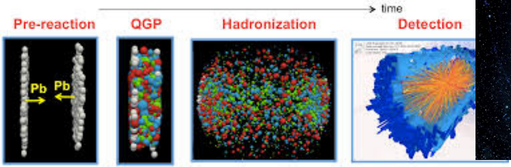
Hadronic, quark matter

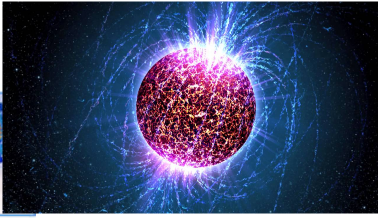


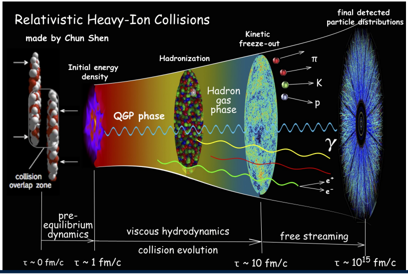


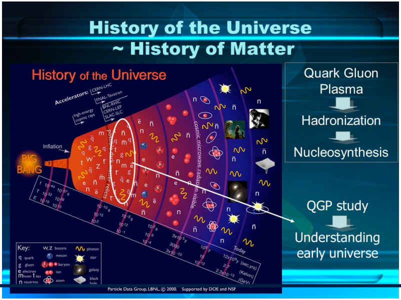


QCD at extreme conditions

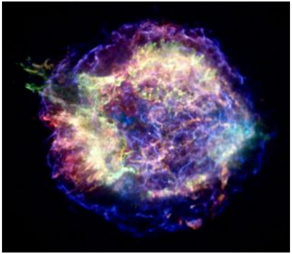


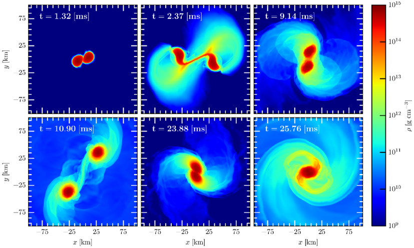


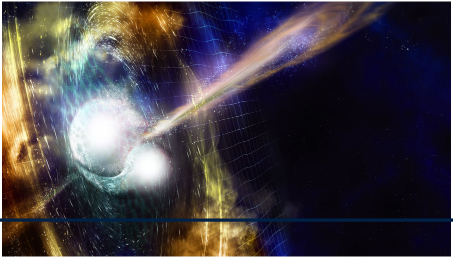


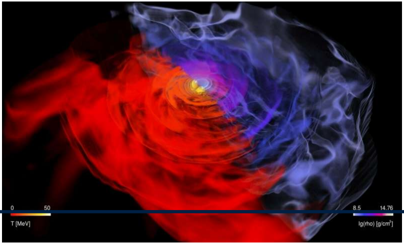


QCD at extreme conditions: even more

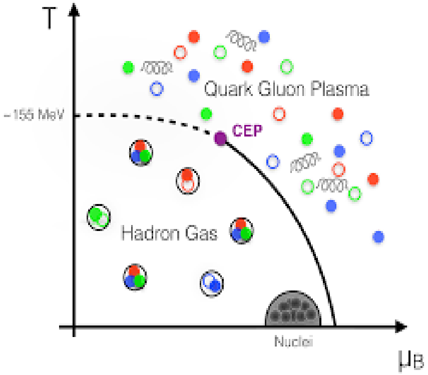








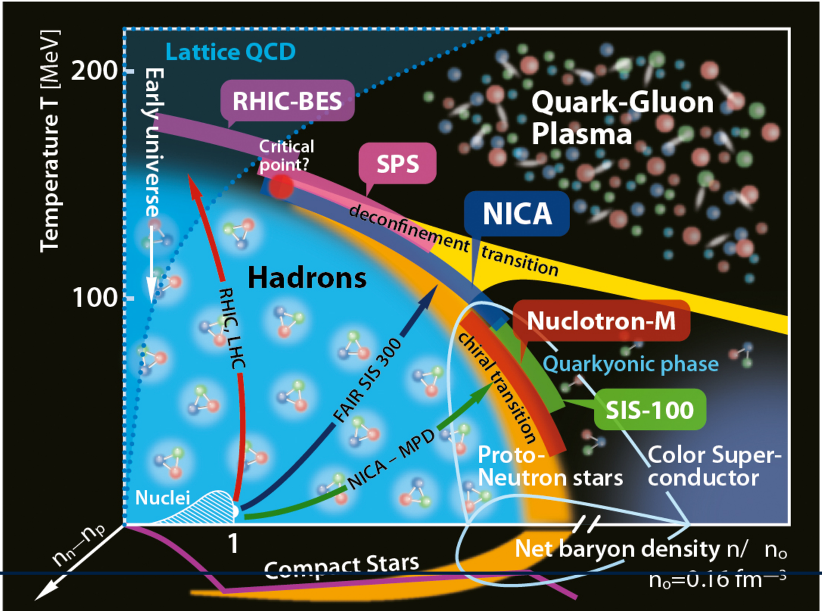
QCD Phase Diagram



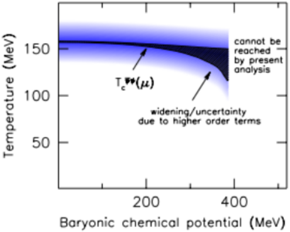
Two phase transitions

- confinement-deconfinement
- chiral symmetry breaking phase—chiral symmetric phase

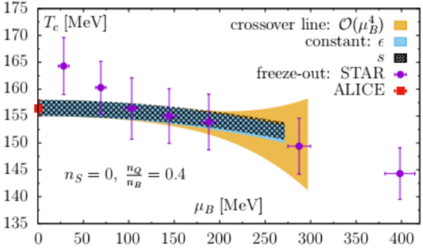
QCD phase diagram



- Curvature of pseudocritical line
  $T_c(\mu_B) = T_c(0) - A_2\mu_B^2 + A_4\mu_B^4 + O(\mu_B^6)$



