

Black Hole Hair:

Quantum and Classical

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Skyrmion Black Hole
Hair: Conservation of
Baryon Number by Black
Holes

with: Alexander
Gubmann

Folk theorem:

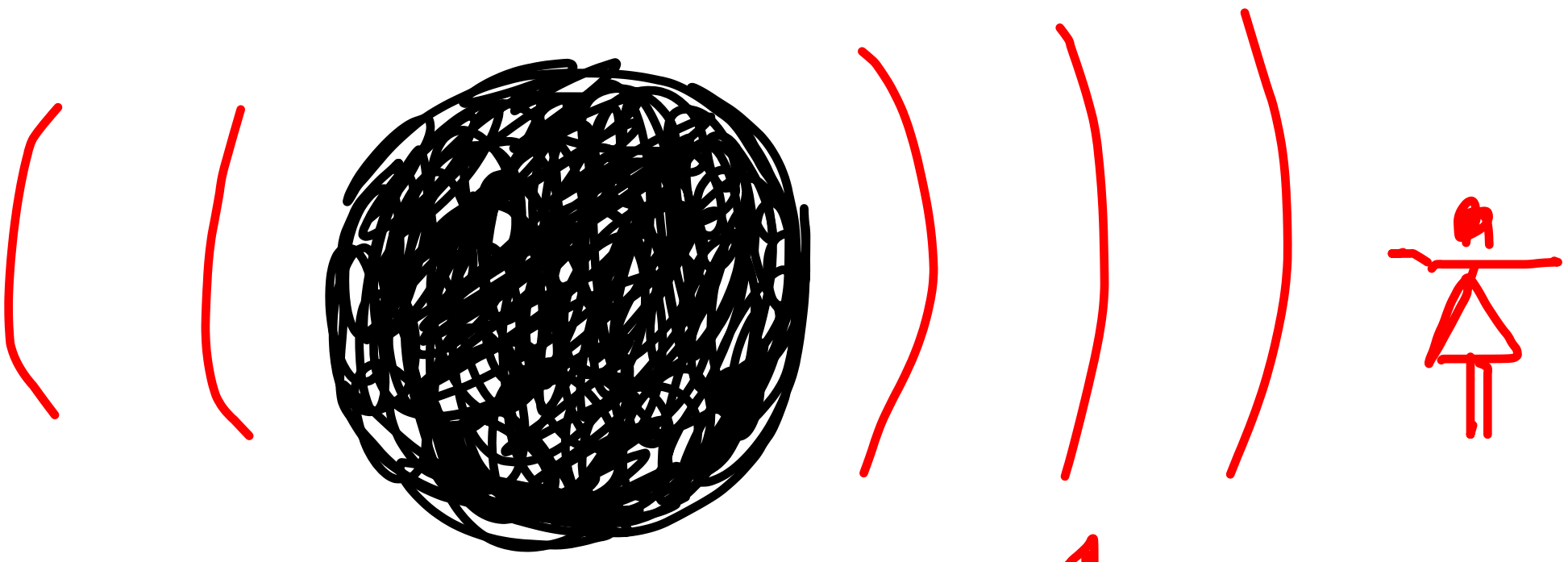
Semi-classical black
hole physics is incompatible
with conserved baryon
number!

We shall find a loophole
in the "proof".

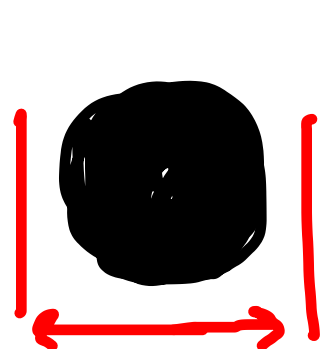
Standard argument:



black hole with baryons
in it



Thermal radiation



$L_p \approx$ Planck length

Assumptions:

- ⊗ No barionic hair;
- ⊗ No deviation from thermality until

$$r_h \sim L_P$$

↑
Black hole horizon

Baryons as Skyrmions

QCD vacuum (take 2
massless quark flavors)

$$\begin{pmatrix} u \\ d \end{pmatrix} \langle \bar{u}u \rangle = \langle \bar{d}d \rangle \neq 0$$

Spontaneous breaking
of symmetry

$$SU(2)_L \otimes SU(2)_R \rightarrow SU(2)_V$$

3 Nambu-Goldstone bosons:

pions $\pi^a \quad a=1,2,3$

Thus, the low energy theory below the QCD scale Λ_{QCD} is a

Chiral theory of

$$U \equiv \exp\left(i \frac{\mathcal{J}^a}{F_\pi} \sigma^a\right)$$

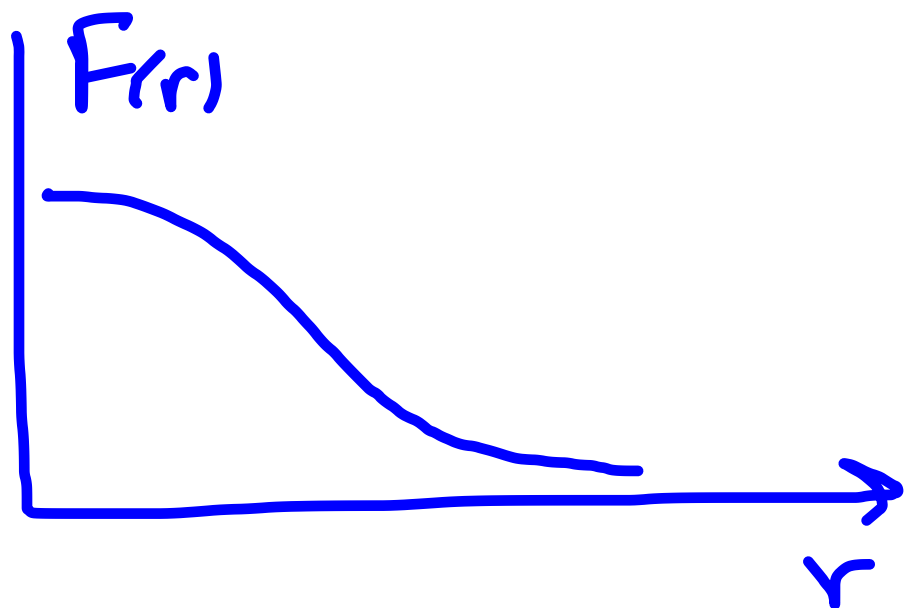
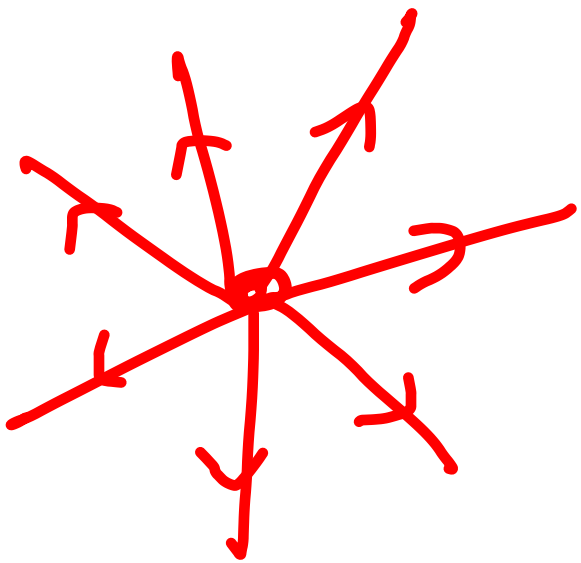
Pauli matrices

Where are the baryons?

The vacuum is S^3
and there are topological
defects (solitons) Skyrmions

$$\frac{1}{F_\pi} \nabla^a = F(r) \hat{n}^a$$

$$\hat{n} \equiv \frac{\vec{r}}{r}$$



Skyrmion hair.

$$L_S = -\frac{F_\pi^2}{4} \text{Tr} (U^\dagger \partial_\mu U U^\dagger \partial^\mu U) \\ + \frac{1}{e^2} \text{Tr} ([\partial_\mu U U^\dagger, \partial_\nu U U^\dagger]^2)$$

where

$$U = e^{i \frac{\vec{T} \cdot \vec{\phi}}{f_\pi}} \leftarrow \underline{SU(2)}.$$

$$\vec{T}^a \leftarrow \text{Pions}$$

Skyrmion size:

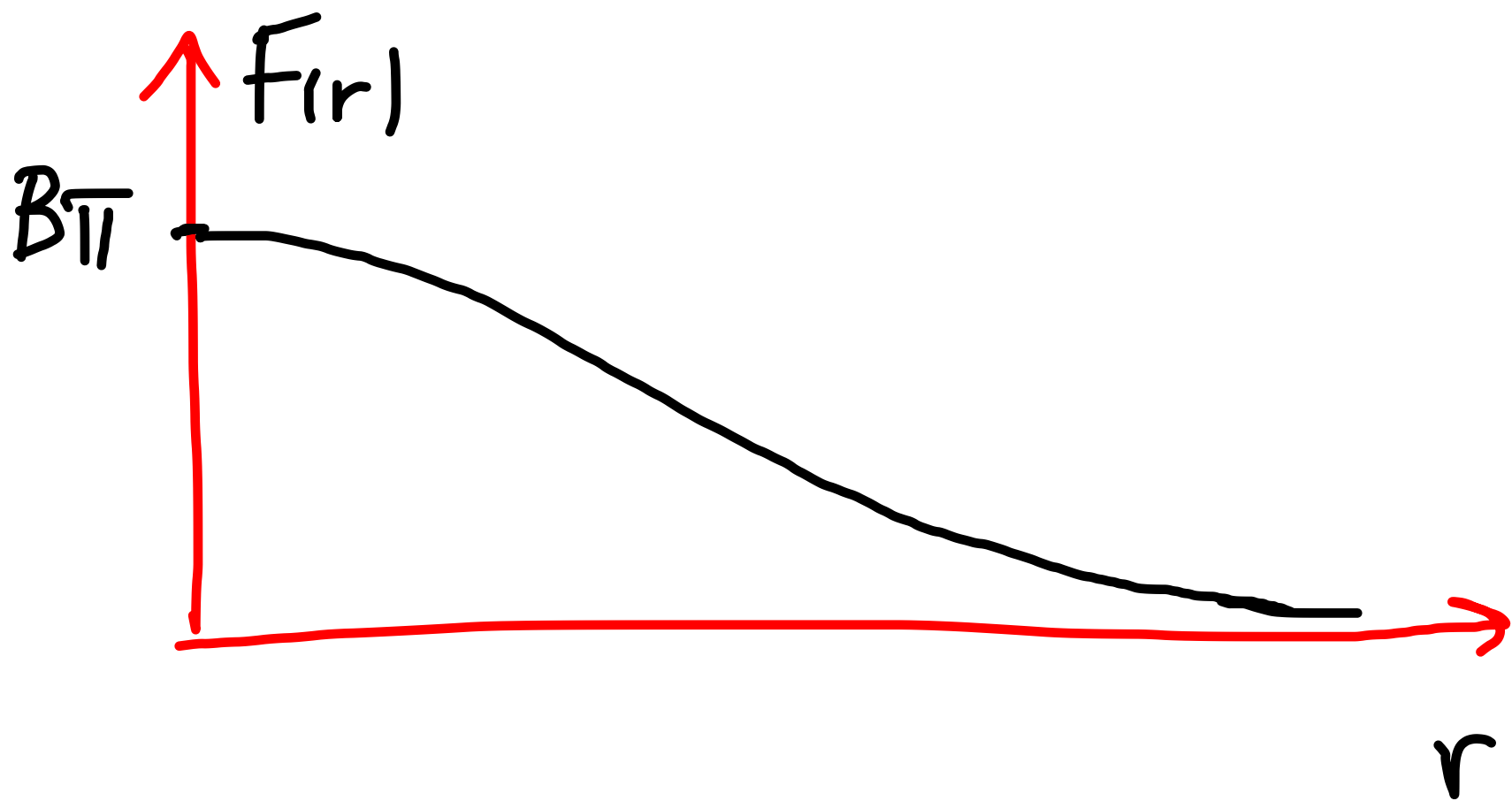
$$L = \frac{1}{e F_{\pi}}$$

Skyrmion Mass

$$M_s = \frac{F_{\pi}}{e}$$

Skyrmion:

$$\frac{J^a}{F_\pi} = F(r) n^a$$



$$F(0) = B\pi, \quad F(\infty) = 0$$

Skyrm topological Chern-
- Simons current

$$J_{\mu} \equiv -\epsilon_{\mu\nu\alpha\beta} \text{Tr} \left(\bar{U}^{-1} \partial^{\nu} U \bar{U}^{-1} \partial^{\alpha} U \right. \\ \left. U^{-1} \partial^{\beta} U \right)$$

Topological charge

$$B = \int d^4x J_0$$

Black holes with
skyrmion hair have
been found for

$$x_h \lesssim L$$



Black holes with Skyrmion
hair (Luckock, Moss; Droz,
Heusler, Straumann; Bizon,
Chmaj; Shiiki, Nawado; ...)

$$ds^2 = N(r)^2 \left(1 - \frac{2M(r)G_N}{r} \right) dt^2 -$$
$$- \left(1 - \frac{2M(r)G_N}{r} \right)^{-1} dr^2 - r^2 d\Omega^2$$

$$N(\infty) = 1 \quad \bar{F}(\infty) = 0$$

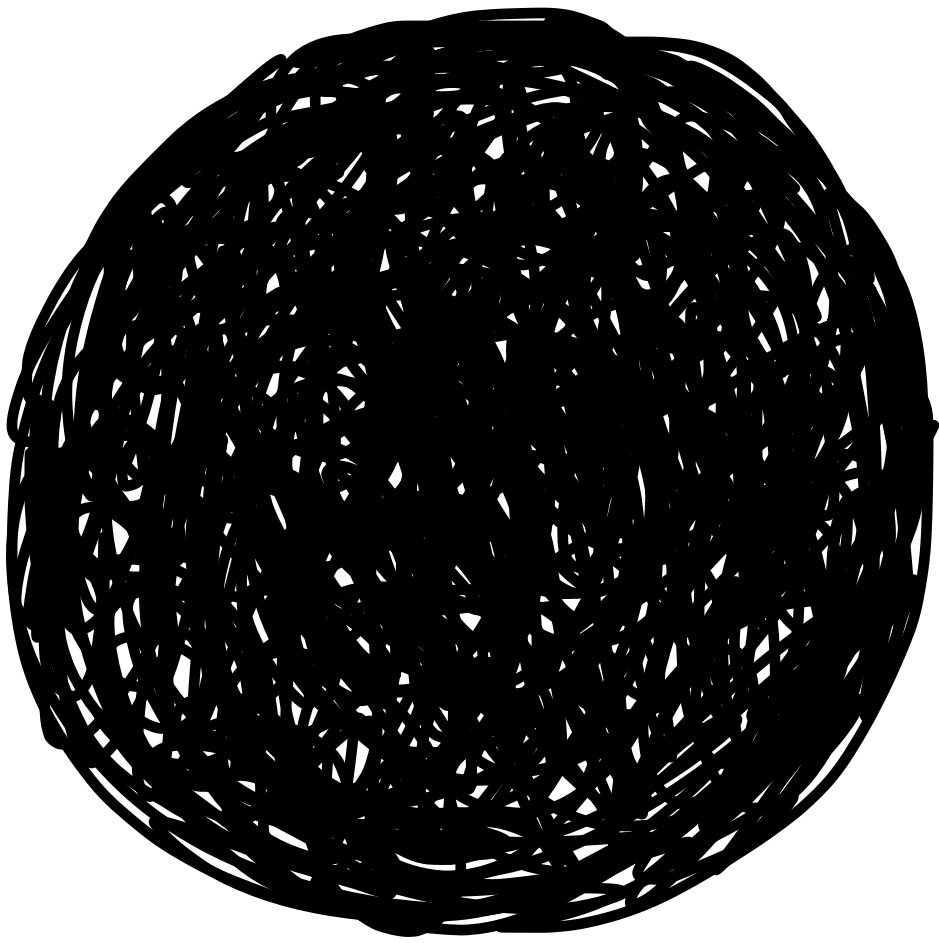
$$F(r) \propto \frac{1}{r^2} \quad r \rightarrow \infty$$