[R. Bellwied et al., 2015]

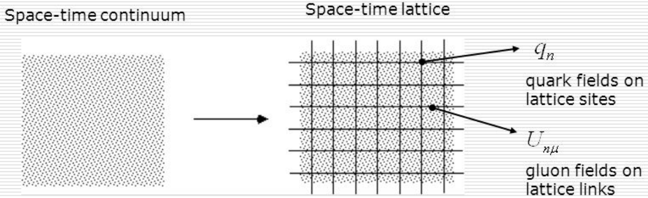


[HotQCD, 2018]

from A. Kotov talk, ITEP lattice group

QCD on a space-time lattice

K. G. Wilson 1974



- Feynman path integral
  - Action $S_{QCD} = \frac{1}{g_s^2}\sum_P tr(UUUU) + \sum_f \bar{q}_f(\gamma\cdot U + m_f)q_f$
  - Physical quantities as integral averages

Monte Carlo Evaluation of the path integral

$$\langle O(U,\bar{q},q)\rangle = \frac{1}{Z}\int \prod_{n\mu} dU_{n\mu} \prod_n d\bar{q}_n dq_n\, O(U,(U,\bar{q},q))\, e^{-S_{QCD}}$$

It is well known that **at non-zero baryon chemical potential** μ_B **lattice simulation** is quite challenging due to the **sign problem**
complex determinant

$$(Det(D(\mu)))^\dagger = Det(D(-\mu^\dagger))$$

NJL model can be considered as **effective field theory for QCD**.

the model is **nonrenormalizable**
Valid up to $E < \Lambda \approx 1$ GeV

Parameters G, Λ, m_0

chiral limit $m_0 = 0$

in many cases chiral limit is a very good approximation

dof– **quarks**
no gluons only **four-fermion interaction**
attractive feature — dynamical CSB
the main drawback – lack of confinement (PNJL)

Nambu–Jona-Lasinio model

$$\mathcal{L} = \bar{q}\gamma^{\nu}\mathrm{i}\partial_{\nu}q + \frac{G}{N_c}\Big[(\bar{q}q)^2 + (\bar{q}\mathrm{i}\gamma^5 q)^2\Big]$$

$$q \to e^{i\gamma_5\alpha}q$$

Equivalent Lagrangian

$$\widetilde{\mathcal{L}} = \bar{q}\Big[\gamma^{\rho}\mathrm{i}\partial_{\rho} - \sigma - \mathrm{i}\gamma^5\pi\Big]q - \frac{N_c}{4G}\Big[\sigma^2 + \pi^2\Big], \qquad \sigma(x) = -2\frac{G}{N_c}(\bar{q}q)$$

Chiral symmetry breaking

$1/N_c$ expansion, leading order or MF

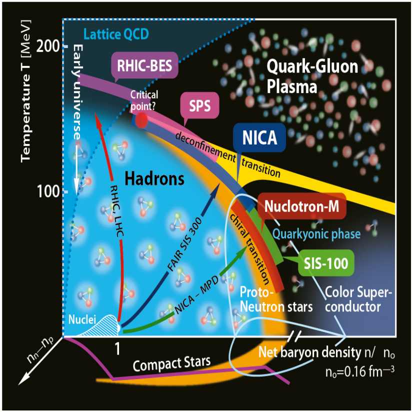
$$\langle\bar{q}q\rangle \neq 0$$

$$\langle\sigma\rangle \neq 0 \quad \longrightarrow \quad \widetilde{\mathcal{L}} = \bar{q}\Big[\gamma^{\rho}\mathrm{i}\partial_{\rho} - \langle\sigma\rangle\Big]q$$

Methods of dealing with QCD

QCD at T and μ
(QCD at extreme conditions)

- neutron stars
- heavy ion collisions
- Early Universe
- Supernovae
- NS merger



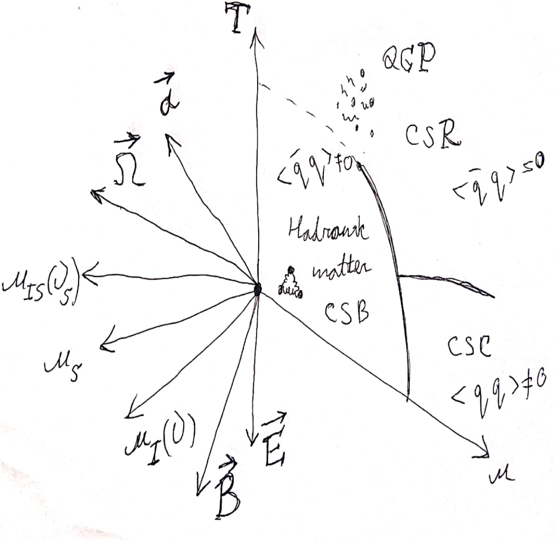
Methods of dealing with QCD

- First principle calcultion – lattice QCD
- Effective models
- .....