Black
hole

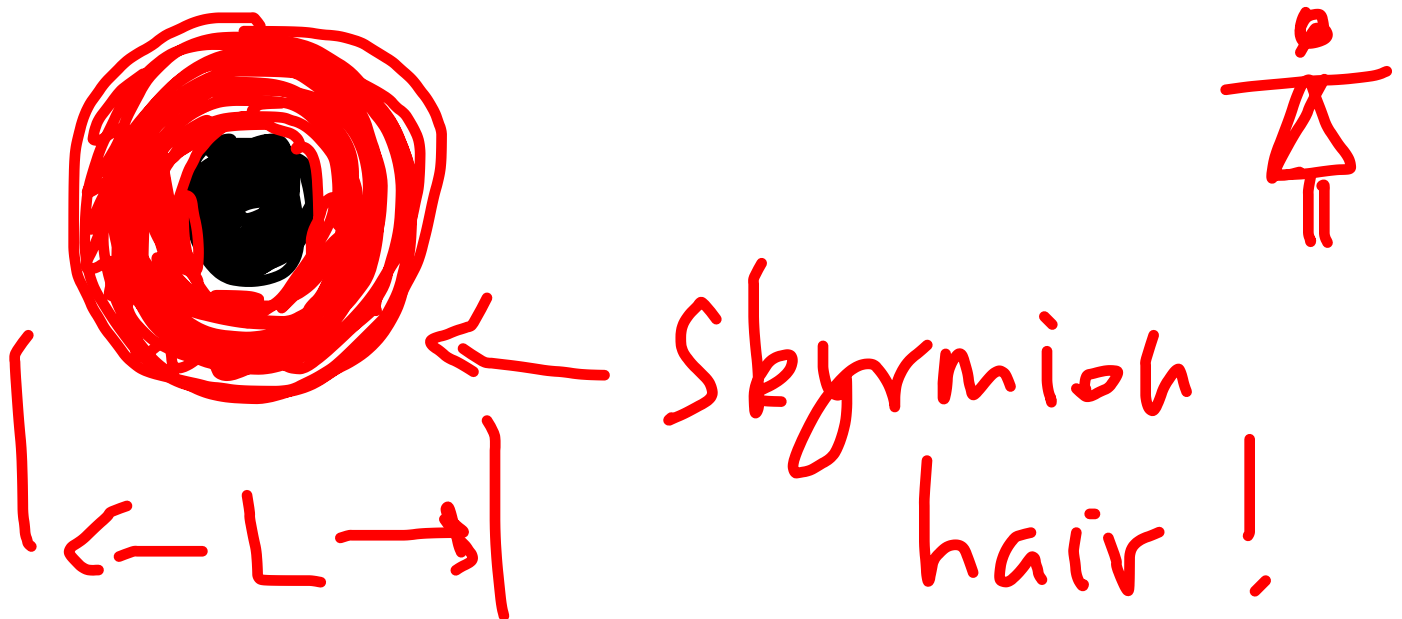
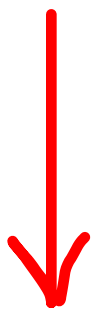
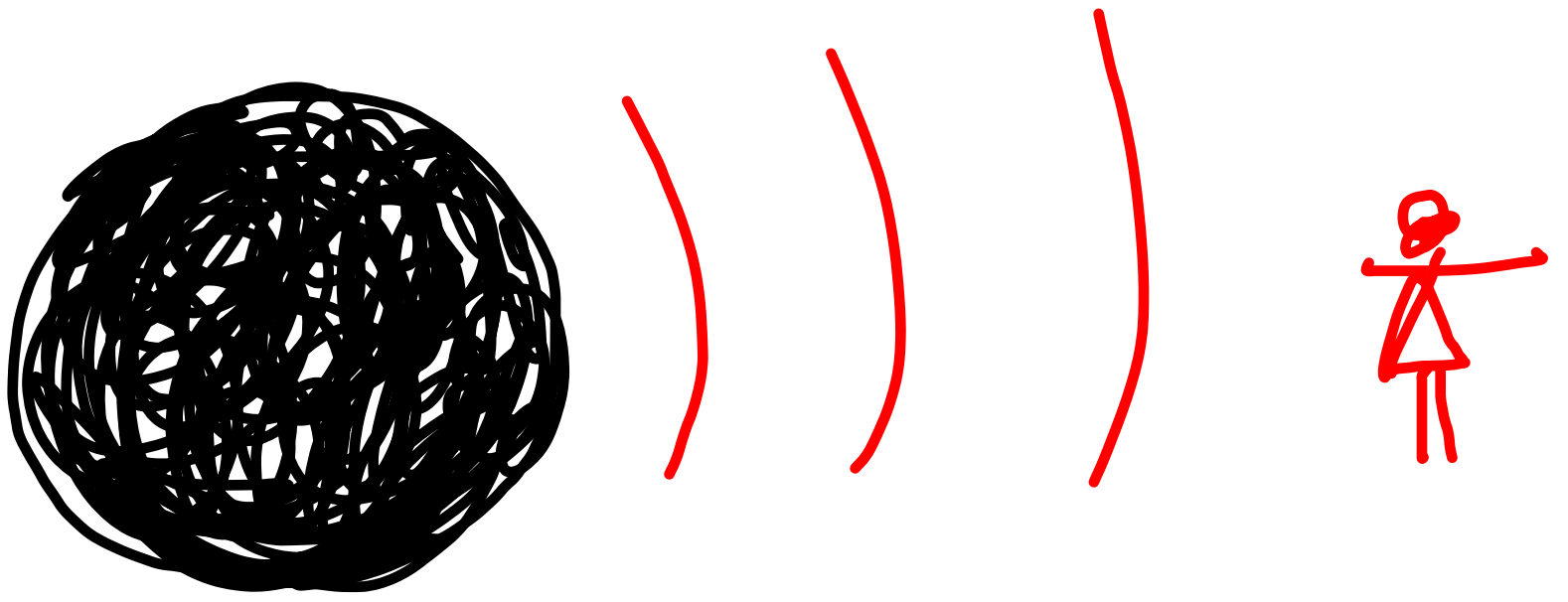
$$\chi_h < L$$

Classical
Skyrmion hair



$$\chi_h \gg L$$

This opens up a logical possibility



Is this only a logical
possibility or is a
must?

What is the implication
for baryon number?

Skyrmion - Baryon
correspondence in QCD
 $SU(N_c)$. Witten

$$U(1)_B \quad q \rightarrow e^{i\alpha} q$$

Baryon current $J_\mu = \frac{1}{N_c} \bar{q} \gamma_\mu q$

Skyrmion current

$$J_\mu = -\epsilon_{\mu\nu\alpha\beta} \text{Tr} (\bar{U} \partial^\nu U \bar{U} \partial^\alpha U \bar{U} \partial^\beta U)$$

Notia

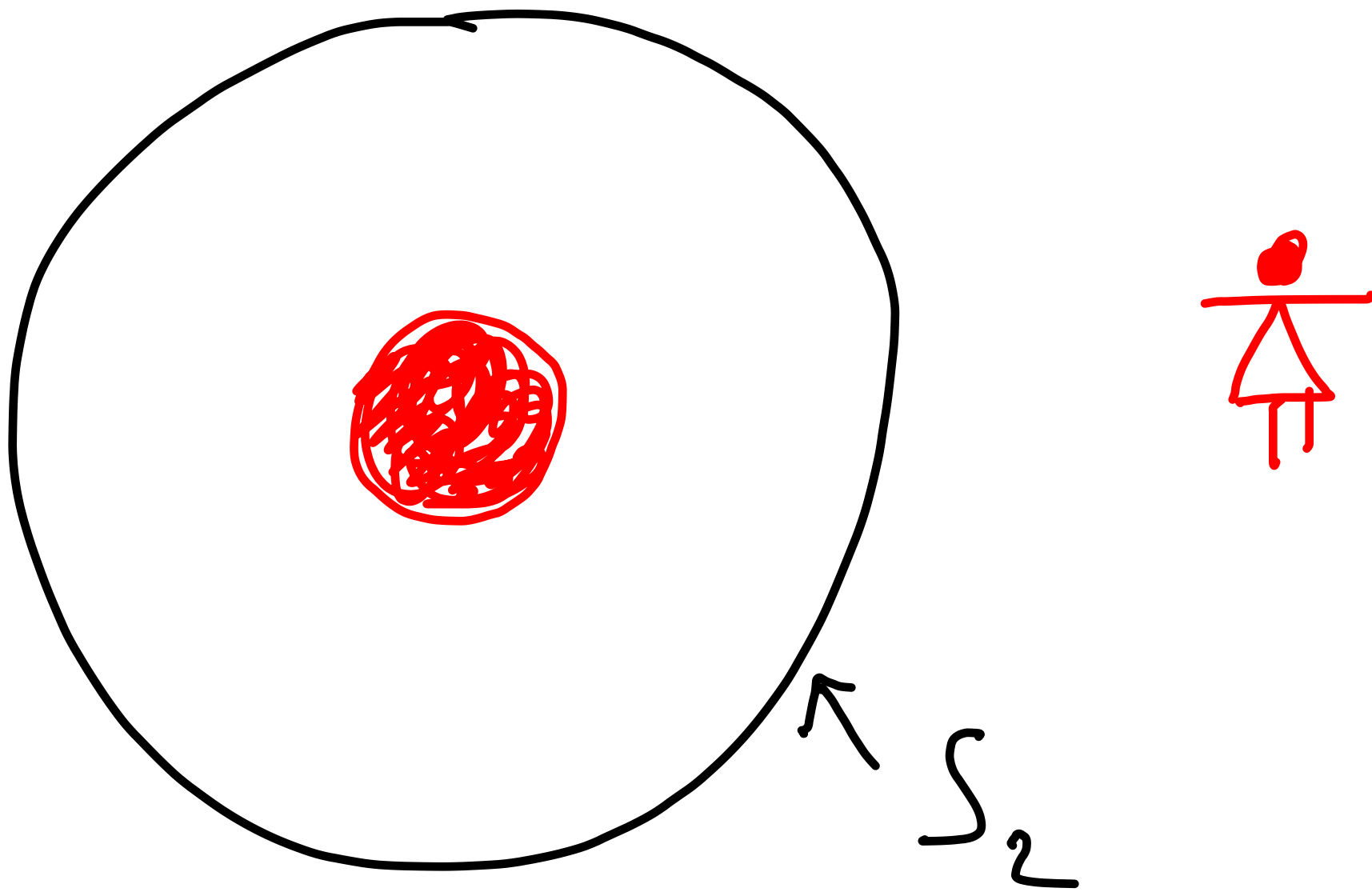
$$J_0 = \int d^3x \mathcal{L}$$

where

$$\mathcal{L} \equiv - \left(F(r) - \frac{1}{2} \sin(2F(r)) - \pi \right) \cdot \partial_\mu \theta \partial_\nu \phi$$

Thus,

$$B = \int d^3x J_0 = \int_{\Sigma_2} dx^\mu dx^\nu \mathcal{L}_{\mu\nu}$$



Thus, Skyrmion/baryon charge can be measured by a surface integral over S_2

$$B = \int_{S_2} dx^M \wedge dx^N S_{MN}$$

For this we couple $S_{\mu\nu}$
to a probe string

$$S = g \int dx^\mu dx^\nu S_{\mu\nu}$$

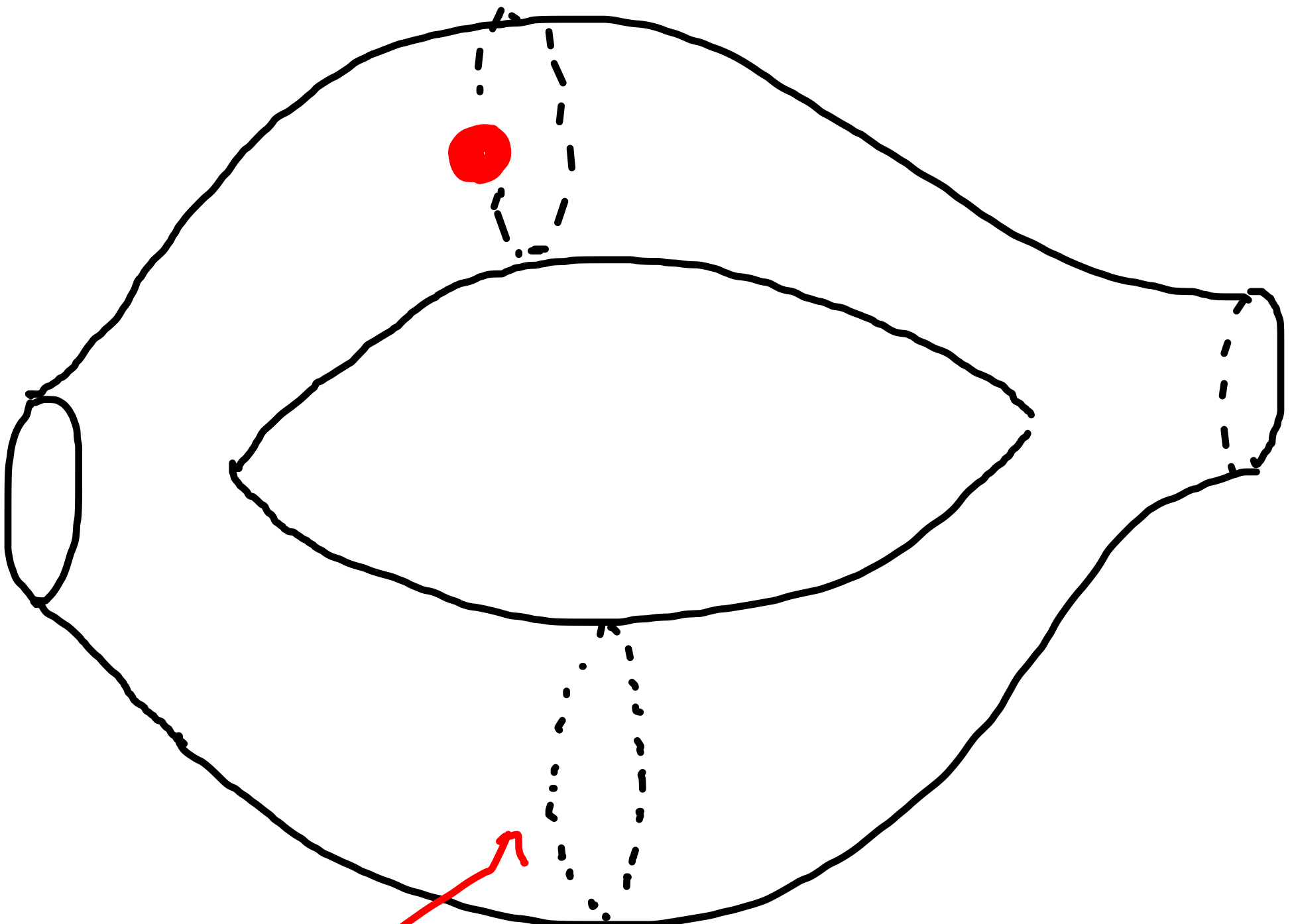
and perform Aharonov-
Bohm type interference
experiment.