More external conditions to QCD

More than just QCD at (μ, T)

- more chemical potentials μ_i
- magnetic fields
- rotation of the system $\vec{\Omega}$
- acceleration $\vec{a}$
- finite size effects (finite volume and boundary conditions)



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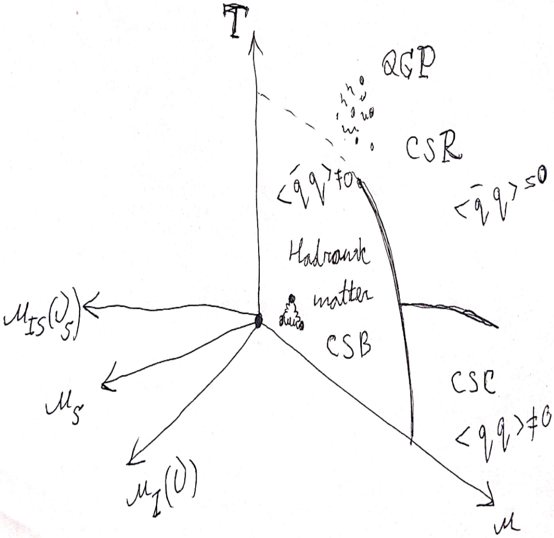


Baryon chemical potential μ_B

Allow to consider systems with non-zero baryon densities.

$$\frac{\mu_B}{3}\bar{q}\gamma^0 q = \mu\bar{q}\gamma^0 q,$$

Baryon chemical potential μ_B

Allow to consider systems with non-zero baryon densities.

$$\frac{\mu_B}{3}\bar{q}\gamma^0 q = \mu\bar{q}\gamma^0 q,$$

Isotopic chemical potential μ_I

Allow to consider systems with isospin imbalance ($n_n \neq n_p$).

$$\frac{\mu_I}{2}\bar{q}\gamma^0\tau_3 q = \nu\left(\bar{q}\gamma^0\tau_3 q\right)$$

$$n_I = n_u - n_d \quad \longleftrightarrow \quad \mu_I = \mu_u - \mu_d$$

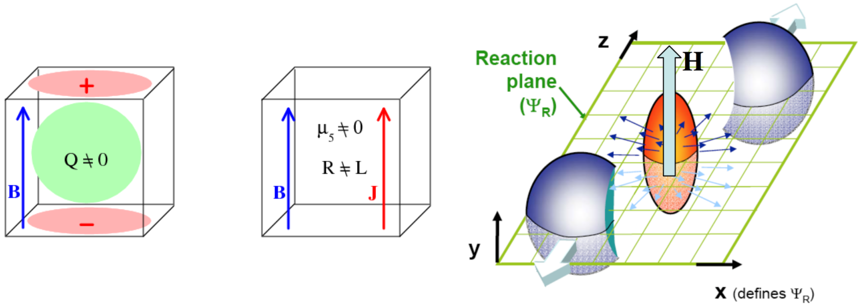
chiral (axial) chemical potential

Allow to consider systems with chiral imbalance (difference between densities of left-handed and right-handed quarks).

$$n_5 = n_R - n_L \quad \longleftrightarrow \quad \mu_5 = \mu_R - \mu_L$$

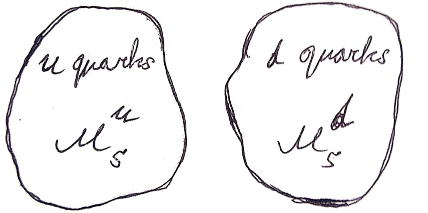
The corresponding term in the Lagrangian is

$$\mu_5 \bar{q} \gamma^0 \gamma^5 q$$



$$\vec{J} \sim \mu_5 \vec{B},$$

A. Vilenkin, PhysRevD.22.3080,
K. Fukushima, D. E. Kharzeev and H. J. Warringa, Phys. Rev. D **78** (2008) 074033



$$\mu_5^u \neq \mu_5^d \qquad \text{and} \qquad \mu_{I5} = \mu_5^u - \mu_5^d$$

Term in the Lagrangian — $\frac{\mu_{I5}}{2}\bar{q}\tau_3\gamma^0\gamma^5 q = \nu_5(\bar{q}\tau_3\gamma^0\gamma^5 q)$

$$n_{I5} = n_{u5} - n_{d5}, \qquad n_{I5} \quad \longleftrightarrow \quad \nu_5$$

Chiral imbalance n_5 and hence μ_5 can be generated in parallel magnetic and electric fileds

$$\vec{E} \,||\, \vec{B}$$

M. Ruggieri, M. Chernodub, H. Warringa et al

Chiral isospin imbalance n_{I5} and hence μ_{I5} can be generated in parallel magnetic and electric fileds $\vec{E}\,||\,\vec{B}$

μ_{I5} and μ_5 are generated by $\vec{E}\,||\,\vec{B}$

Generation of Chiral imbalance in dense quark matter

Chiral imbalance could appear in dense matter

- Chiral separation effect
  (Thanks to Igor Shovkovy)
- Chiral vortical effect

Different chemical potentials and matter content

$$\mu = \frac{\mu_B}{3}, \qquad \nu = \frac{\mu_I}{2}, \qquad \mu_5, \qquad \nu_5 = \frac{\mu_{I5}}{2}$$

We consider a NJL model, which describes dense quark matter with two massless quark flavors (u and d quarks).

$$\mathcal{L} = \bar{q}\Big[\gamma^\nu \mathrm{i}\partial_\nu + \frac{\mu_B}{3}\gamma^0 + \frac{\mu_I}{2}\tau_3\gamma^0 + \frac{\mu_{I5}}{2}\tau_3\gamma^0\gamma^5 + \mu_5\gamma^0\gamma^5\Big]q +$$

$$\frac{G}{N_c}\Big[(\bar{q}q)^2 + (\bar{q}\mathrm{i}\gamma^5\vec{\tau}q)^2\Big]$$

q is the flavor doublet, $q = (q_u, q_d)^T$, where q_u and q_d are four-component Dirac spinors as well as color N_c-plets; τ_k ($k = 1, 2, 3$) are Pauli matrices.