Skyrmion



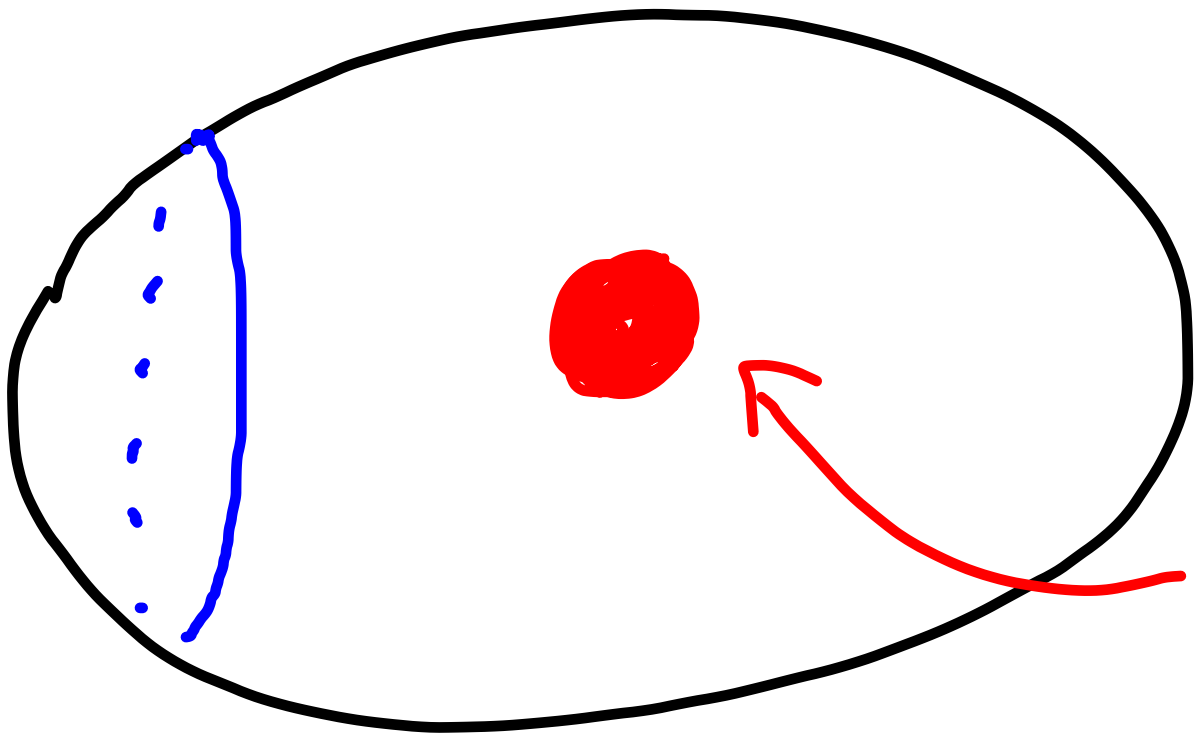
Aharonov-Bohm phase-shift

$$\Delta\Phi = 2\pi g_B$$

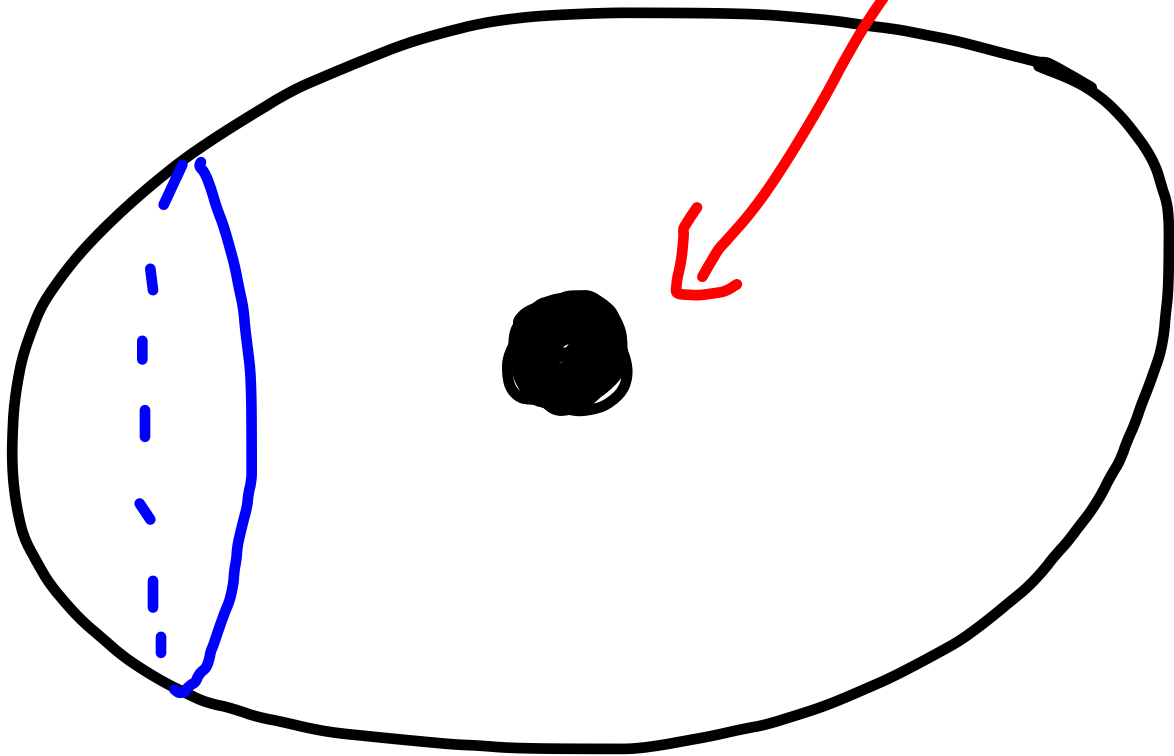
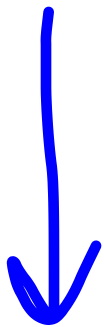


String loop

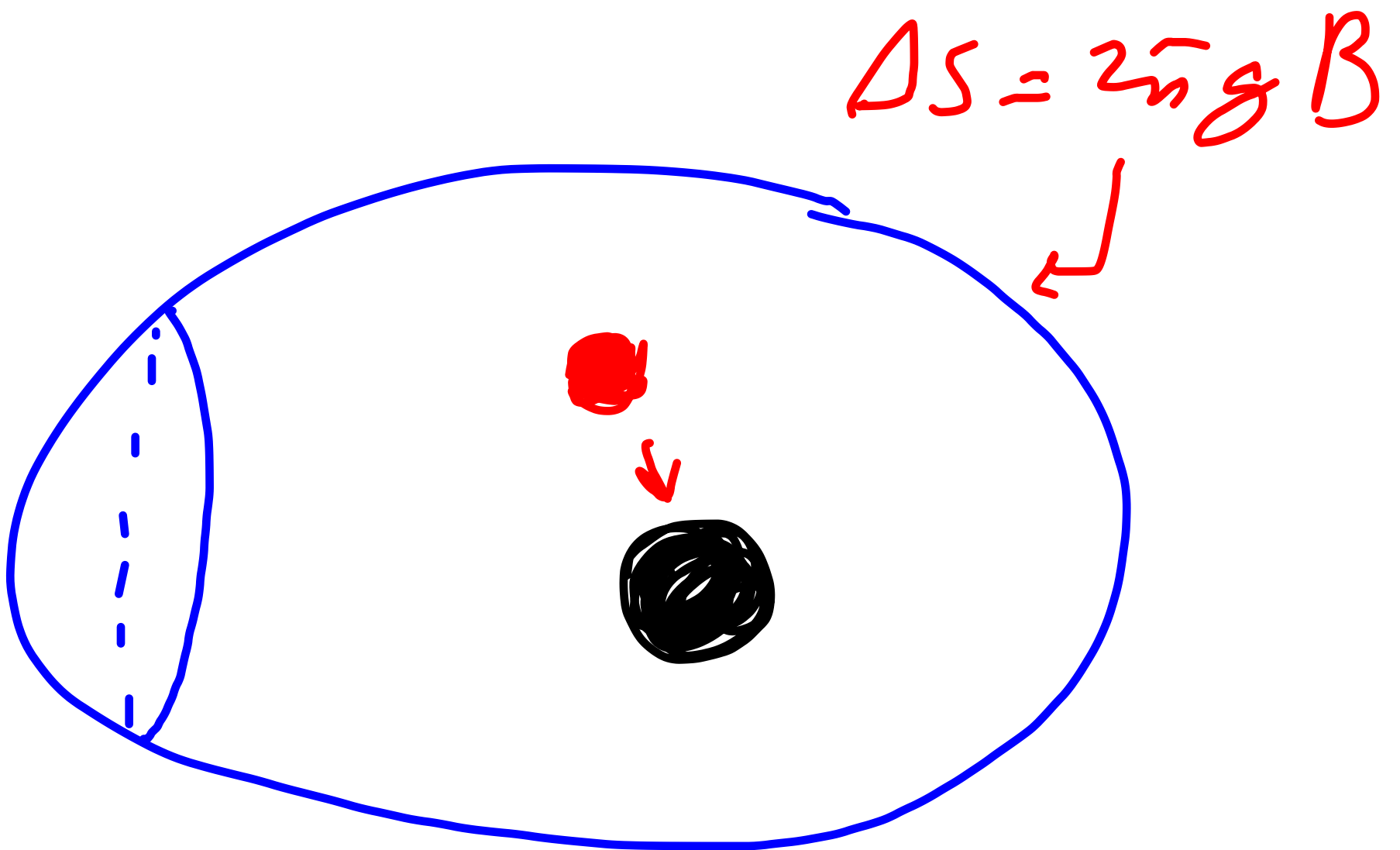
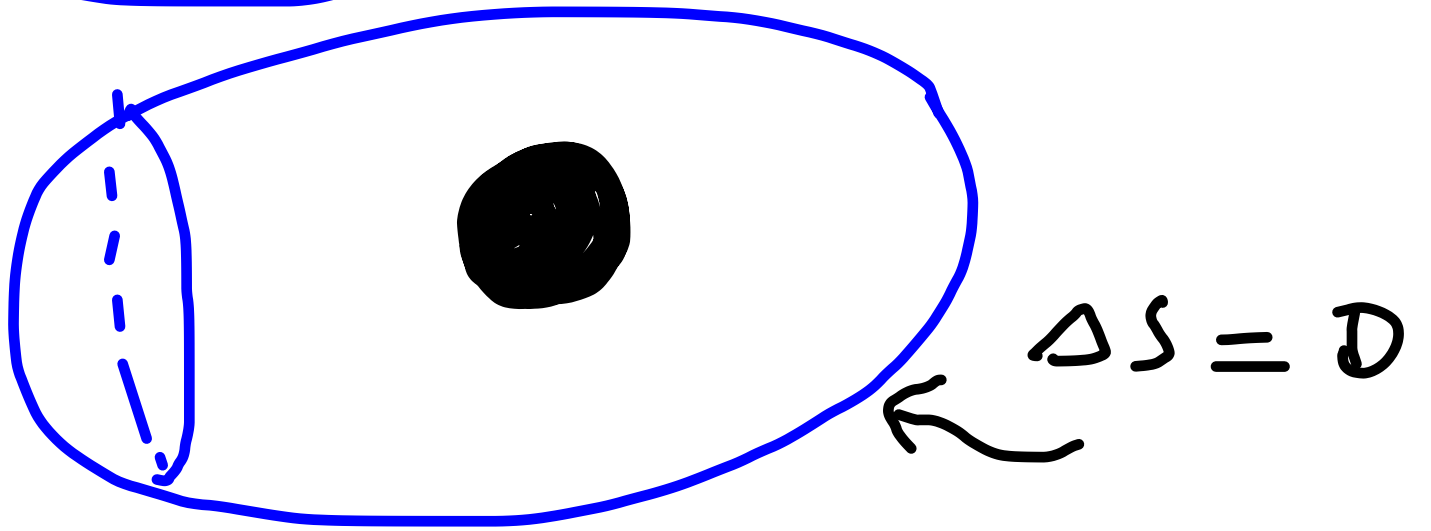
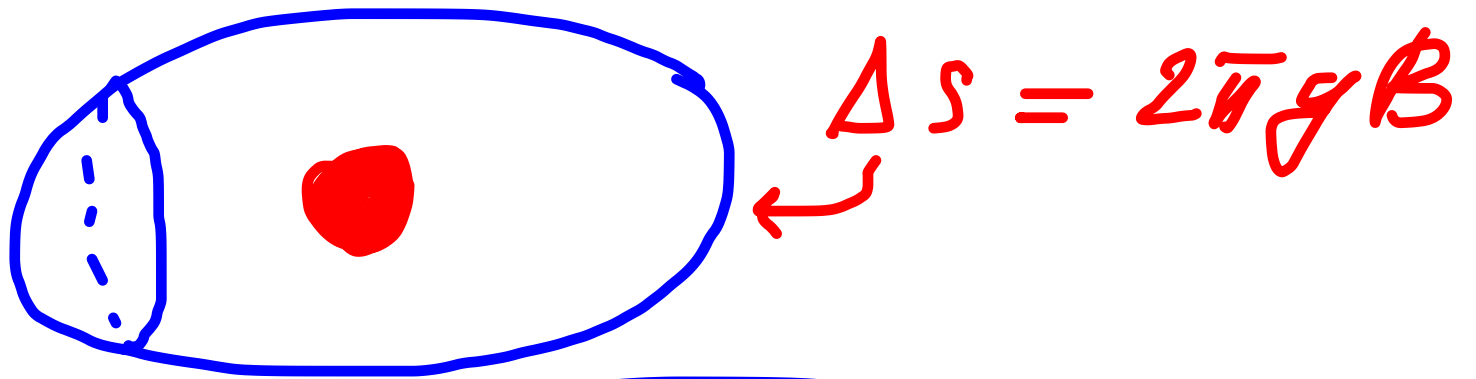