To find the thermodynamic potential we use a semi-bosonized version of the Lagrangian

$$\widetilde{L} = \bar{q}\Big[i\partial\!\!\!/ + \mu\gamma^0 + \nu\tau_3\gamma^0 + \nu_5\tau_3\gamma^0\gamma^5 + \mu_5\gamma^0\gamma^5 - \sigma - i\gamma^5\pi_a\tau_a\Big]q - \frac{N_c}{4G}\Big[\sigma^2 + \pi_a^2\Big].$$

$$\sigma(x) = -2\frac{G}{N_c}(\bar{q}q); \quad \pi_a(x) = -2\frac{G}{N_c}(\bar{q}i\gamma^5\tau_a q).$$

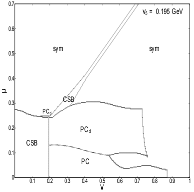
Condansates ansatz $\langle\sigma(x)\rangle$ and $\langle\pi_a(x)\rangle$ do not depend on spacetime coordinates

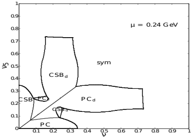
$$\langle\sigma(x)\rangle = M, \quad \langle\pi_1(x)\rangle = \Delta, \quad \langle\pi_2(x)\rangle = 0, \quad \langle\pi_3(x)\rangle = 0.$$

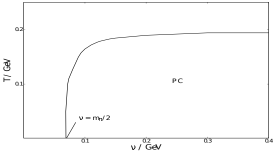
where M and Δ are already constant quantities.

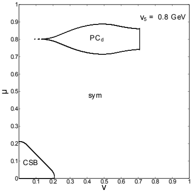
Phase diagram, lots of plots

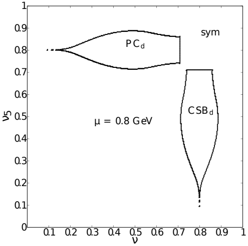
Symmetry 2019, 11(6), 778

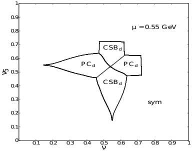












Chiral imbalance lead to the generation of PC in dense quark matter (PC_d)

Dualities

Dualities

It is not related to holography or gauge/gravity duality

it is the dualities of the phase structures of different systems

The TDP

$$\Omega(T, \mu, \mu_i, ..., \langle \bar{q}q \rangle, ...)$$

The TDP

$$\Omega(T, \mu, \mu_i, ..., \langle \bar{q} q \rangle, ...) \qquad \Omega(T, \mu, \mu_i, ..., M, \Delta, ...)$$

The TDP

$$\Omega(T, \mu, \mu_i, ..., \langle\bar{q}q\rangle, ...) \qquad \Omega(T, \mu, \mu_i, ..., M, \Delta, ...)$$

The TDP (phase daigram) is invariant under

Interchange of - condensates - matter content

$$\Omega(M, \Delta, \mu_I, \mu_{I5})$$

$$M \longleftrightarrow \Delta, \qquad \nu \longleftrightarrow \nu_5$$

$$\Omega(M, \Delta, \mu_I, \mu_{I5}) = \Omega(\Delta, M, \mu_{I5}, \mu_I)$$

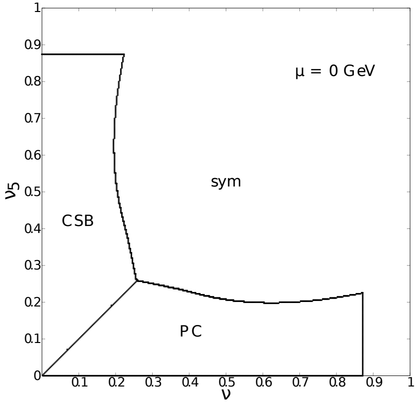


Figure: NJL model results

$$\mathcal{D}: \; M \longleftrightarrow \Delta, \quad \nu \longleftrightarrow \nu_5$$

Duality between chiral symmetry breaking and pion condensation

$$\text{PC} \longleftrightarrow \text{CSB} \quad \nu \longleftrightarrow \nu_5$$

Duality, Chiral limit

$$\mathcal{L}_{\text{QCD}} = \sum_{f=u,d} \bar{q}_f (i\not{D} - m_f) q_f - \frac{1}{4} \mathcal{G}_{\mu\nu,a} \mathcal{G}^{\mu\nu}_a .$$

$$\mathcal{L}_{\text{NJL}} = \sum_{f=u,d} \bar{q}_f \Big[\mathrm{i}\gamma^\nu \partial_\nu - m_f \Big] q_f + \frac{G}{N_c} \Big[(\bar{q}q)^2 + (\bar{q}\mathrm{i}\gamma^5 \vec{\tau} q)^2 \Big]$$

m_f is current quark masses

In the chiral limit $m_f = 0$ the Duality $\mathcal{D}$ is exact

m_f : $\quad \frac{m_u + m_d}{2} \approx 3.5\text{MeV}$
In our case typical values of $\mu, \nu, ..., T, .. \sim 10 - 100s$ MeV, for example, 200-400 MeV
One can work in the chiral limit $m_f \to 0$
$m_f = 0 \quad \to \quad m_\pi = 0$
physical m_f a few MeV $\quad \to \quad$ physical $m_\pi \sim 140$ MeV

Duality between CSB and PC is **approximate** in **physical point**

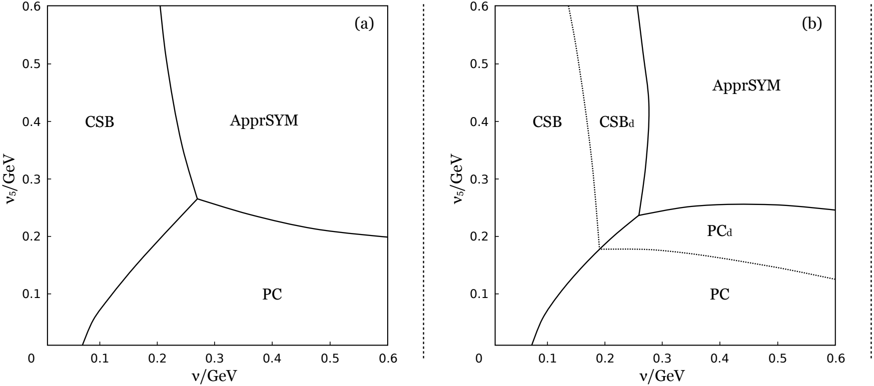


Figure: (ν, ν_5) phase diagram

Other Dualities

They are not that strong but still...

They could still be usefull

The TDP

$$\Omega(T, \mu, \nu, \nu_5, \mu_5;\ M, \Delta)$$

Let us assume that there is no PC

$$\Delta = 0$$

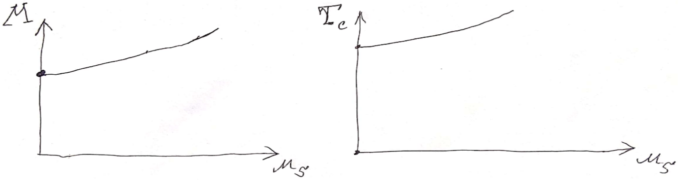
The TDP (phase diagram) is invariant under

$$\mu_5 \longleftrightarrow \nu_5$$

Chiral symmetry breaking phenomenon does not feel the difference between μ_5 and ν_5

Example of use of duality: catalysis of CSB

QCD at non-zero μ_5



catalysis of CSB by chiral imbalance:

- increase of $\langle \bar{q}q \rangle$ as μ_5 increases
- increase of critical temperature T_c of chiral phase transition (crossover) as μ_5 increases

M(ν_5) the same as M(μ_5)

all the results can be obtained by duality only (no PC)

catalysis of CSB by chiral isospin imbalance:

- increase of $\langle \bar{q}q \rangle$ as μ_{I5} increases
- increase of critical temperature T_c of chiral phase transition (crossover) as μ_{I5} increases

The TDP

$$\Omega(T, \mu, \nu, \nu_5, \mu_5;\ M, \Delta)$$

Let us assume that there is no CSB

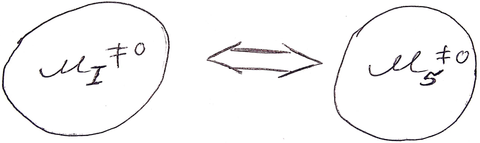
$$M = 0$$

The TDP (phase diagram) is invariant under

$$\mu_5 \longleftrightarrow \nu$$

Pion condensation phenomenon does not feel the difference between ν and μ_5

Two completely different systems



It was shown that chiral imbalance generates pion condensation in dense matter $n_B \neq 0$

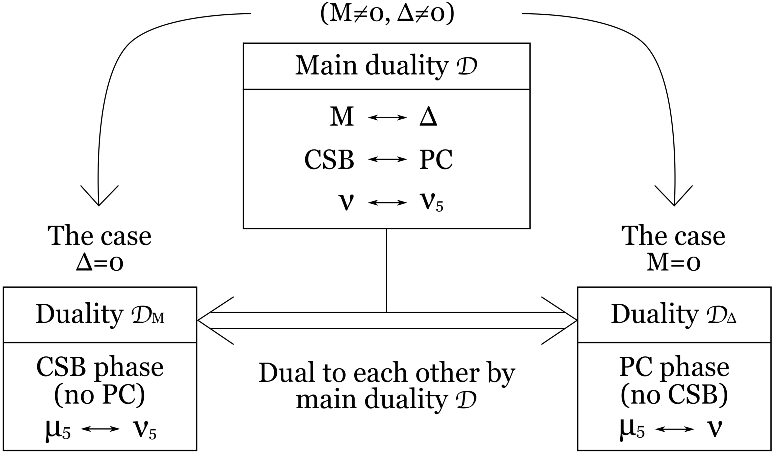
$\mu_5 \rightarrow$ **PC with** $n_B \neq 0$

- ν **and** μ_5 **has the same effect on PC**
- $\nu \rightarrow$ **PC**

So it can be

$\mu_5 \rightarrow$ **PC and** $\mu_5 \rightarrow$ **PC with** $n_B \neq 0$

Dualities of the phase diagram



Dualities on the lattice

Dualities on the lattice

$(\mu_B, \mu_I, \mu_{I5}, \mu_5)$

$\mu_B \neq 0$ impossible on lattice but if $\mu_B = 0$

$\mu_B \neq 0$ impossible on lattice but if $\mu_B = 0$

- **QCD at** μ_5 —- (μ_5, T)

  V. Braguta, A. Kotov et al, ITEP lattice group

- **QCD at** μ_I —- (μ_I, T)

  G. Endrodi, B. Brandt et al, Emmy Noether junior research group, Goethe-University Frankfurt, Institute for Theoretical Physics ()

No lattice calculations at μ_{I5}

QCD at μ_{I5}

But

there is duality $\mathcal{D}_M:\ \mu_5 \longleftrightarrow \nu_5$ if $\Delta = 0$

at $\mu_I = 0$ there is no PC ($\Delta = 0$)

- $M(\mu_5) = M(\nu_5)$
- $T_c(\mu_5) = T_c(\nu_5)$

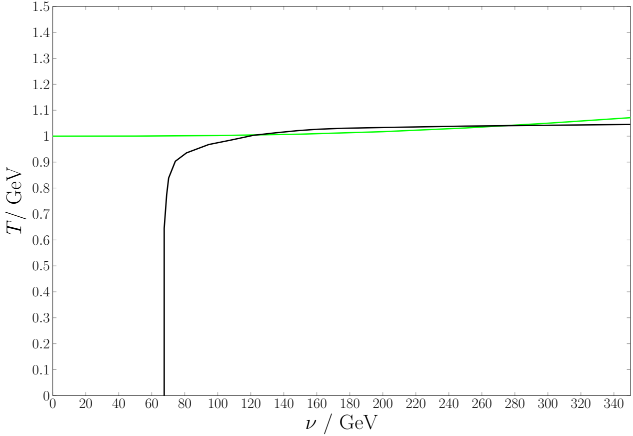
$$M \longleftrightarrow \Delta, \qquad \nu_5 \longleftrightarrow \nu$$

$$M \longleftrightarrow M \longleftrightarrow \Delta, \qquad \mu_5 \longleftrightarrow \nu_5 \longleftrightarrow \nu$$