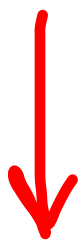
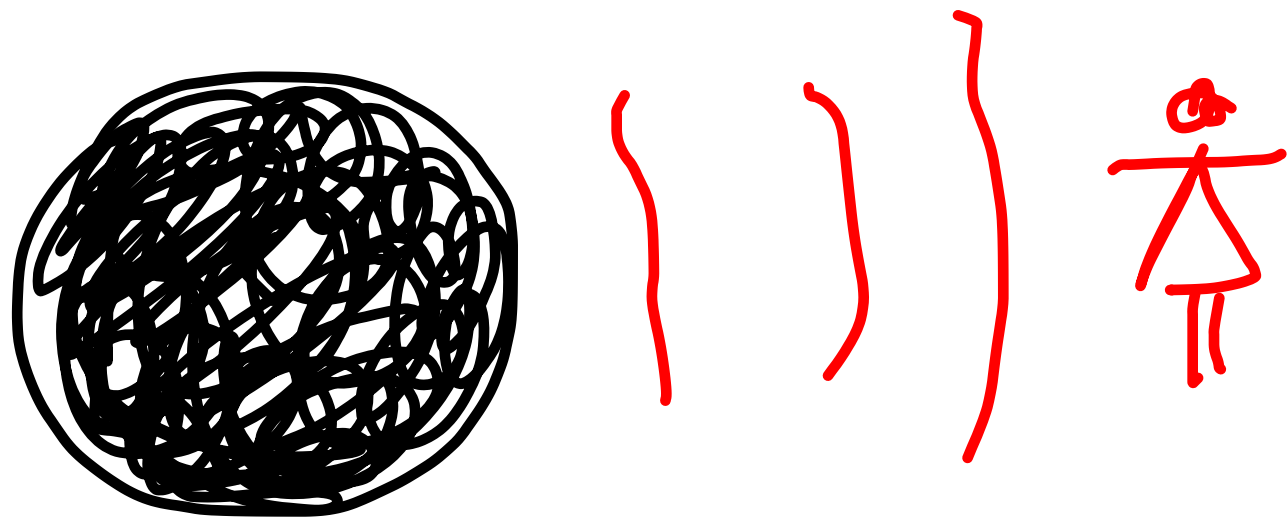
Skyrmion



Black hole



Since Alice can measure
AB-phase-shift she can
monitor Skyrminion/baryon
number at $r = \infty$.



Measurement of black hole
gauge hair via AB

phase shift was considered:

① For $B_{\mu\nu}$ (Bowick, Giddings,
Harvey, Horowitz and Strominger)

② For gauged Z_N (Krauss &
Wilczek; Pruski & Krauss)

③ For massive spin-2 $h_{\mu\nu}$
(Dvali)

In all these examples the
hair is carried by gauge field!

It thus appears that
Skyrmion 2-form $S_{\mu\nu}$
acts like a gauge potential
and skyrmion-baryon charge
acquires a secret topological
protection.

Obviously $B = \int d^3x J_0$

is invariant under a gauge
shift

$$S_{\mu\nu} \rightarrow S_{\mu\nu} + \partial_\mu \xi_\nu - \partial_\nu \xi_\mu$$

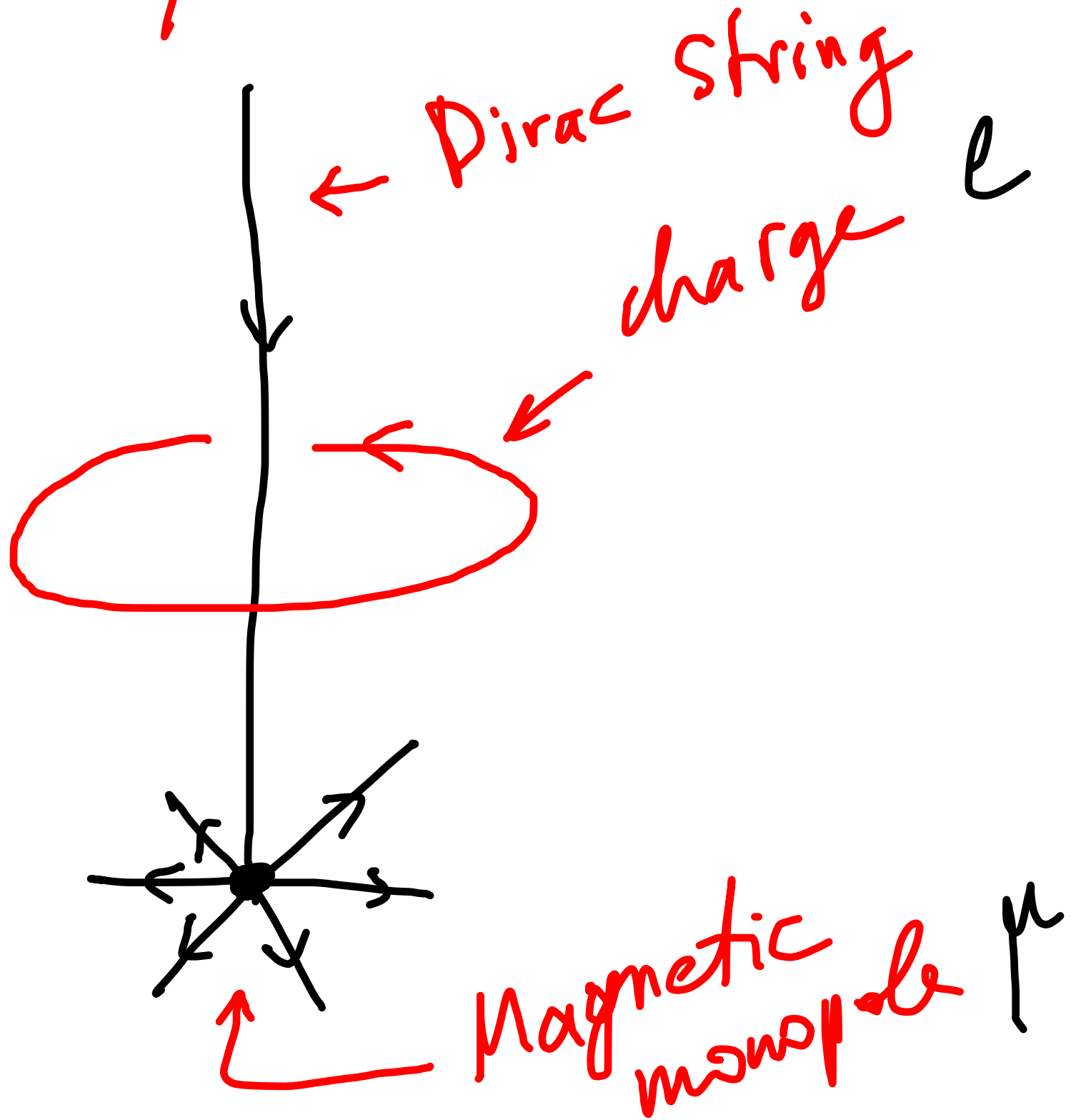
Physicality of $S_{\mu\nu}$ depends on the charge of a probe string.

$$S_{\text{string}} = g \int dX^M dX^\nu S_{\mu\nu}$$

We must have

$$g \neq \frac{\beta}{h}$$

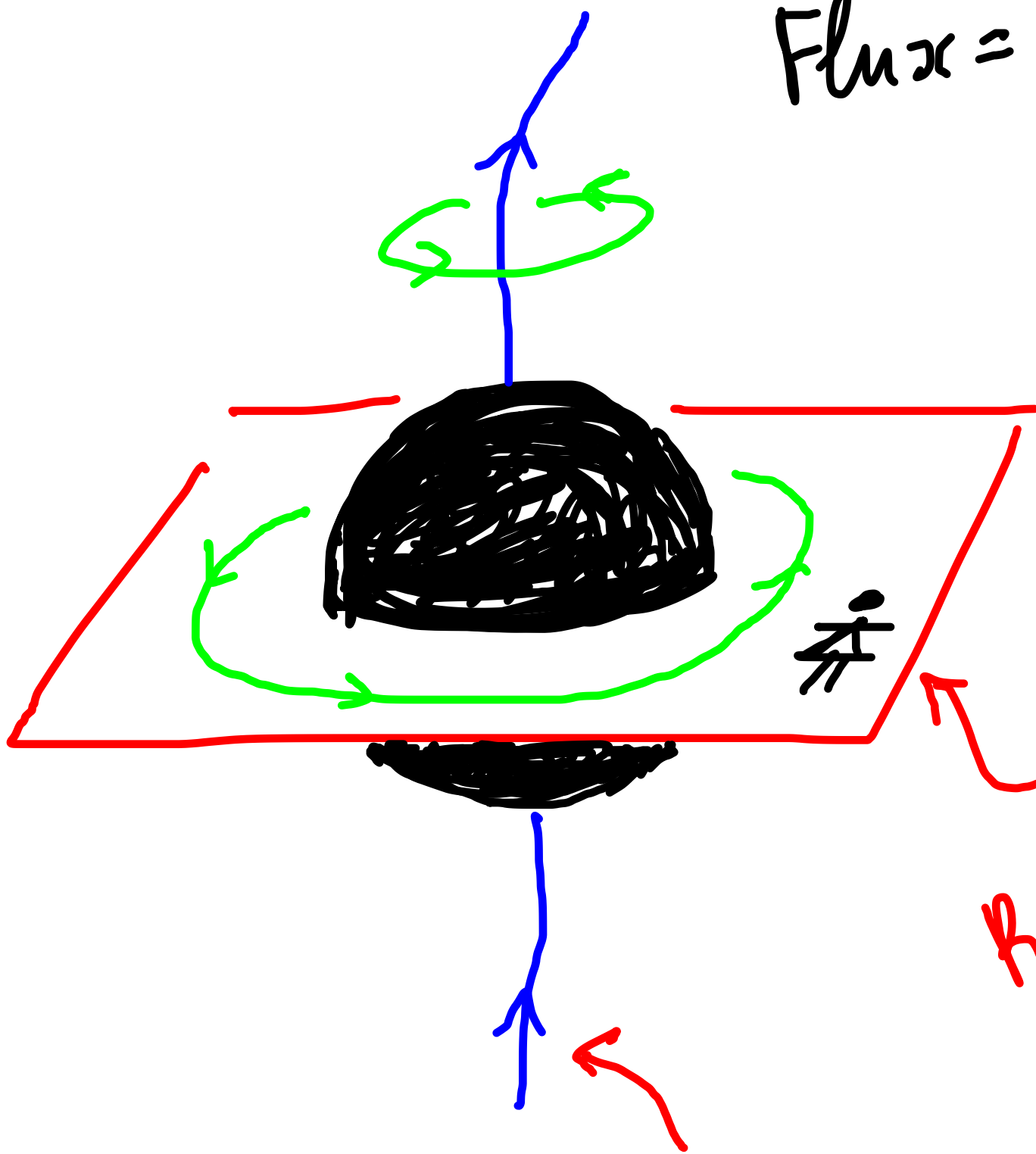
Dirac's quantization condition



Phase-shift $\Delta S = 2\pi \mu g$

Cosmic string AB-hair

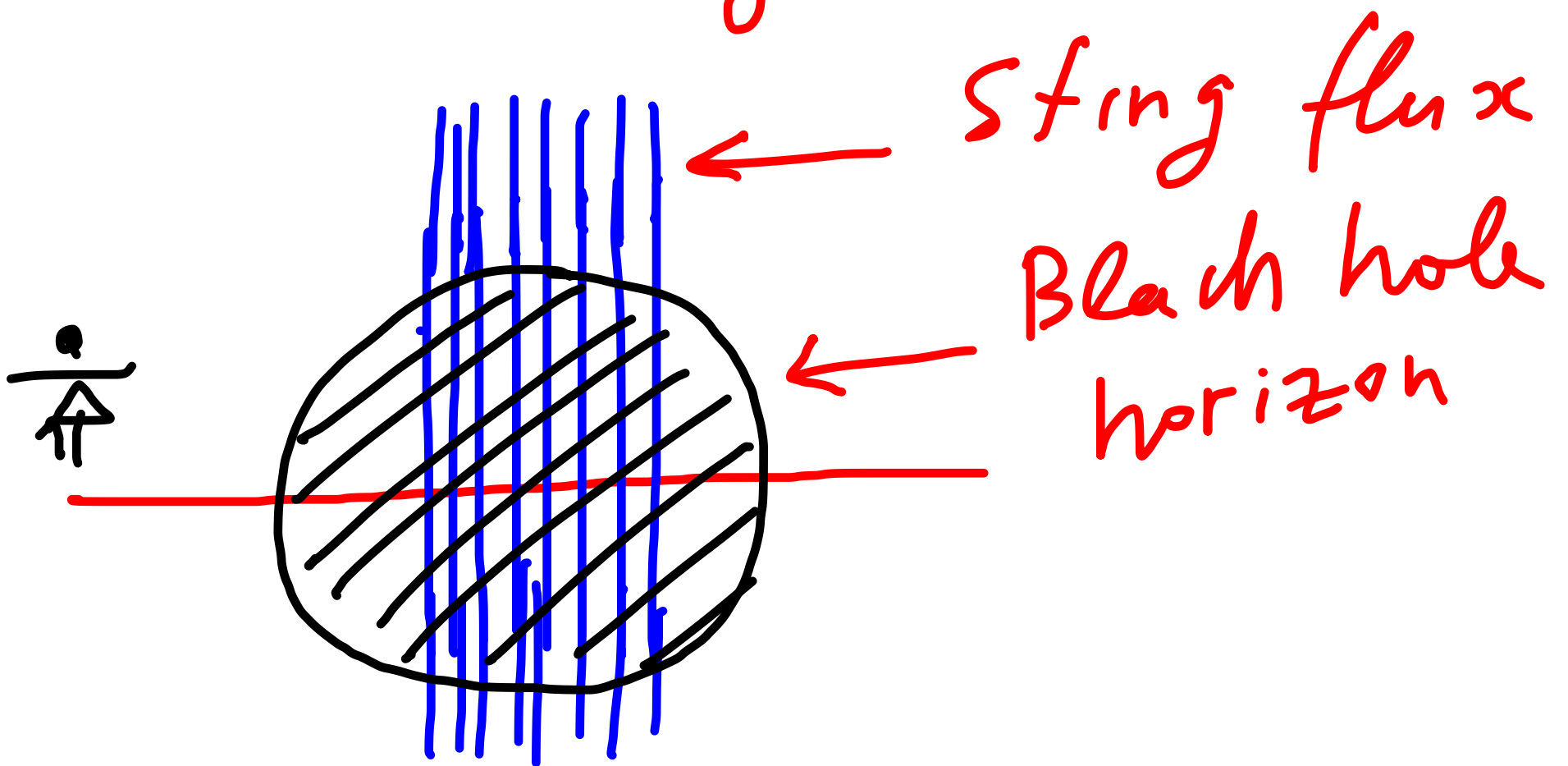
$$\text{Flux} = \oint A_\mu dx^\mu$$



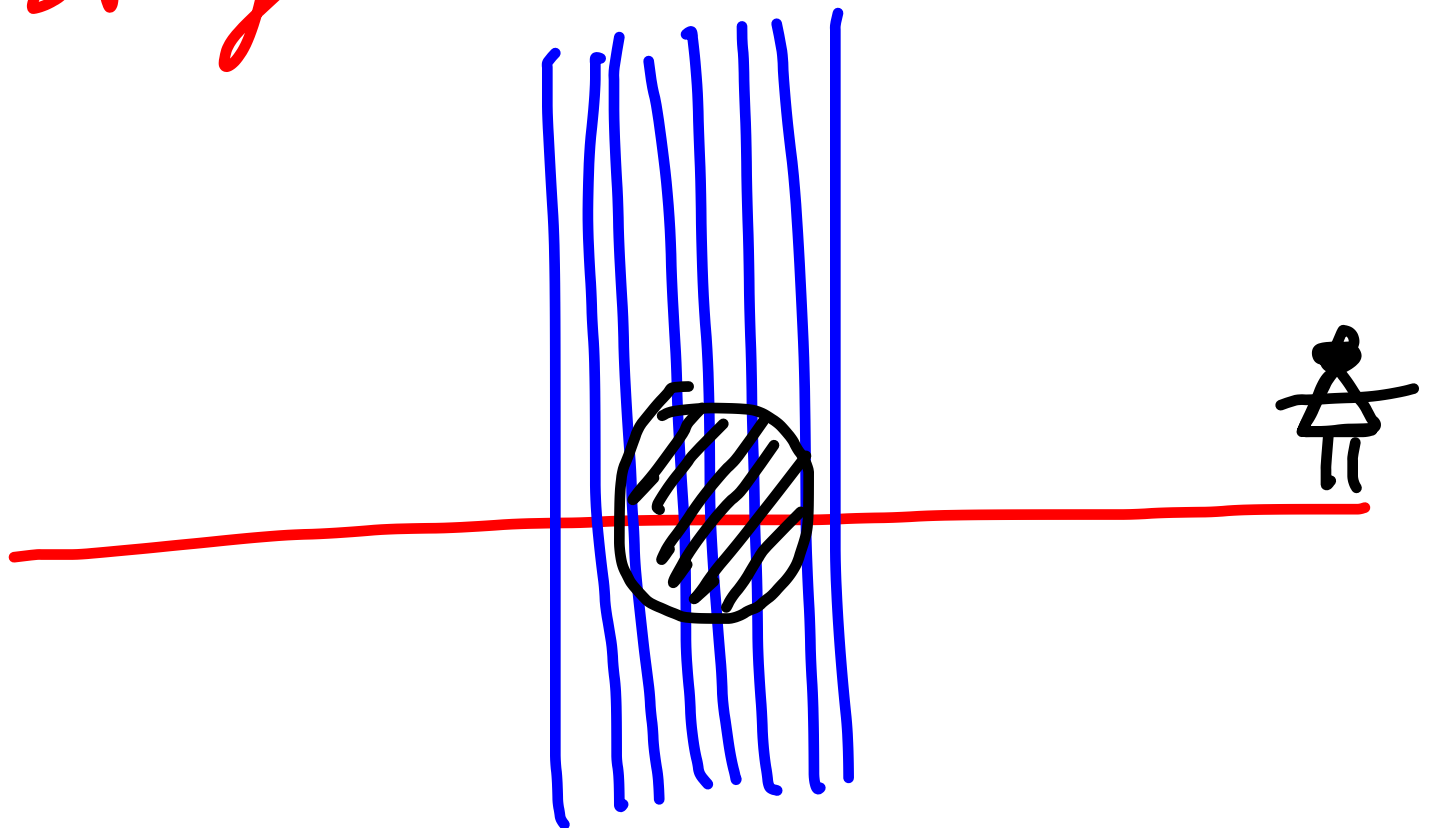