So in this particular case you have a duality

$$M \longleftrightarrow \Delta, \qquad \mu_5 \longleftrightarrow \nu$$

- $M(\mu_5) = \Delta(\nu)$
- $T_c^M(\mu_5) = T_c^\Delta(\nu_5)$



T_c^M as a function of μ_5 (green line) and $T_c^\Delta(\nu)$ (black)

A number of papers predicted **anticatalysis** (T_c decrease with μ_5) of dynamical chiral symmetry breaking

A number of papers predicted **catalysis** (T_c increase with μ_5) of dynamical chiral symmetry breaking

lattice results show the **catalysis**

(ITEP lattice group, V. Braguta, A. Kotov, et al)

But unphysically large pion mass

Duality $\Rightarrow$ catalysis of chiral symmetry beaking

- **Large N$_c$ orbifold equivalences** connect gauge theories with different gauge groups and **matter content** in the large N$_c$ limit.

M. Hanada and N. Yamamoto, JHEP 1202 (2012) 138, arXiv:1103.5480 [hep-ph], PoS LATTICE **2011** (2011), arXiv:1111.3391 [hep-lat]

- two gauge theories with gauge groups G_1 and G_2

- but with different μ_1 and μ_2

Duality

$$G_1 \longleftrightarrow G_2, \qquad \mu_1 \longleftrightarrow \mu_2$$

or

$$\text{Phase structure } (G_1 \text{ at } \mu_1) \quad \longleftrightarrow \quad \text{Phase structure } (G_2 \text{ at } \mu_2)$$

Circumvent the sign problem

Duality

$$\text{QCD at } \mu_1 \longleftrightarrow \text{QCD at } \mu_2$$

- QCD with μ_2 —- sign problem free,
- QCD with μ_1 —- sign problem (no lattice)

Investigations of $(\text{QCD with } \mu_2)_{\text{on lattice}} \Longrightarrow (\text{QCD with } \mu_1)$

Inhomogeneous phases (case)

Homogeneous case

$$\langle \sigma(x) \rangle \text{ and } \langle \pi_a(x) \rangle$$

$$\langle \sigma(x) \rangle = M, \quad \langle \pi_\pm(x) \rangle = \Delta, \quad \langle \pi_3(x) \rangle = 0.$$

In vacuum the quantities $\langle\sigma(x)\rangle$ and $\langle\pi_a(x)\rangle$ do not depend on space coordinate x.

in a dense medium the ground state expectation values of bosonic fields might depend on spatial coordinates

CDW ansatz for CSB
the single-plane-wave LOFF ansatz for PC

$$\langle\sigma(x)\rangle = M\cos(2kx^1), \quad \langle\pi_3(x)\rangle = M\sin(2kx^1),$$
$$\langle\pi_1(x)\rangle = \Delta\cos(2k'x^1), \quad \langle\pi_2(x)\rangle = \Delta\sin(2k'x^1)$$

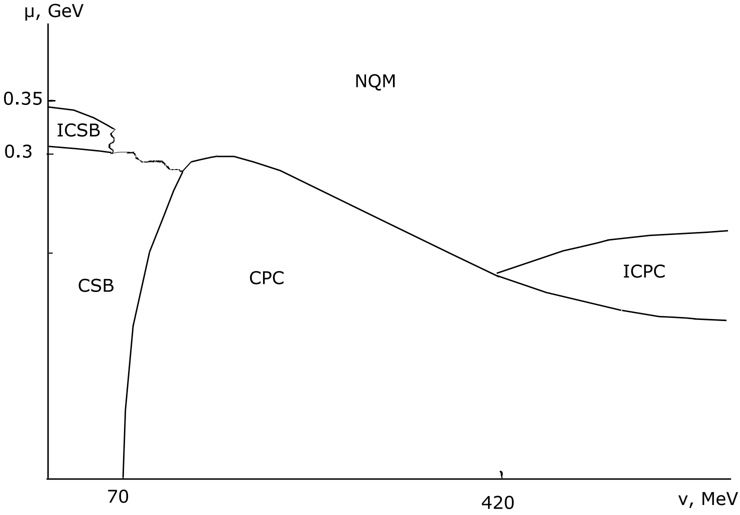
equivalently

$$\langle\pi_\pm(x)\rangle = \Delta e^{\pm 2k'x^1}$$

Duality in inhomogeneous case is shown

$$\mathcal{D}_I : \qquad M \longleftrightarrow \Delta, \quad \nu \longleftrightarrow \nu_5, \quad k \longleftrightarrow k'$$

schematic (ν, μ)-phase diagram



Combined schematic (ν, μ)-phase diagram.

schematic (ν_5, μ)-phase diagram

- exchange axis ν to the axis ν_5,
- rename the phases ICSB $\leftrightarrow$ ICPC, CSB $\leftrightarrow$ CPC, and NQM phase stays intact here