Alice's
brane world

Cosmic string
magnetic flux

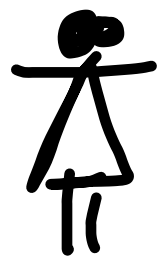
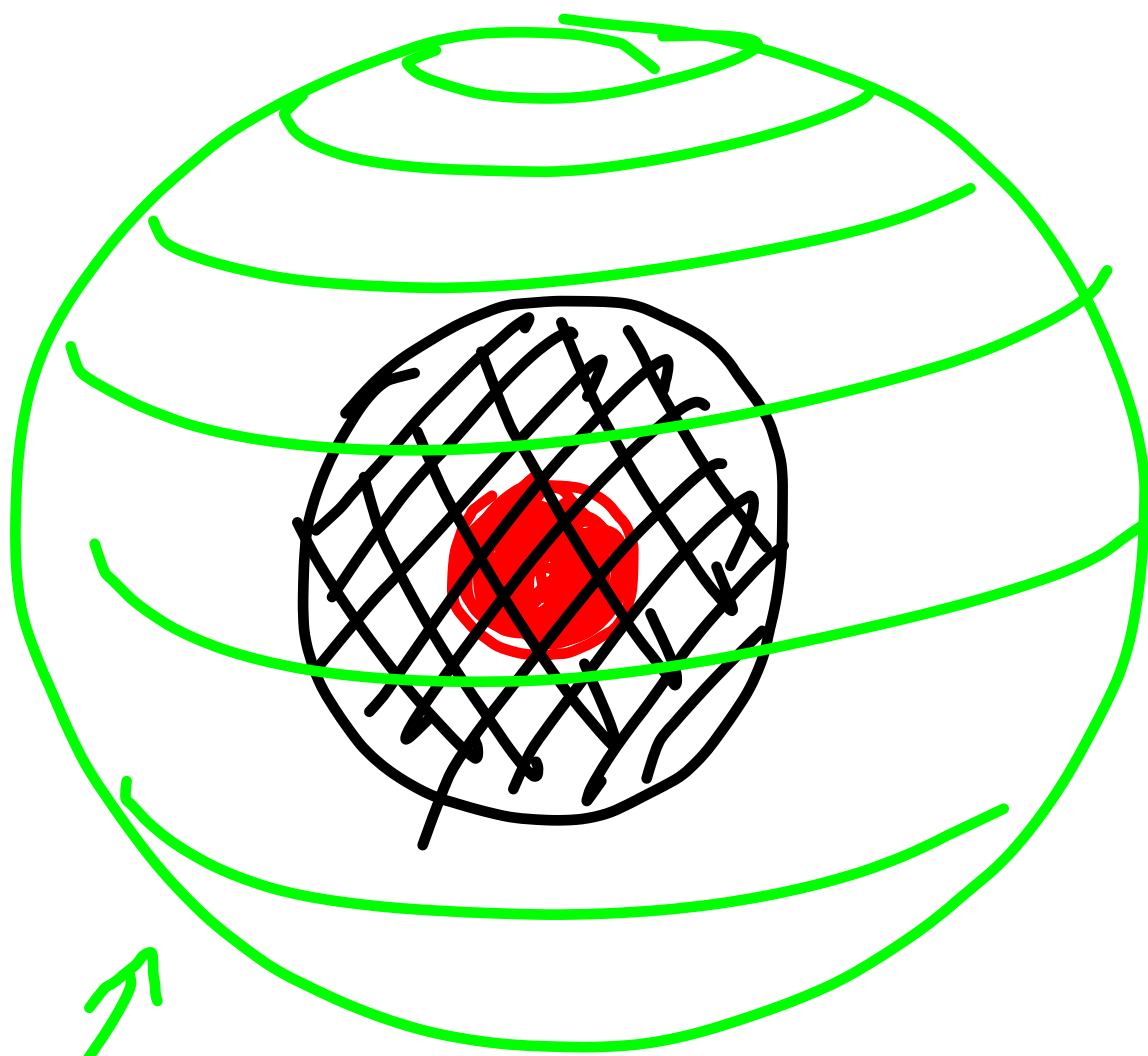
Initial stage



Later stage



For Skyrmion we are
in similar situation except
our "brane-world" is
3-dimensional



↖ $B = \int dx^0 dx^1 dx^2 S_{\mu\nu}$

So the baryon
number cannot be
lost!

Astrophysical complications?