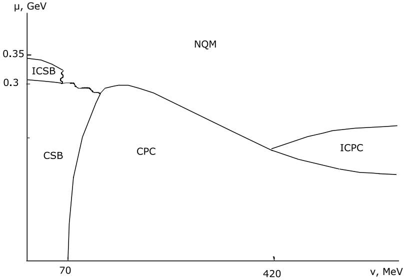


Figure: (ν, μ)-phase diagram

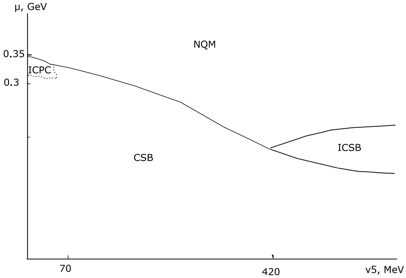


Figure: (ν_5, μ)-phase diagram

Duality between CSB and PC was found in

- In the framework of NJL model
- In the large N_c approximation (or mean field)
- In the chiral limit

QCD Lagrangian is

$$\mathcal{L} = i\bar{\psi}\gamma^{\mu}D_{\mu}\psi + \bar{\psi}\Big[\mu\gamma^0 + \frac{\mu_I}{2}\tau_3\gamma^0 + \frac{\mu_{I5}}{2}\tau_3\gamma^0\gamma^5 + \mu_5\gamma^0\gamma^5\Big]\psi$$

$$\mathcal{D}: \qquad \psi_R \to i\tau_1\psi_R$$

$$\mu_I \leftrightarrow \mu_{I5}$$

$$\bar{\psi}\psi \leftrightarrow i\bar{\psi}\gamma^5\tau_1\psi$$

$$M \longleftrightarrow \Delta, \qquad \nu \longleftrightarrow \nu_5, \quad \mu_I \longleftrightarrow \mu_{I5}$$

$$i\bar{\psi}^C\sigma_2\tau_2\gamma^5\psi \leftrightarrow i\bar{\psi}^C\sigma_2\tau_2\gamma^5\psi, \qquad \bar{\psi}^C\sigma_2\tau_2\psi \leftrightarrow \bar{\psi}^C\sigma_2\tau_2\psi$$

$$\bar{\psi}\tau_2\psi \leftrightarrow \bar{\psi}\tau_3\psi, \qquad \bar{\psi}\tau_1\psi \leftrightarrow i\bar{\psi}\gamma^5\psi, \qquad i\bar{\psi}\gamma^5\tau_2\psi \leftrightarrow i\bar{\psi}\gamma^5\tau_3\psi$$

$$\mathcal{D} \in \quad SU_R(2) \quad \in SU_L(2) \times SU_R(2)$$

$$\mu_I \leftrightarrow \mu_{I5}$$

$M \neq 0$ breaks the chiral symmetry

Duality $\mathcal{D}$ is a remnant of chiral symmetry

$$\tilde{\mathcal{D}} \in \quad SU_R(2) \times U_A(1)$$

$$\mu_I \leftrightarrow \mu_{I5}$$

$U_A(1)$ is anomalous

$$\tilde{\mathcal{D}} \in \quad SU_R(2) \times U_A(1)$$

$$\mu_I \leftrightarrow \mu_{I5}$$

$U_A(1)$ is anomalous

The NJL Lagrangian

$$\mathcal{L} = i\bar{\psi}\gamma^\mu\partial_\mu\psi + G_1\Big\{(\bar{\psi}\psi)^2 + (i\bar{\psi}\vec{\tau}\gamma^5\psi)^2\Big\} + G_2\Big\{(i\bar{\psi}\gamma^5\psi)^2 + (\bar{\psi}\vec{\tau}\psi)^2\Big\}$$

$$\bar{\psi}\psi \leftrightarrow \bar{\psi}\tau_1\psi, \quad \bar{\psi}\tau_1\psi \leftrightarrow \bar{\psi}\psi, \quad i\bar{\psi}\gamma^5\tau_1\psi \leftrightarrow i\bar{\psi}\gamma^5\psi, \quad \bar{\psi}\tau_3\psi \leftrightarrow i\bar{\psi}\gamma^5\tau_2\psi$$

The transformation should be

$$\mu_I \leftrightarrow \mu_{I5}, \qquad G_1 \leftrightarrow G_2 \tag{1}$$

$$\mathcal{L} = i\bar{\psi}\gamma^\mu\partial_\mu\psi + G_1\Big\{(\bar{\psi}\psi)^2 + (i\bar{\psi}\vec{\tau}\gamma^5\psi)^2\Big\} + G_2\Big\{(i\bar{\psi}\gamma^5\psi)^2 + (\bar{\psi}\vec{\tau}\psi)^2\Big\}$$

$$+H_1(i\bar{\psi}\sigma_2\lambda_2\gamma^5\psi^C)(i\bar{\psi}^C\sigma_2\lambda_2\gamma^5\psi) + H_2(\bar{\psi}\sigma_2\lambda_2\psi^C)(\bar{\psi}^C\sigma_2\lambda_2\psi)$$

$$\bar{\psi}\psi \leftrightarrow \bar{\psi}\tau_1\psi, \quad \bar{\psi}\tau_1\psi \leftrightarrow \bar{\psi}\psi, \quad i\bar{\psi}\gamma^5\tau_1\psi \leftrightarrow i\bar{\psi}\gamma^5\psi, \quad \bar{\psi}\tau_3\psi \leftrightarrow i\bar{\psi}\gamma^5\tau_2\psi$$

$$|\bar{\psi}^C\sigma_2\lambda_2\psi|^2 \leftrightarrow |i\bar{\psi}^C\sigma_2\lambda_2\gamma^5\psi|^2$$

The transformation should be

$$\mu_I \leftrightarrow \mu_{I5}, \qquad G_1 \leftrightarrow G_2, \qquad H_1 \leftrightarrow H_2$$

$$\text{QCD:} \qquad SU_L(2) \times SU_R(2)$$

$$\text{QC}_2\text{D:} \qquad SU(4)$$

$$\mathcal{D} \in SU(4)$$

$$\mu \leftrightarrow \nu, \qquad \Delta \leftrightarrow \Delta_{CSC}$$

$$\bar{\psi}\psi \to \bar{\psi}\psi,\ i\bar{\psi}\gamma^5\tau_3\psi \to i\bar{\psi}\gamma^5\tau_3\psi,\ i\bar{\psi}\gamma^5\psi \to i\bar{\psi}\gamma^5\psi,\ \bar{\psi}\tau_3\psi \to \bar{\psi}\tau_3\psi$$

$$i\bar{\psi}\gamma^5\tau_1\psi \to \frac{-i\bar{\psi}\sigma_2\gamma^5\tau_2\psi^C - i\bar{\psi^C}\sigma_2\gamma^5\tau_2\psi}{2}$$

$$i\bar{\psi}\gamma^5\tau_2\psi \to \frac{-\bar{\psi}\sigma_2\gamma^5\tau_2\psi^C + \bar{\psi^C}\sigma_2\gamma^5\tau_2\psi}{2}$$

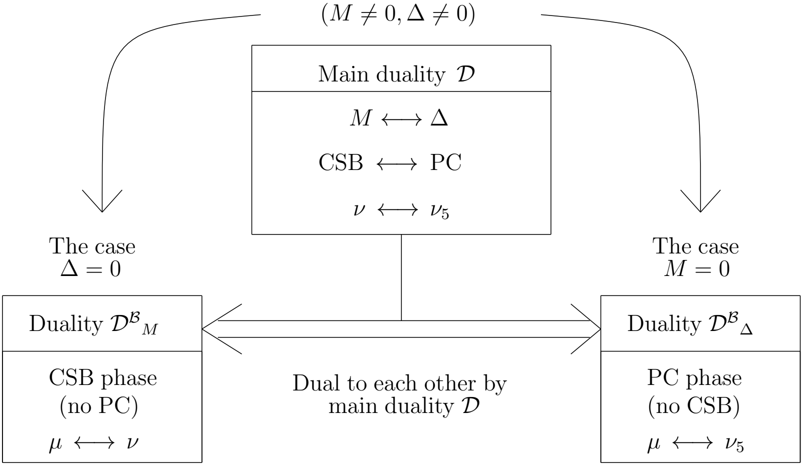
$$\mathcal{D} \in SU(4)$$

$$\mu \leftrightarrow \nu_5, \qquad M \leftrightarrow \Delta_{CSC}$$

$$\bar{\psi}\psi \to \frac{-i\bar{\psi}\sigma_2\gamma^5\tau_2\psi^C - i\bar{\psi^C}\sigma_2\gamma^5\tau_2\psi}{2},\ i\bar{\psi}\gamma^5\tau_3\psi \to \frac{-\bar{\psi}\sigma_2\gamma^5\tau_2\psi^C + \bar{\psi^C}\sigma_2\gamma^5\tau_2\psi}{2}$$

Dualities concerning baryon density

They could be even more usefull


$(M \neq 0, \Delta \neq 0)$
Main duality $\mathcal{D}$
$M \longleftrightarrow \Delta$
CSB $\longleftrightarrow$ PC
$\nu \longleftrightarrow \nu_5$
The case $\Delta = 0$
The case $M = 0$
Duality ${\mathcal{D}^{\mathcal{B}}}_M$
CSB phase (no PC)
$\mu \longleftrightarrow \nu$
Dual to each other by main duality $\mathcal{D}$
Duality ${\mathcal{D}^{\mathcal{B}}}_\Delta$
PC phase (no CSB)
$\mu \longleftrightarrow \nu_5$

CSB phenomenon
does not feel the difference between μ and ν
if there is no pion condensation phenomenon

in finite volume there cannot be a breaking of continuous symmetry

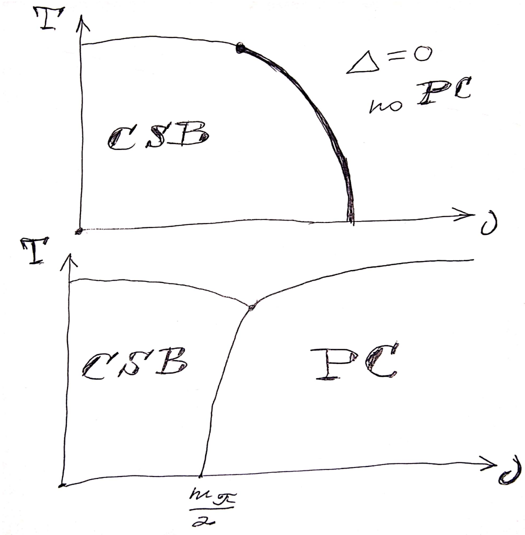
on lattice to get pion condensation (breaking of $U_{\tau_3}(1)$) one adds the pion source λ

CSB phenomenon
does not feel the difference between μ and ν
if there is no pion condensation phenomenon

probably

if $\lambda = 0$ then there is no symmetry breaking and pion condensation

it is unphysical and there should be (M, Δ)
but artificially one can probe the local minimum $(M, 0)$


T
Δ=0
no PC
CSB
T
CSB
PC

Let us assume that $\nu \neq 0$
as a rule there is PC and there is no CSB

$$M = 0, \qquad \mu \longleftrightarrow \nu_5$$

PC at non-zero μ $\longleftrightarrow$ PC at non-zero ν_5

Conclusions

- $(\mu_B, \mu_I, \nu_5, \mu_5)$ phase diagram was studied
  PC *in dense matter with chiral imbalance*
- It was shown that there exist dualities
- There have been shown ideas how dualities can be used
- lattice QCD support the idea of duality
- Duality is **not just entertaining mathematical property** but an **instrument with very high predictivity power**
- (μ_B, ν_5) phase diagram is **quite rich** and contains various **inhomogeneous phases**
- It was shown that the duality holds not only in large N_c limit and in NJL model but in QCD

Thanks for the attention